

# Dark Continuum in the Spectral Function of the Resonant Fermi Polaron



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## Summary

- controlled numerical study of the ground state **spectral function** of the **resonant Fermi polaron** in 3d
- establish the existence of a **"dark continuum"** — a region of anomalously low spectral weight between the narrow polaron peak and the broad maximum at positive energy
- the dark continuum develops around  $k_F a \lesssim 1$ , i.e. in the absence of a small interaction-related parameter
- detailed analysis of what features of the spectrum can be recovered reliably (polaron peak, dark continuum) and what features are inaccessible (width and structure of the higher-frequency part of the spectrum)

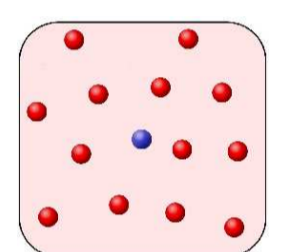
[1] O. Goulko, A. S. Mishchenko, N. Prokof'ev, B. Svistunov, Phys. Rev. A 94, 051605(R) (2016), eprint arXiv:1603.06963

[2] O. Goulko, A. S. Mishchenko, L. Pollet, N. Prokof'ev, B. Svistunov, Phys. Rev. B 95, 014102 (2017), eprint arXiv:1609.01260

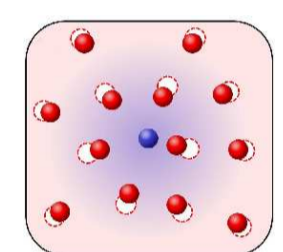
## The Fermi Polaron

A **single spin**  $\downarrow$  fermion interacting with a **sea of spin**  $\uparrow$  fermions

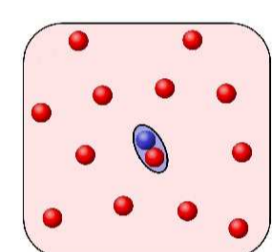
- extremal case of an imbalanced Fermi mixture
- impurity problem
- direct realization of a quasiparticle



weak attraction:  
mean-field of the  
medium



stronger attraction:  
polaron "dressed" with  
cloud of spin  $\uparrow$  atoms

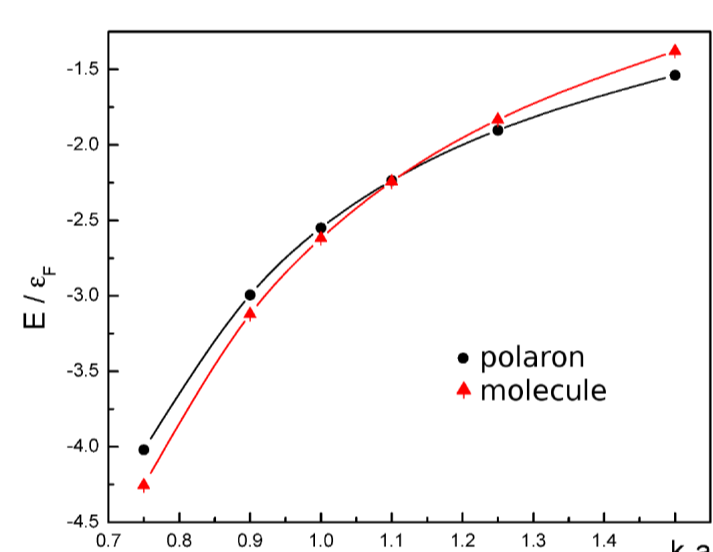


strongest attraction:  
molecular bound  
state

image credit: Schirotzek, Wu, Sommer, Zwerlein, PRL102, 230402 (2009)

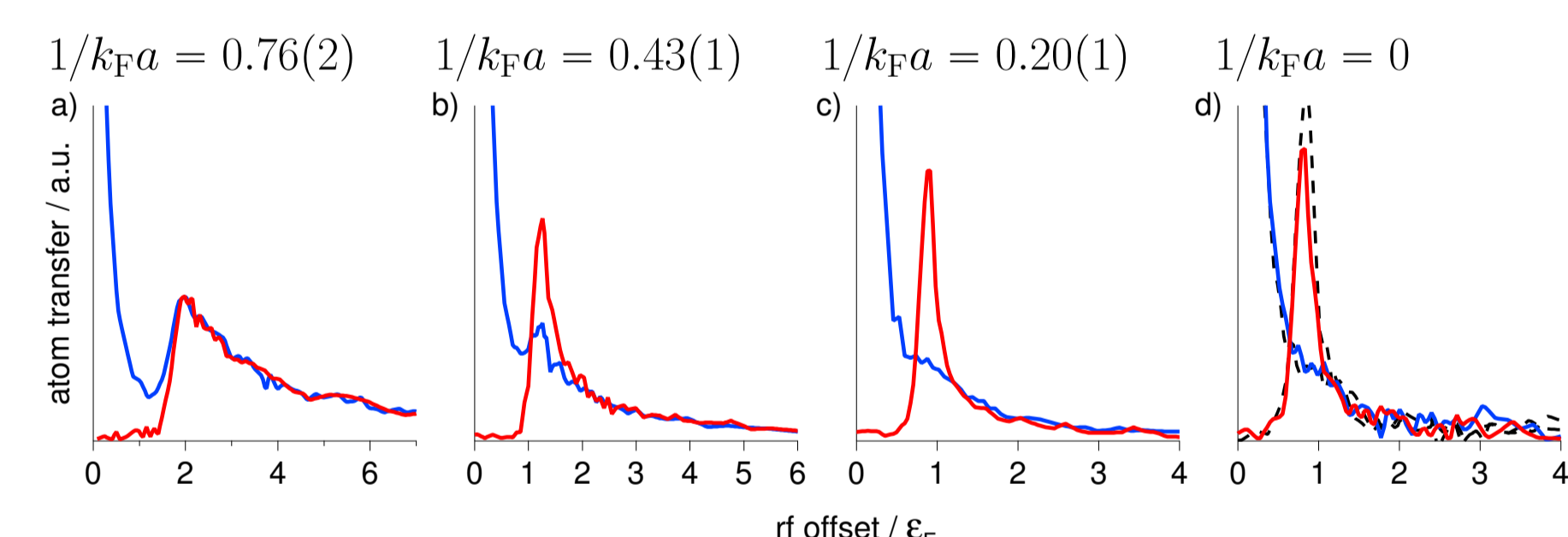
## Interesting questions:

- ground state energy and  $Z$ -factor,
- polaron to molecule transition,
- polaron effective mass,
- unequal masses, trimers,
- spectral function, kinetics,...



[Prokof'ev & Svistunov, PRB77, 125101 and 020408 (2008)]

Experiments measure the **spectral function** via rf-spectroscopy



[Schirotzek, Wu, Sommer and Zwerlein, PRL102, 230402 (2009)]

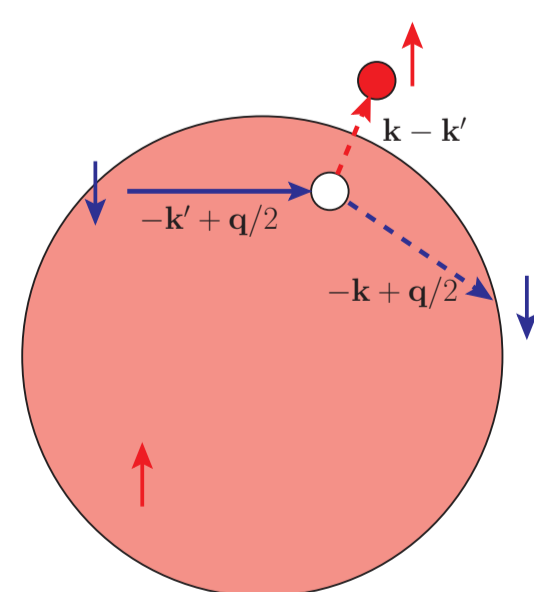
Our goals:

- Quantitative** and **controlled** calculation of the polaron spectral function
- Rigorous interpretation of experimental spectral functions

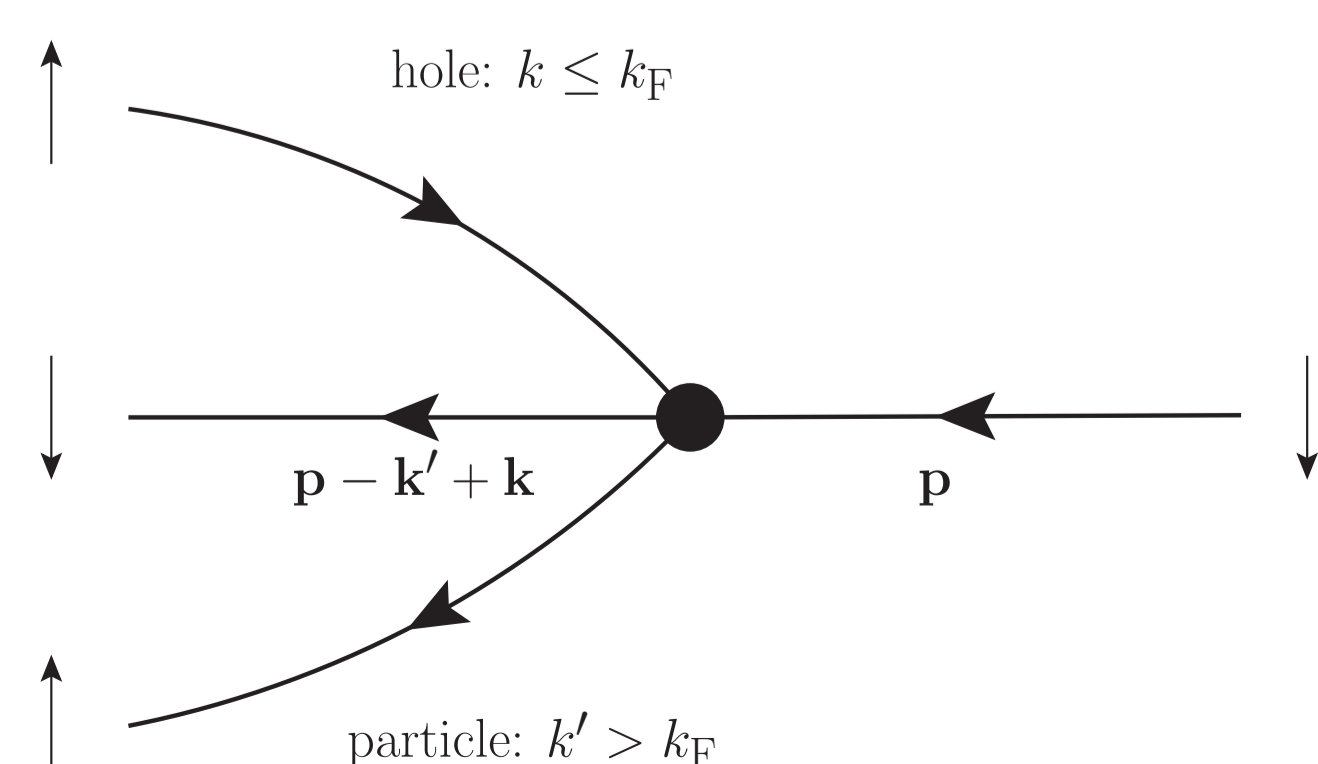
## Model

$$H = \sum_{\mathbf{k}, \sigma=\uparrow, \downarrow} \epsilon_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + g_0 \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} c_{\mathbf{k}+\mathbf{q}/2, \uparrow}^\dagger c_{-\mathbf{k}+\mathbf{q}/2, \downarrow}^\dagger c_{-\mathbf{k}+\mathbf{q}/2, \downarrow} c_{\mathbf{k}+\mathbf{q}/2, \uparrow}$$

- zero temperature
- 3D
- attractive interaction
- $m_\uparrow = m_\downarrow$
- short-ranged potential



Spin  $\downarrow$  impurity excites **particle-hole pairs** via momentum exchange with the spin  $\uparrow$  Fermi sea



## Spectral function

- gives the system's response to the impurity creation operator  $c_{\mathbf{k}\downarrow}^\dagger$
- probability of adding or removing a state with given momentum at energy cost  $\omega$

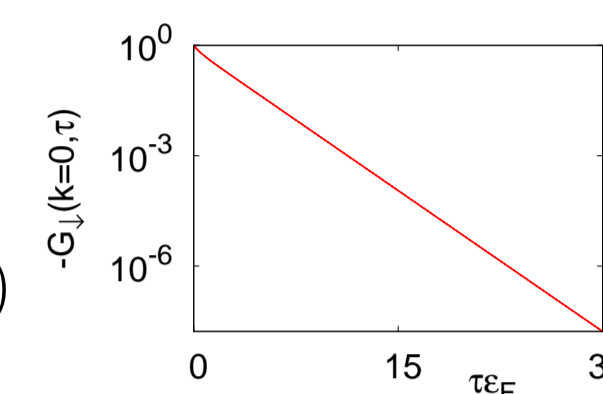
From full impurity Green's function:

$$-G_\downarrow(\mathbf{k}, \tau) = \theta(\tau) \langle c_{\mathbf{k}\downarrow}(\tau) c_{\mathbf{k}\downarrow}^\dagger(0) \rangle = \int_0^\infty A_\downarrow(\mathbf{k}, \omega) e^{-\omega\tau} d\omega$$

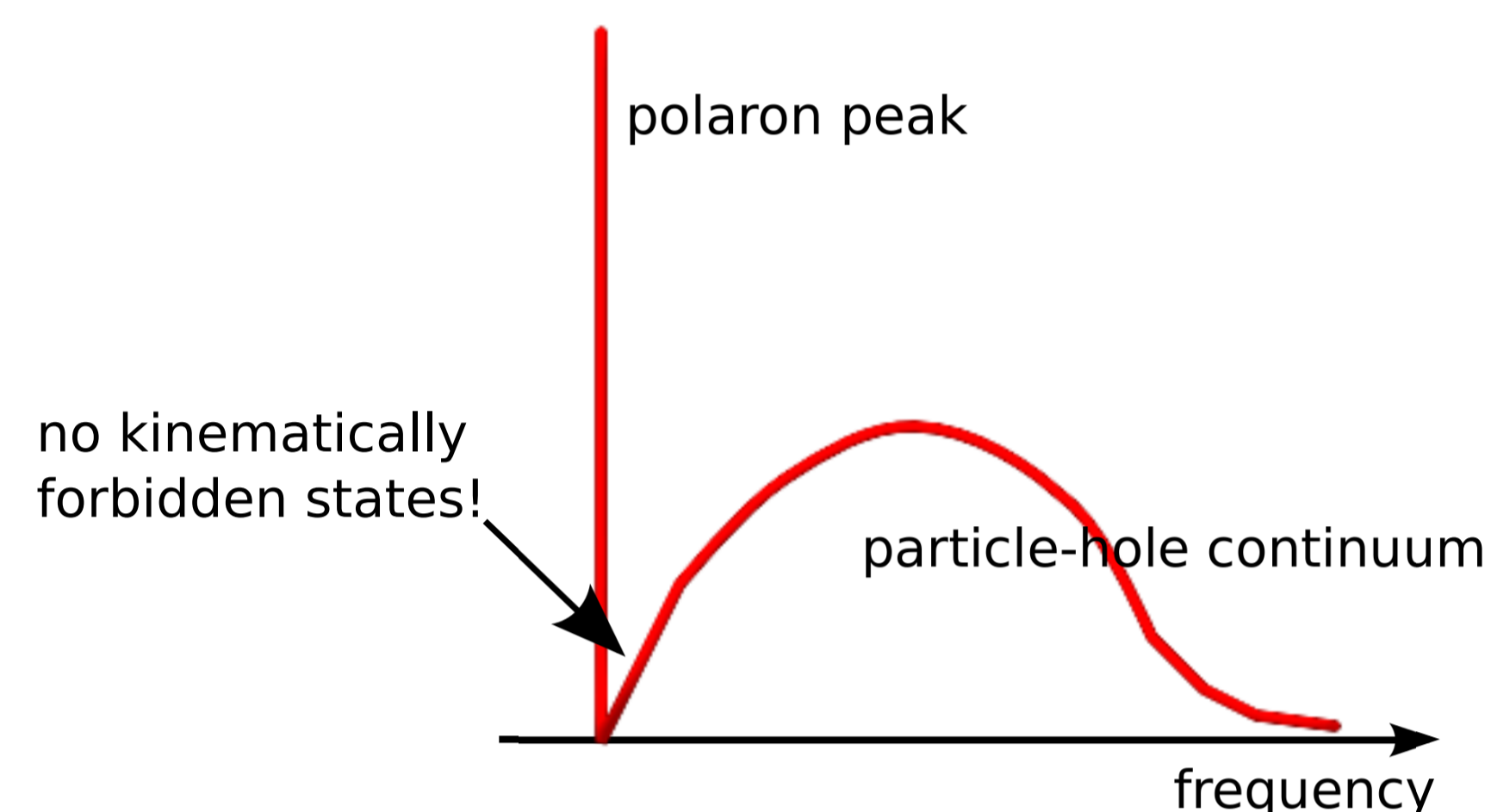
Asymptotic behavior:

$$-G_\downarrow(0, \tau \rightarrow \infty) \rightarrow Z e^{-(E_p - \mu_\downarrow)\tau}$$

$\Rightarrow$  polaron peak in spectrum  $Z\delta(\omega + \mu_\downarrow - E_p)$

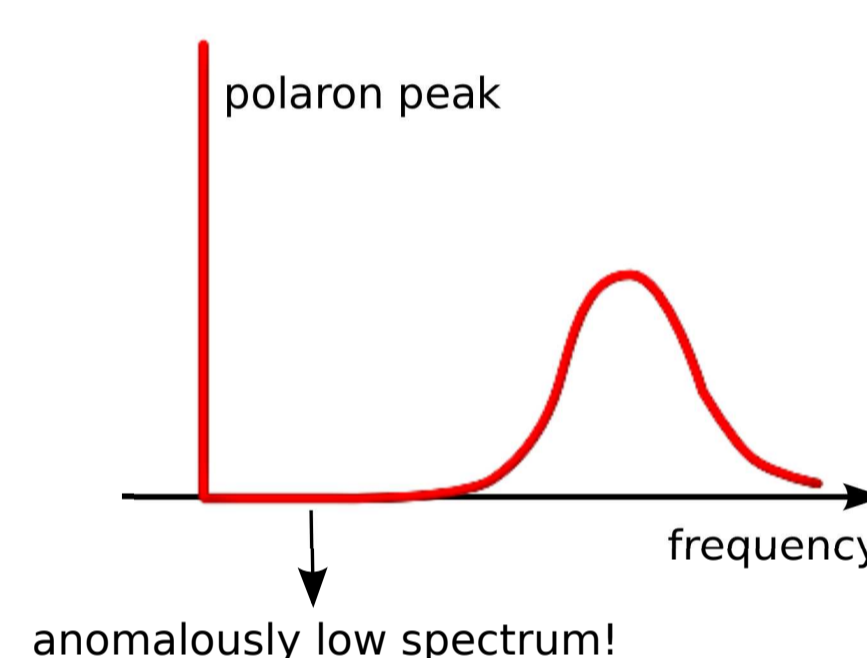


**What we expect:**



(matrix element controls the slope of the spectral function)

**Instead for  $k_F a \lesssim 1$  we find:**



- spectral weight after the polaron peak strongly suppressed
- this happens despite the absence of a small interaction-related parameter!

## Analytic continuation

discretisation:  $A(\omega) = \sum_{n=1}^N A(\omega_n) \delta(\omega - \omega_n) \Rightarrow$  matrix equation

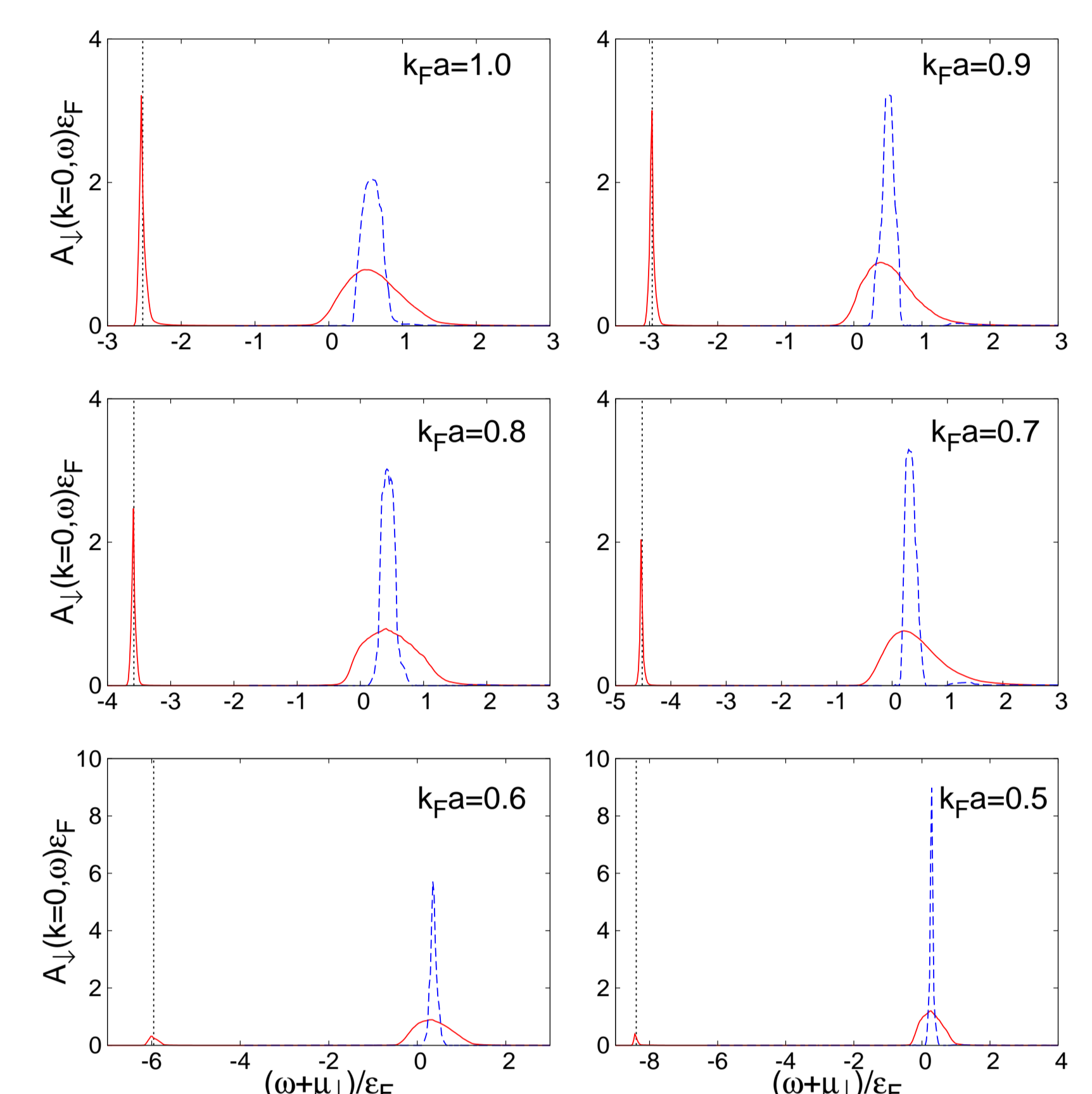
$$-G(\tau) = \int_0^\infty e^{-\omega\tau} A(\omega) d\omega \leftrightarrow -G_m = \sum_{n=1}^N e^{-\omega_n \tau} A(\omega_n)$$

infinite number of solutions satisfying this equation within errors of  $G_m$

### Stochastic Optimization Method

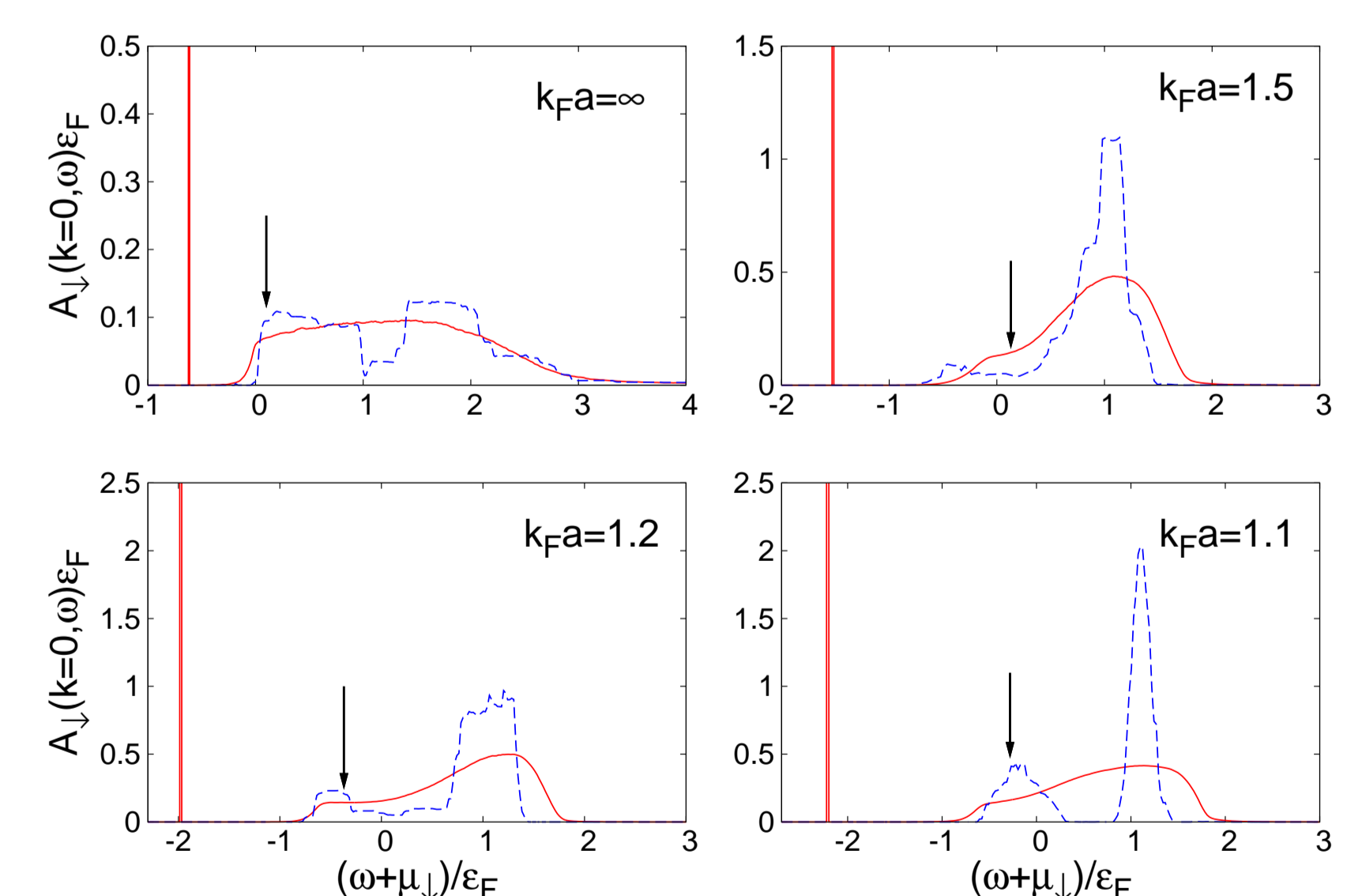
- no default model, no prior except normalization and positivity
- optimize deviation measure:  $D[\tilde{A}] = \sum_{m=1}^M \left| \frac{G_m - \tilde{G}_m}{\tilde{G}_m} \right|$
- linear combinations of good solutions are also good solutions
- "smoothest" linear combination of good solutions? the solution can be **ambiguous**  
 $\Rightarrow$  use target function to force the solution through a set of points  
 $\Rightarrow$  check if any given feature can be distorted in this way

### Polaron spectral function in the molecule regime ( $k = 0$ )

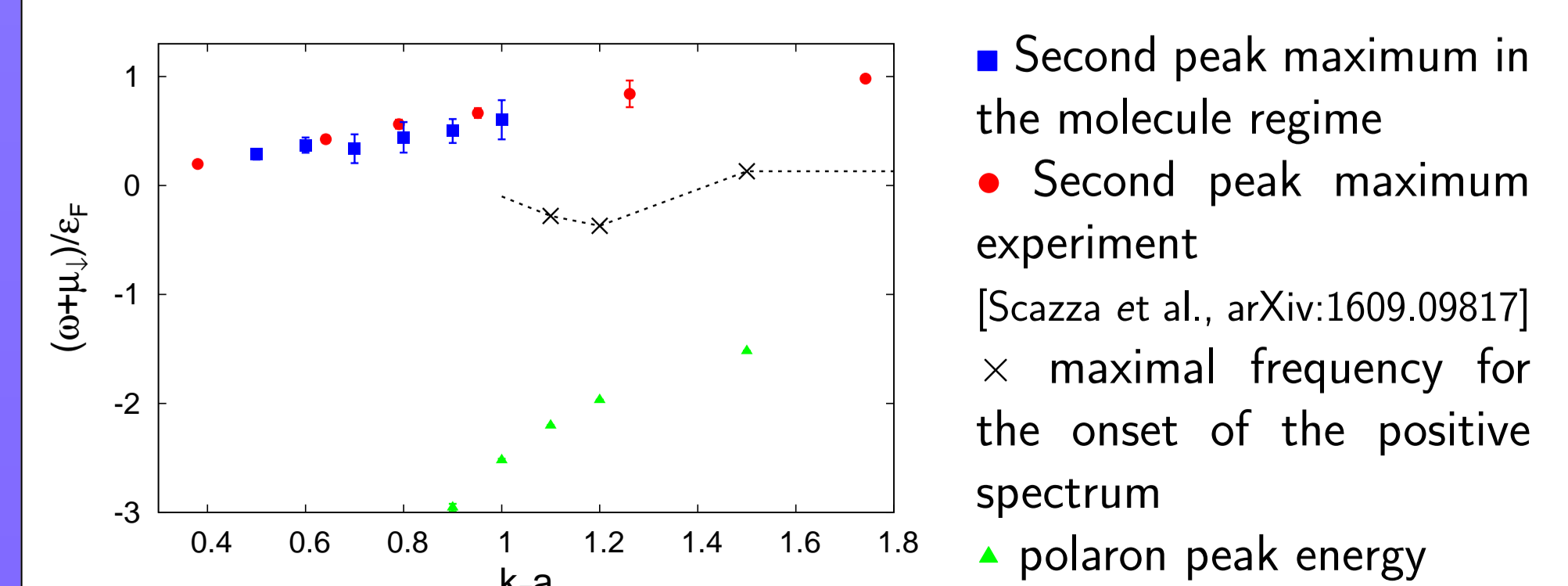


- The maximally smooth spectrum (red solid lines) is broad, but there are solutions consistent with a narrow peak (blue dashed lines)  
 $\Rightarrow$  the width of the second peak cannot be unambiguously established!
- The polaron peak is consistent with a  $\delta$ -function (black short-dashed lines) for all values of  $k_F a$  studied  
 $\Rightarrow$  polaron is a well-defined quasiparticle even on the molecule side
- with decreasing  $k_F a$  the spectral weight of the polaron peak ( $Z$ -factor) diminishes and the peak moves to more negative frequencies  
 $\Rightarrow$  gap-like region of very low spectral weight between peaks  
 $\Rightarrow$  kinematically, these states should get excited, but remain strongly suppressed  
 $\Rightarrow$  without a small parameter at  $k_F a \sim 1$ , the weight of the few-body continuum should be comparable to the  $Z$ -factor

### Polaron spectral function in the polaron regime ( $k = 0$ )

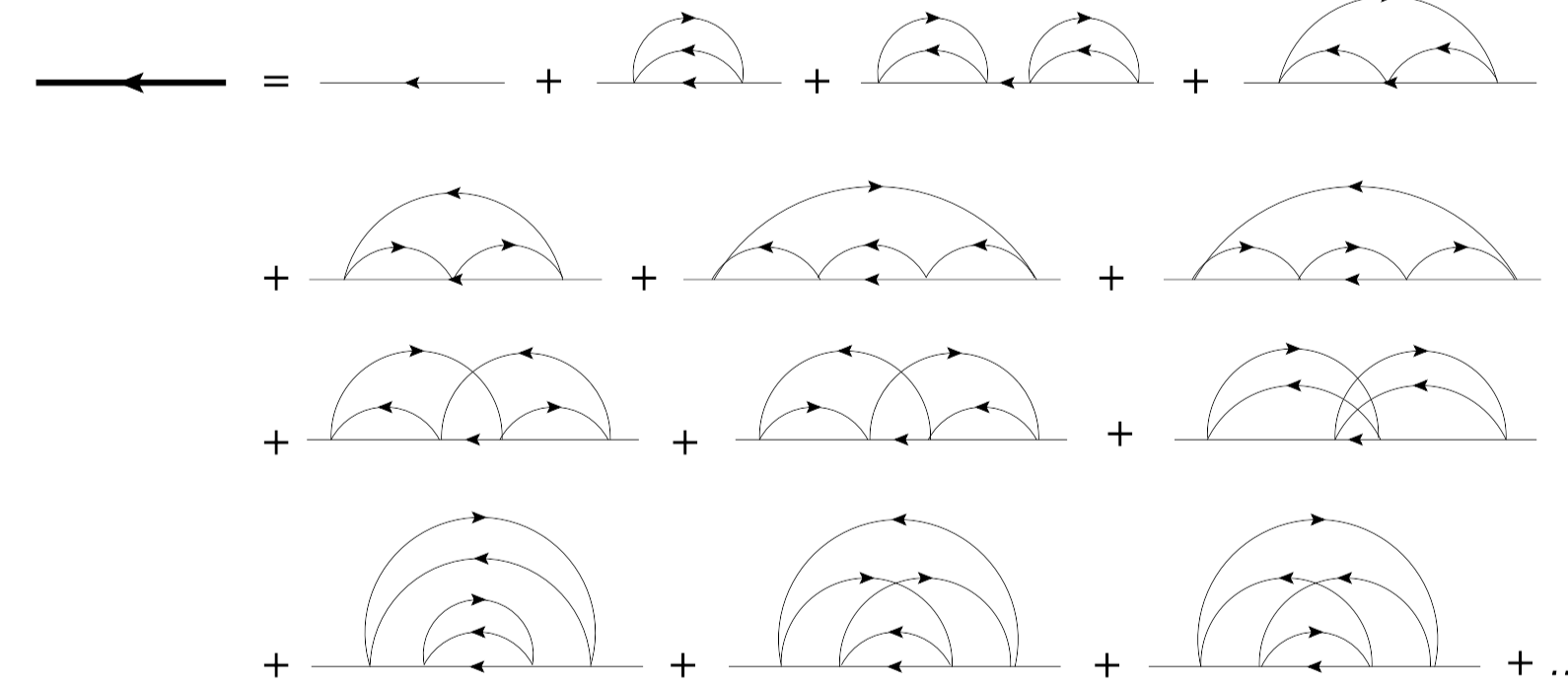


- The maximally smooth spectrum (red solid lines) is a broad peak with preceding plateau, but there are solutions with two separate peaks (blue dashed lines)!
- Either the shoulder or the small intermediate peak must be present
- The highest frequency after which the spectrum must be larger than zero (arrow) is much lower than the position of the major second peak, in contrast to the molecule side.



## Diagrammatic Monte Carlo

Expansion in the free Green's function and coupling constant:



- excitation of multiple particle-hole pairs
- impurity exchanges momentum with particles and holes

**Factorially growing number of diagrams!**

- Partial summation: replace interaction vertex with **pair propagator**  $\Gamma_0$  (sum of ladder diagrams)

- Sample **self-energy**: no 1-particle reducible diagrams

Dyson equation:

$$\frac{G}{\Gamma} = \frac{G^0}{\Gamma^0} + \frac{\Sigma}{\Gamma}$$

- Boldification**: expand in terms of the full Green's function (and full  $\Gamma$ -function) [Prokof'ev and Svistunov, PRB77, 125101 and 020408 (2008)]

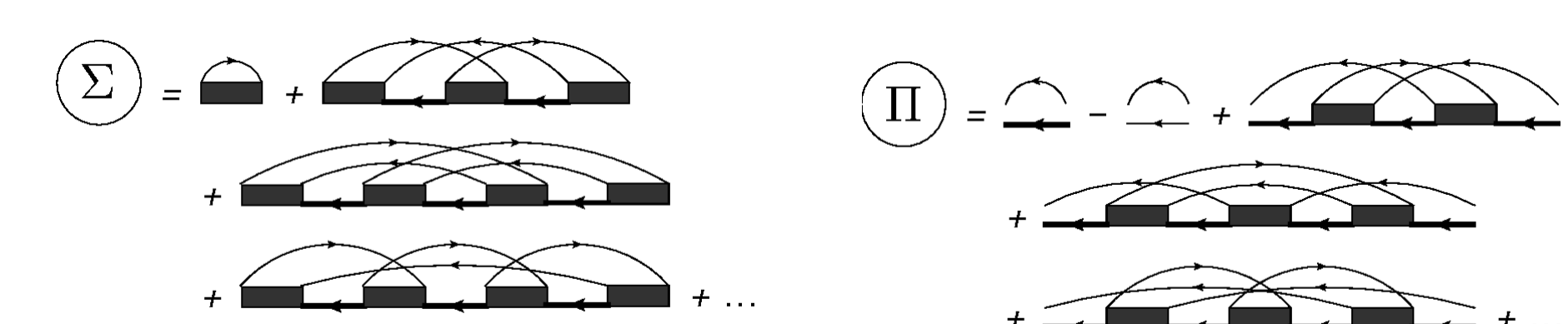


image credit: Vlietinck, Ryckebusch, Van Houcke, PRB87, 115133 (2013)

**Worm trick**: introduce artificial elements to simplify updates



- weight of configurations with worms can be adjusted to make updates maximally efficient
- free worm ends are a convenient intermediate stage between different complex configurations  
 $\Rightarrow$  chains of several small (local) updates more efficient than attempting large (global) updates