

Standard Model β spectra

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ACFI Workshop, November 1st 2018

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Where to look for it?

Quirky weak interaction!

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Where to look for it?

Quirky weak interaction!

How to do this?

β Spectrum shape

1. Direct BSM sensitivity
2. Enters into reactor anomaly

General Hamiltonian

$$\mathcal{H} = \sum_{j=V,A,S,P,T} \langle f | \mathcal{O}_j | i \rangle \langle e | \mathcal{O}_j [C_j + C_j' \gamma_5] | \nu \rangle + h.c.$$

Introduction

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QCD influences \rightarrow *induced* currents, influenced through **nuclear structure**?

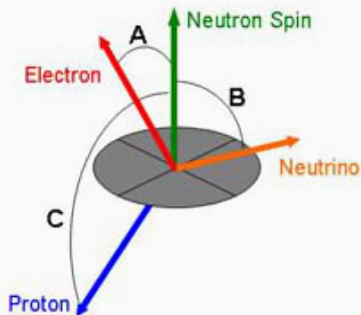
BSM Observables in β decay

Typical BSM searches through correlations

$$\frac{d\Gamma}{dE_e d\Omega_e d\Omega_\nu} \propto 1 + a_{\beta\nu} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b_F \frac{m_e}{E_e} + A \frac{\vec{p}_e}{E_e} \langle \vec{l} \rangle + \dots$$

Measure effective correlations

$$\tilde{X} = \frac{X}{1 + b_F \langle \frac{m_e}{E_e} \rangle}$$



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Sensitivity comes from b_F

$$b_F = \pm \frac{1}{1 + \rho^2} \left[\text{Re} \left(\frac{C_S + C'_S}{C_V} \right) + \rho^2 \text{Re} \left(\frac{C_T + C'_T}{C_A} \right) \right]$$

because it's **linear** in coupling constants

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because it's **linear** in coupling constants

→ measure β spectrum directly & fit for $1/E_e$

Beta Spectrum Shape

Exploring the Standard Model and Beyond via the allowed β spectrum shape:

$$\frac{dN}{dE_e} \propto 1 + b_{\text{Fierz}} \frac{m_e}{E_e} + b_{\text{WM}} E_e$$

b_{Fierz} : Proportional to scalar (Fermi) and tensor (Gamow-Teller) couplings

b_{WM} : Weak Magnetism (main induced current), poorly known for $A > 60$, forbidden decays

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This requires knowledge of the theoretical spectrum shape to $\leq 10^{-3}$ level!

Beta spectrum shape

Beta Spectrum Shape

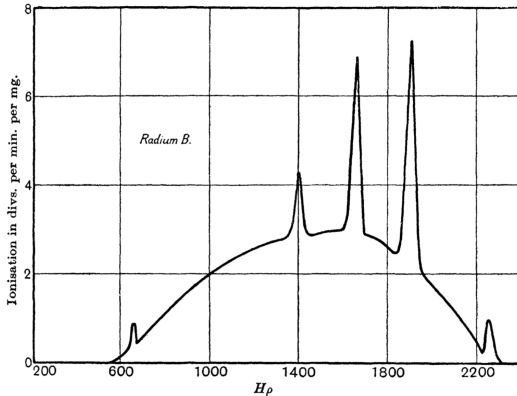
3-body decay

$$P(E_e) = (E_0 - E_e)^2 E_e p_e \approx (E_0 - E_e)^2 E_e^2$$

Beta Spectrum Shape

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Ellis & Chadwick, 1914

Beta Spectrum Shape

Active participation of QED, QCD & WI \rightarrow Complicated system

Weak Hamiltonian is **modified**

1. Emitted β particle immersed in Coulomb field: (electroweak) radiative corrections

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Large scale gap to cross

Quark \rightarrow Nucleon \rightarrow Nucleus \rightarrow Atom \rightarrow Molecule

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Whole slew of **approximations** introduced

Standard Model Calculation: Quark

Starting from the Standard Model $SU(2)_L \times U(1)_Y$ EW sector

$$\mathcal{M} = \frac{g^2}{8} V_{ud} \bar{u} \gamma^\mu (1 - \gamma^5) d \frac{g_{\mu\nu} - q_\mu q_\nu / M_W^2}{q^2 - M_W^2} \bar{e} \gamma^\nu (1 - \gamma^5) \nu$$

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Since $q \ll M_W$, identify Fermi coupling constant

$$G_F = \frac{g^2}{8M_W^2}$$

Moving to the nucleon system, we face

$$\langle p | \bar{u} \gamma^\mu (1 - \gamma^5) d | n \rangle$$

Standard Model Calculation: Nucleon

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$$\langle p | \bar{u} \gamma^\mu (1 - \gamma^5) d | n \rangle$$

Symmetries to the rescue! CVC & PCAC define new *nucleon* currents

$$V^\mu + A^\mu \approx g_V(q^2) \gamma^\mu (1 - \lambda \gamma^5)$$

where $g_V(q^2) \approx 1$ and λ from the lattice

Standard Model Calculation: Nucleon

Strong interaction introduces extra terms into the vertex \rightarrow
Construct all Lorentz invariants

$$\langle p|V^\mu|n\rangle = \bar{p} \left[g_V \gamma^\mu + \frac{g_M - g_V}{2M} \sigma^{\mu\nu} q_\nu + i \frac{g_S}{2M} q^\mu \right] n$$
$$\langle p|A^\mu|n\rangle = \bar{p} \left[g_A \gamma^\mu \gamma^5 + \frac{g_T}{2M} \sigma^{\mu\nu} q_\nu \gamma^5 + i \frac{g_P}{2M} q^\mu \gamma^5 \right] n$$

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Introduction of recoil ($\sim q/M$) terms

Standard Model Calculation: Nucleon

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Introduction of recoil ($\sim q/M$) terms

CVC requires $g_S = 0$ & $g_M = \mu_p^{an} - \mu_n = 4.7$

Standard Model Calculation: Nucleus

Nucleus is spherical system \rightarrow multipole decomposition,
elementary particle

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Relativistic generalization in **Breit frame**

$$\langle f | V^0 + A^0 | i \rangle \propto \sum_{LM} (-)^{J_f - M_f} \begin{pmatrix} J_f & L & J_i \\ -M_f & M & M_i \end{pmatrix} (Y_L^M)^* F_L(q^2)$$

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Form factors \sim reduced matrix elements

Standard Model Calculation: Nucleus

Require transformation from form factors to matrix elements

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Immediately faced with several issues:

- Weak current in strongly bound system?

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Here the going gets rough → **severe approximations**

Standard Model Calculation: Nucleus

Weak current in strongly bound system?

→ **Impulse approximation**, non-interacting nucleons

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- Neglects meson exchange
- Nucleon-nucleon interaction present in many-body methods

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Relativistic nuclear wave functions

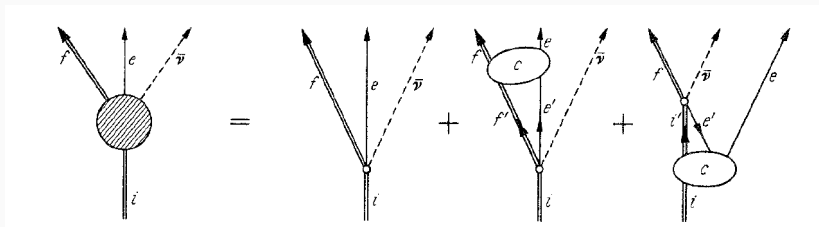
→ **Non-relativistic** nucleons

- expand operator to $\mathcal{O}(v/c)$
- Incomplete wave function basis, core polarization

Standard Model Calculation: Nucleus

Final state interactions

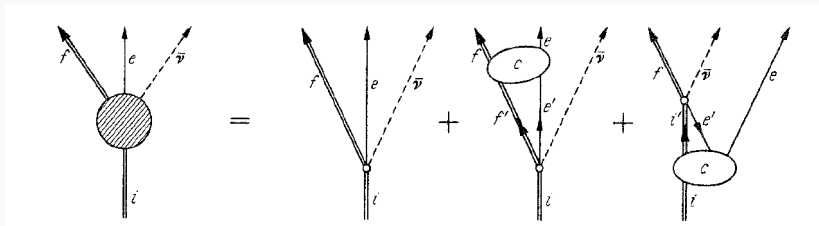
1. Coulomb interaction



Standard Model Calculation: Nucleus

Final state interactions

1. Coulomb interaction



→ Fermi function, *induced Coulomb* terms

Standard Model Calculation: Nucleus

Final state interactions

1. Coulomb interaction

Make several **approximations**

- Initial & Final Coulomb potentials are same

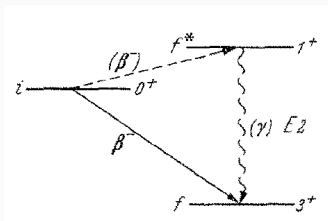
Standard Model Calculation: Nucleus

Final state interactions

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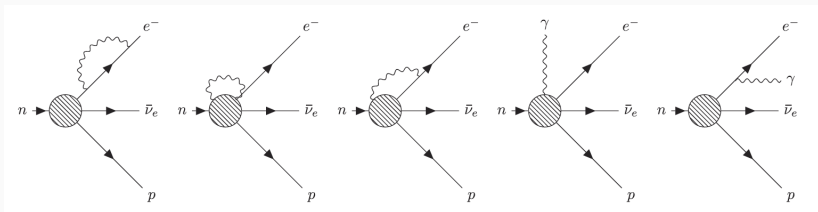
- Initial & Final Coulomb potentials are same
- Typically neglect intermediate decays



Standard Model Calculation: Nucleus

Final state interactions

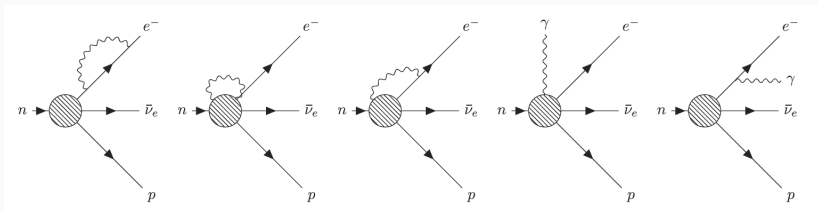
2. EW Radiative corrections



Standard Model Calculation: Nucleus

Final state interactions

2. EW Radiative corrections



+ higher orders, γW boxes: see previous talks

Standard Model Calculation: Atom

Must consider total nuclear + atomic Hamiltonian

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Changes

- Available phase space
- Final state interactions
- Opens new decay modes (bound & exchange)

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Changes

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- Final state interactions
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Require *atomic* wave functions

- Central & static potential
- Sudden approximation

Standard Model Calculation: Molecule

Similar as atomic system, but changes

- Available phase space
- Molecular excitation, ionization
- Recoil correction & distribution

Standard Model Calculation: Molecule

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- Available phase space
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Enter quantum chemistry

- Born-Oppenheimer approximation
- MOLCAO

Current status

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$$\begin{aligned} N(W)dW = & \frac{G_V^2 V_{ud}^2}{2\pi^3} F_0(Z, W) L_0(Z, W) U(Z, W) R_N(W, W_0, M) \\ & \times Q(Z, W, M) R(W, W_0) S(Z, W) X(Z, W) r(Z, W) \\ & \times C(Z, W) D_C(Z, W, \beta_2) D_{FS}(Z, W, \beta_2) \\ & \times pW(W_0 - W)^2 dW \end{aligned}$$

LH *et al.*, Rev. Mod. Phys. 90 (2018) 015008; 1709.07530



Analytical β spectrum shape

80 years of history, in detail

Item	Effect	Formula	Magnitude
1	Phase space factor	$pW(W_0 - W)^2$	Unity or larger
2	Traditional Fermi function	F_0	
3	Finite size of the nucleus	L_0	
4	Radiative corrections	R	
5	Shape factor	C	10^{-1} - 10^{-2}
6	Atomic exchange	X	
7	Atomic mismatch	r	

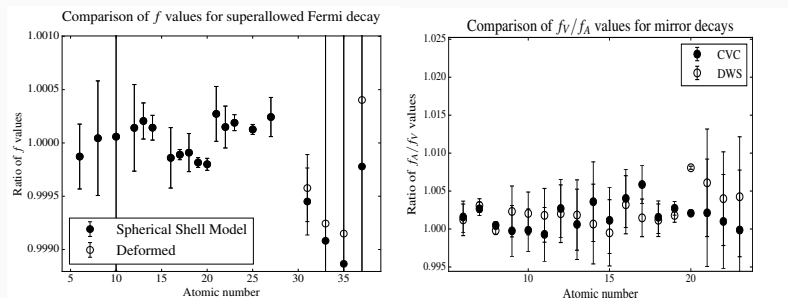
Analytical β spectrum shape

Item	Effect	Formula	Magnitude
8	Atomic screening	S	
9	Shake-up	See 7	
10	Shake-off	See 7	
11	Isovector correction	C_I	
12	Recoil Coulomb correction	Q	$10^{-3}-10^{-4}$
13	Diffuse nuclear surface	U	
14	Nuclear deformation	D_{FS}	
15	Recoiling nucleus	R_N	
16	Molecular screening	ΔS_{Mol}	
17	Molecular exchange	Case by case	

Added/Improved/Didactic

Performance summary

Comparison against numerical results for superallowed & mirror transitions



Agreement is very good

Serves as input for several experiments, C++ code available

Order of magnitude estimates

Nuclear structure sensitivity in shape factor

$$C(Z, W) \sim 1 \pm \frac{4}{3} \frac{W}{M_N} \frac{\mathbf{b}}{A_c} \pm \frac{4\sqrt{2}}{21} \alpha ZWR\Lambda - \frac{1}{3WM_c} (\pm 2\mathbf{b} + \mathbf{d})$$

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Fill in typical numbers to obtain

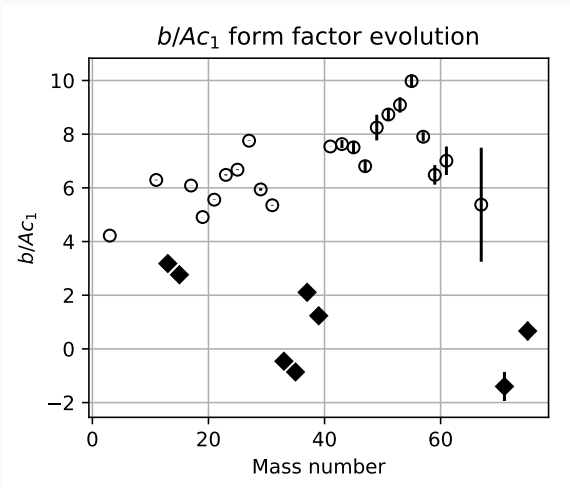
Matrix element	Name	Slope (% MeV ⁻¹)
b	Weak Magnetism	0.5
d	Induced Tensor	0.1
Λ	Induced Pseudoscalar	0.1

Weak magnetism is generally more stable than others

→ **essential** to get this right

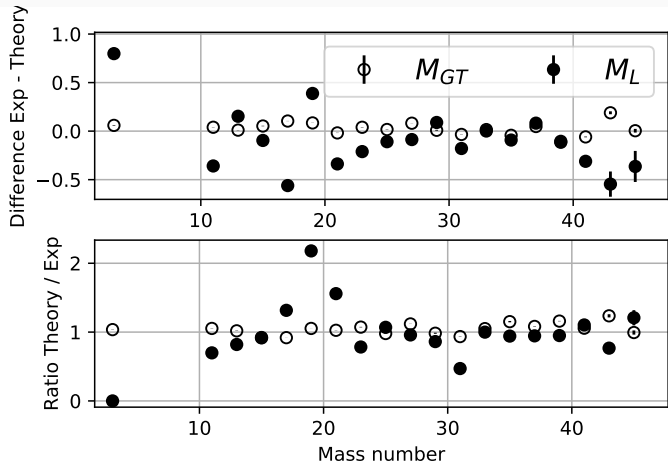
Weak magnetism

Mirror nuclei have CVC-determined WM



open: $I + 1/2$, closed: $I - 1/2$

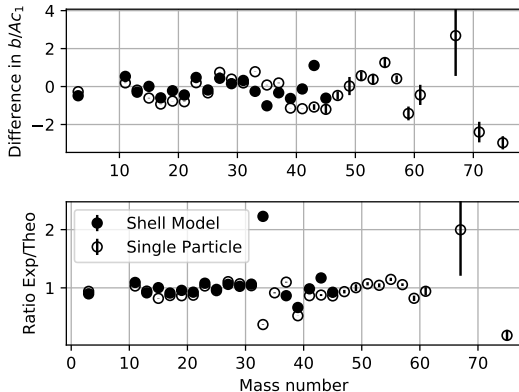
Weak magnetism



'Easy' matrix elements only accurate to 10-20%

Weak magnetism

How does shell model perform right now?



$$\Delta b/Ac = 1 \rightarrow 0.1\% \text{ MeV}^{-1}$$

Still large discrepancies for d/Ac

PHYSICAL REVIEW C **95**, 035501 (2017)

2_1^+ to 3_1^+ γ width in ^{22}Na and second class currents

S. Triambak,^{1,2,*} L. Phuthu,¹ A. García,³ G. C. Harper,³ J. N. Orce,¹ D. A. Short,³ S. P. R. Steininger,³ A. Diaz Varela,⁴
R. Dunlop,⁴ D. S. Jamieson,⁴ W. A. Richter,¹ G. C. Ball,⁵ P. E. Garrett,⁴ C. E. Svensson,⁴ and C. Wrede^{3,6}

$$21(6) \geq d/Ac \geq 3(6)$$

Factor 7 differences depending on shell model results \rightarrow killer!

Challenges

At $\mathcal{O}(10^{-3})$, nuclear structure is main culprit

- Nuclear matrix elements only precise to 10-20%
- Generally: large meson exchange corrections on induced currents
- Isospin multiplet decays are way to go: WM from CVC, induced tensor = 0

At $\leq \mathcal{O}(10^{-4})$, *everything* breaks

At $\leq \mathcal{O}(10^{-4})$, *everything* breaks , but not in the same place!

- Low energy: Atomic & Molecular effects (exchange)
- Endpoint: Final state interactions, excitations
- Radiative corrections: higher order, model dependence
- Low Z : recoil corrections to matrix elements
- High Z : everything electromagnetic

Spectrum shape measurements are valuable tests for S, T currents

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Theoretical spectrum is *theoretically* valid to few 10^{-4}

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Further, radiative & recoil corrections become bottleneck even for nuclear-structure-favorable transitions