

# Sterile Neutrino Portals

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ACFI Neutrino Workshop

July 19, 2017

# Sterile Neutrino Portals [and Dark Matter(s)]

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# Model Building a Dark Sector or: what is a “portal”?

Standard Model  
symmetries

$$SU(3)_c \times SU(2)_L \times U(1)_Y \rightarrow SU(3)_c \times U(1)_{em}$$

Standard Model  
particle content

$$\left. \begin{array}{l} \ell = \left( \begin{array}{l} \nu_L \\ e_L \end{array} \right) \quad e_R \\ q = \left( \begin{array}{l} u_L \\ d_L \end{array} \right) \quad u_R \quad d_R \end{array} \right\} \times 3$$

$$H = \left( \begin{array}{l} \rho^+ \\ v + h + \rho^0 \end{array} \right) \quad G_\mu^a, W_\mu^b, B_\mu \rightarrow G_\mu^a, A_\mu$$

Renormalization: lower dim. operators (fewer fields/particles)  
more important

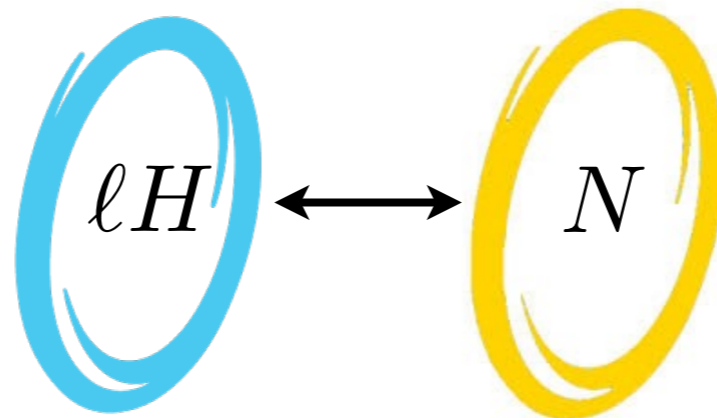
# Model Building a Dark Sector or: what is a “portal”?

Standard Model symmetries  $SU(3)_c \times SU(2)_L \times U(1)_Y \rightarrow SU(3)_c \times U(1)_{em}$

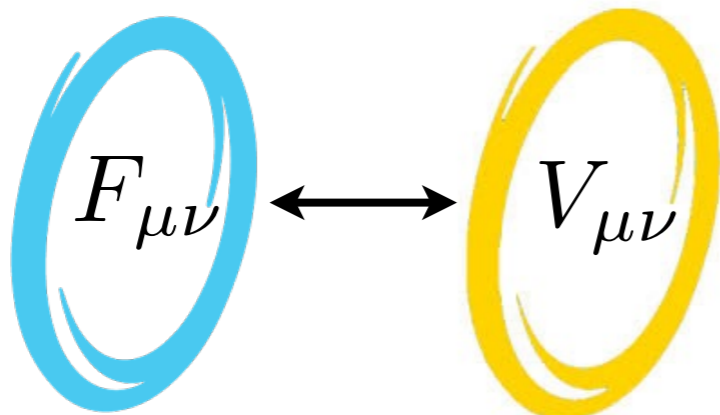
Portals: coupling via stuff uncharged w.r.t. SM

Lead to minimal difficulties incorporating hidden sectors

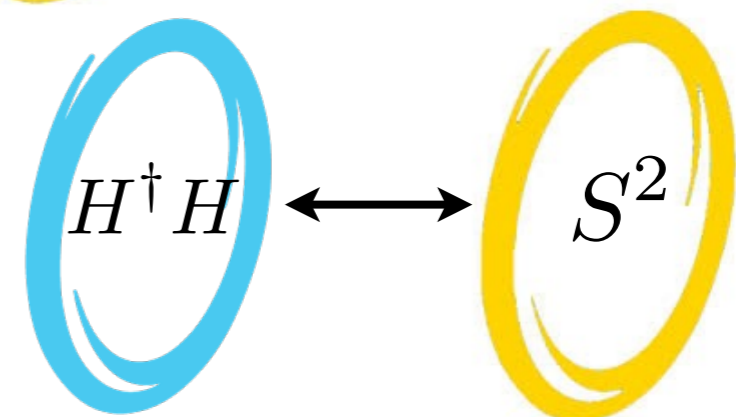
neutrino:



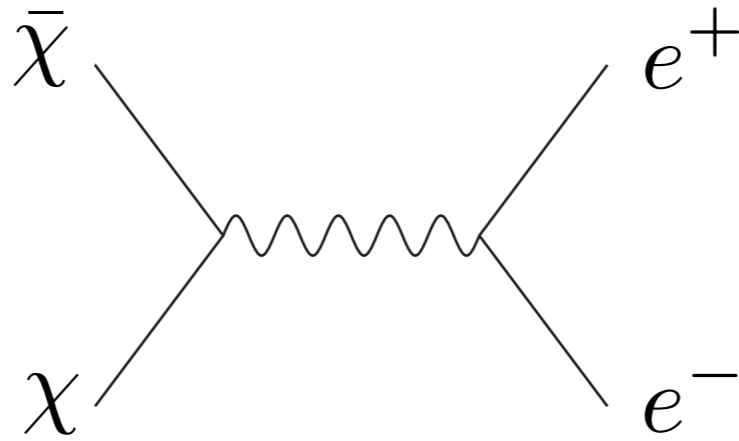
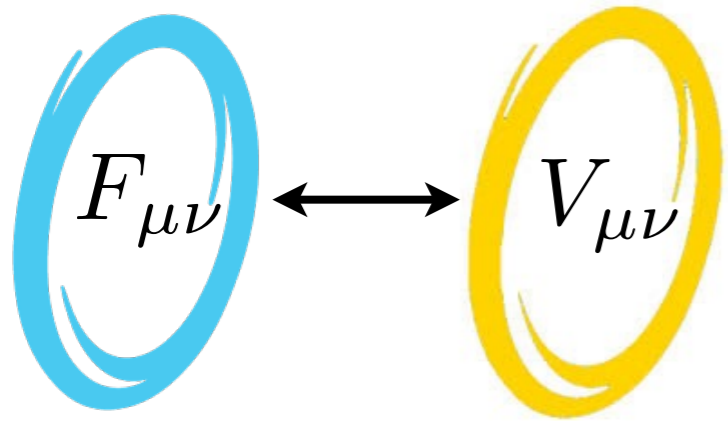
kinetic mixing:



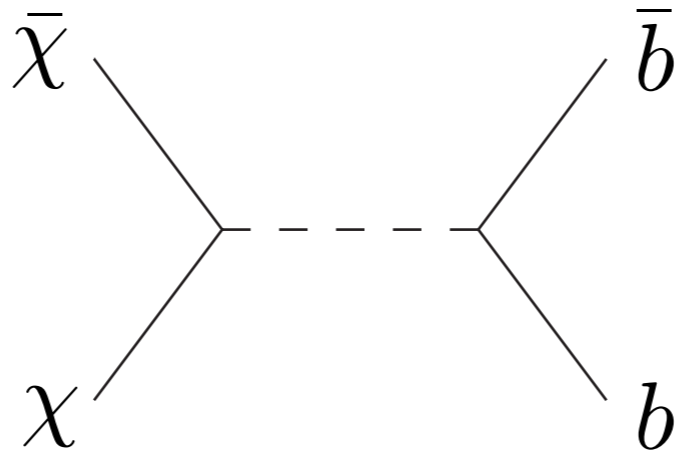
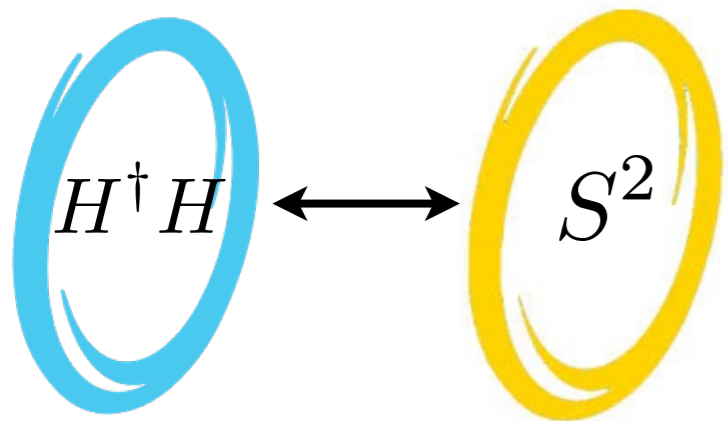
Higgs:



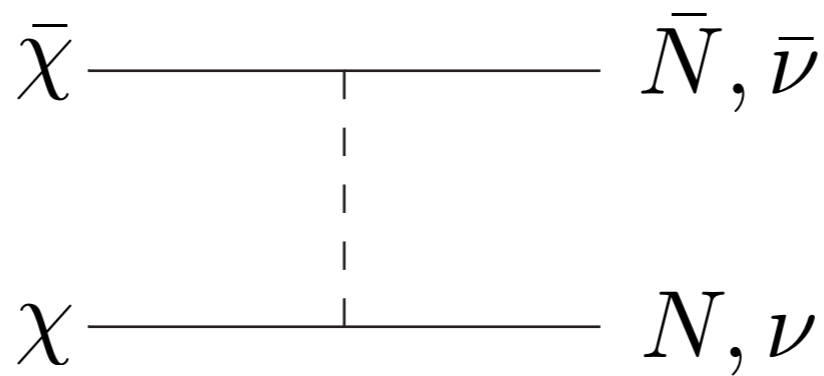
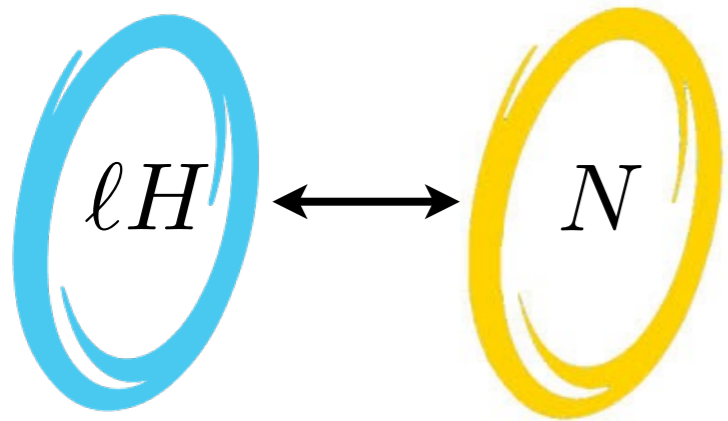




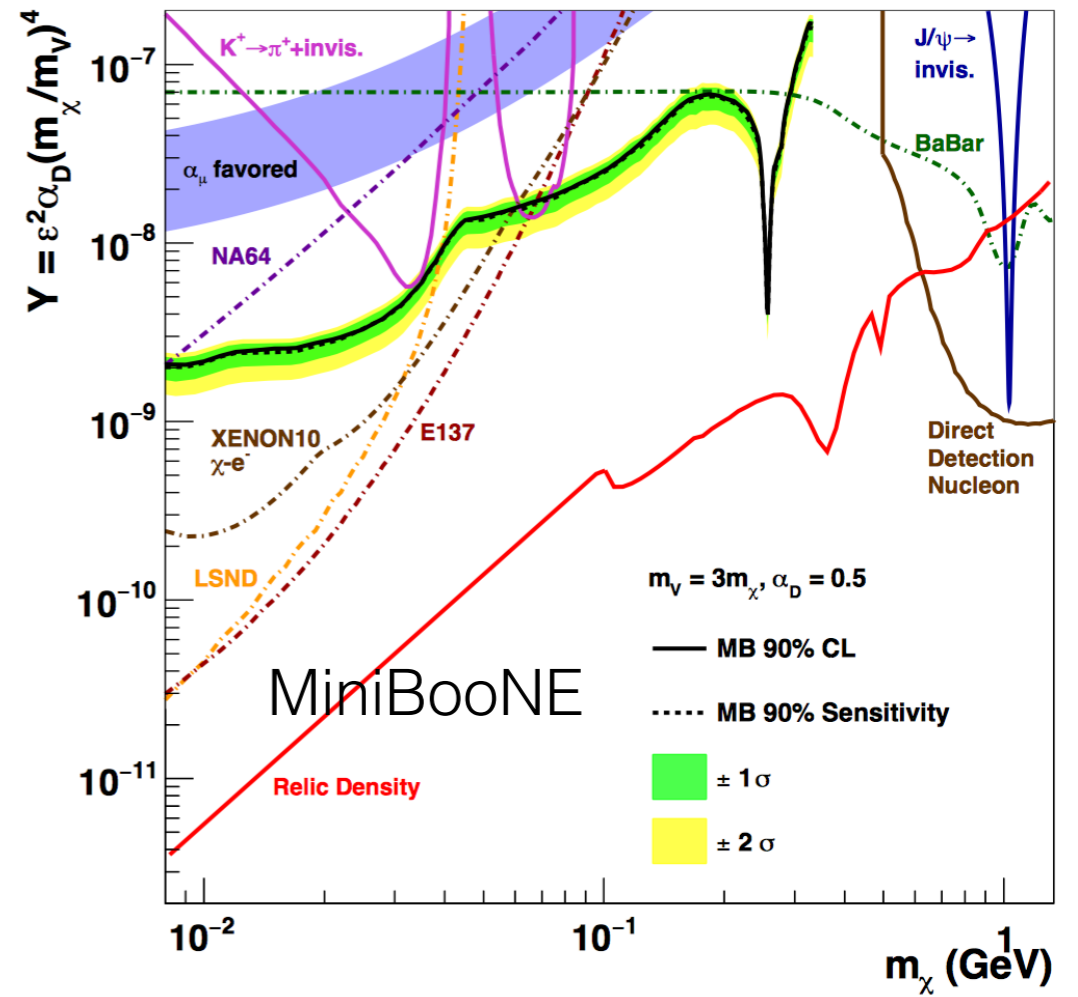
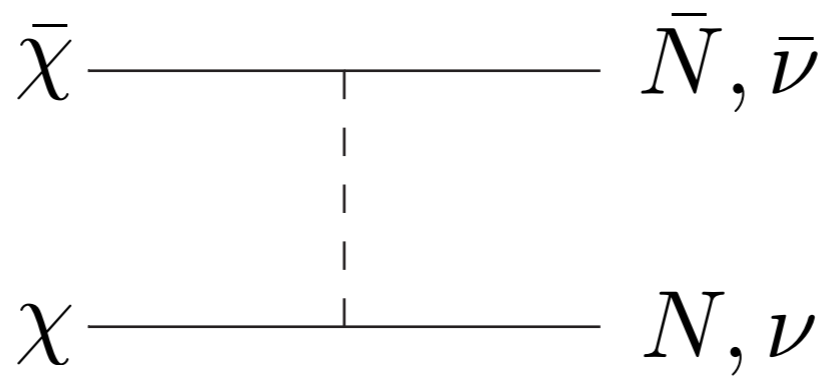
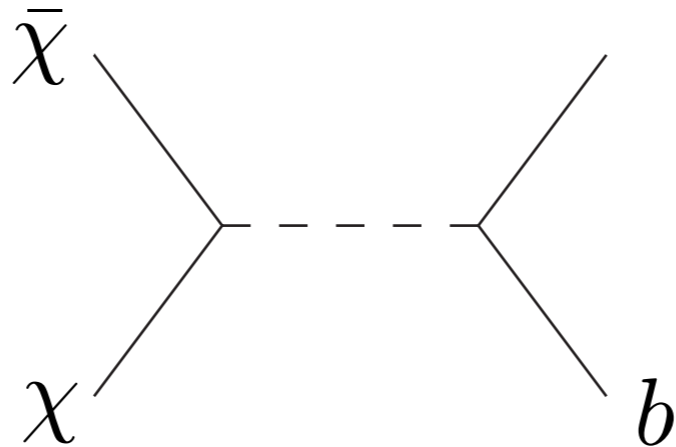
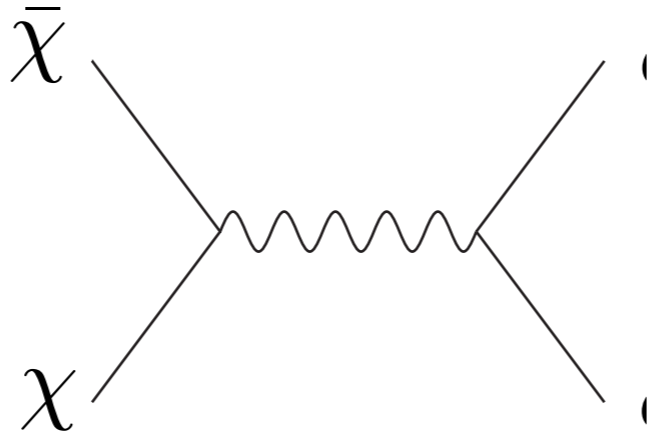
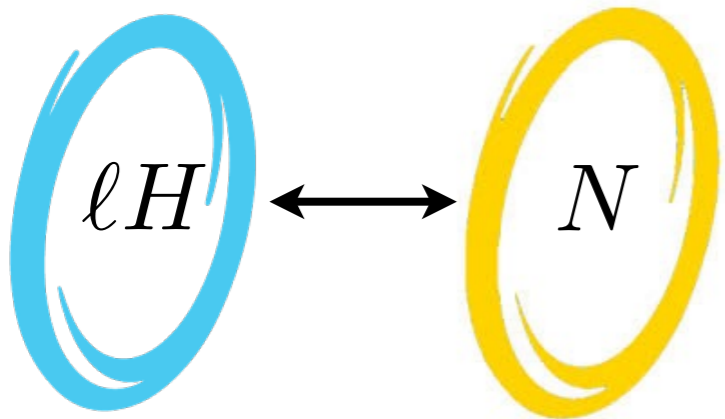
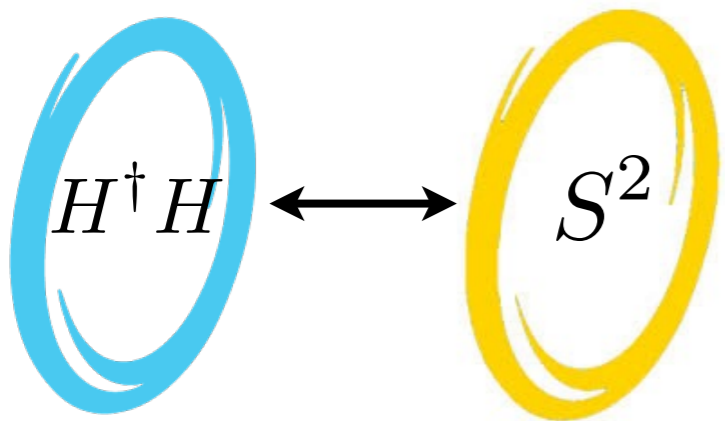
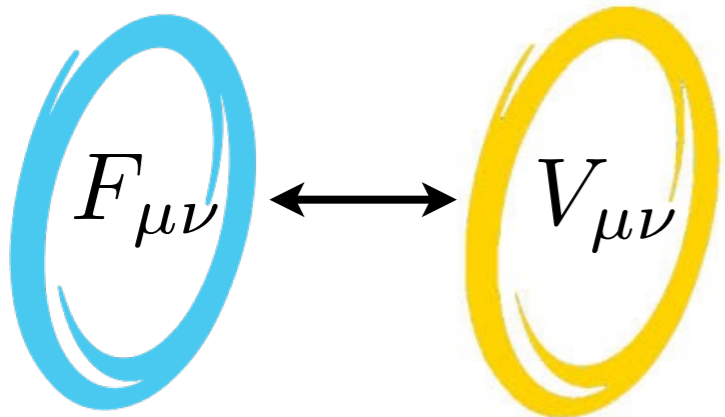
Vector Portal



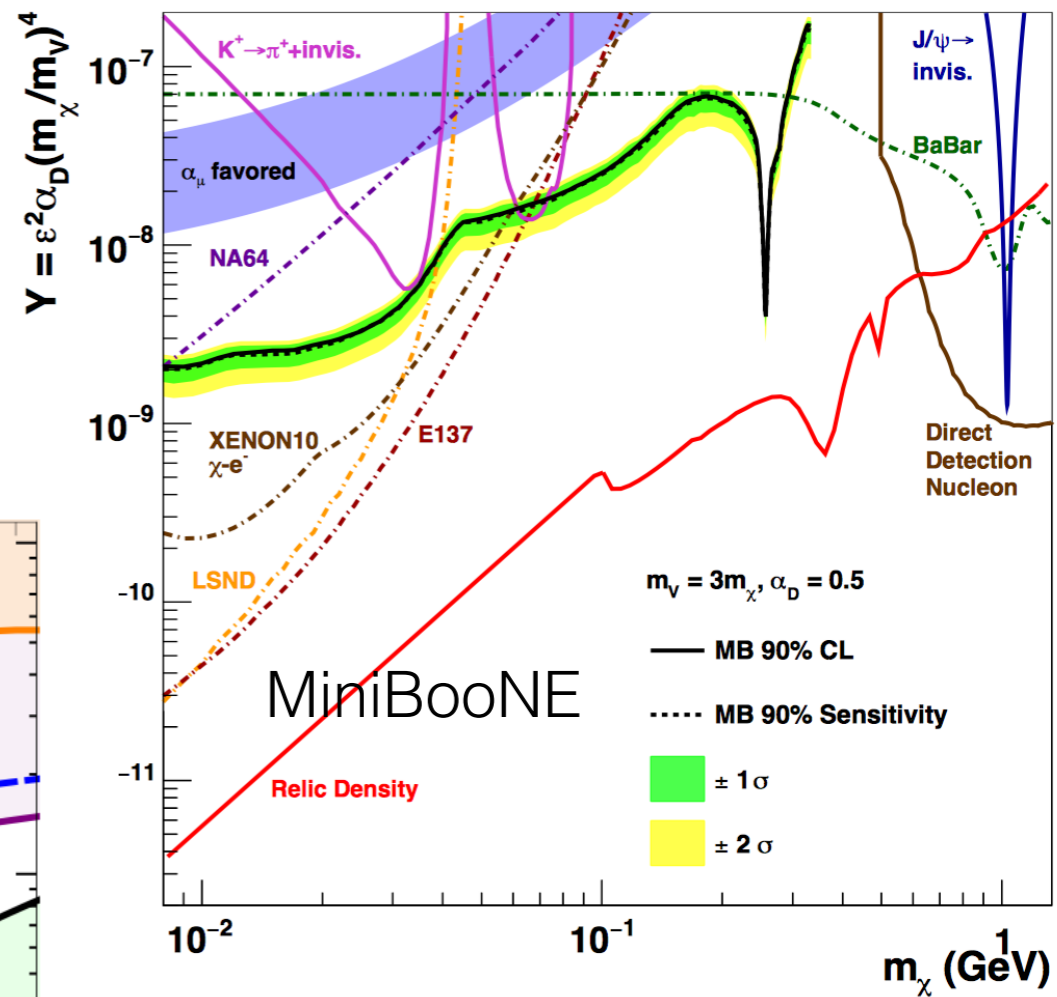
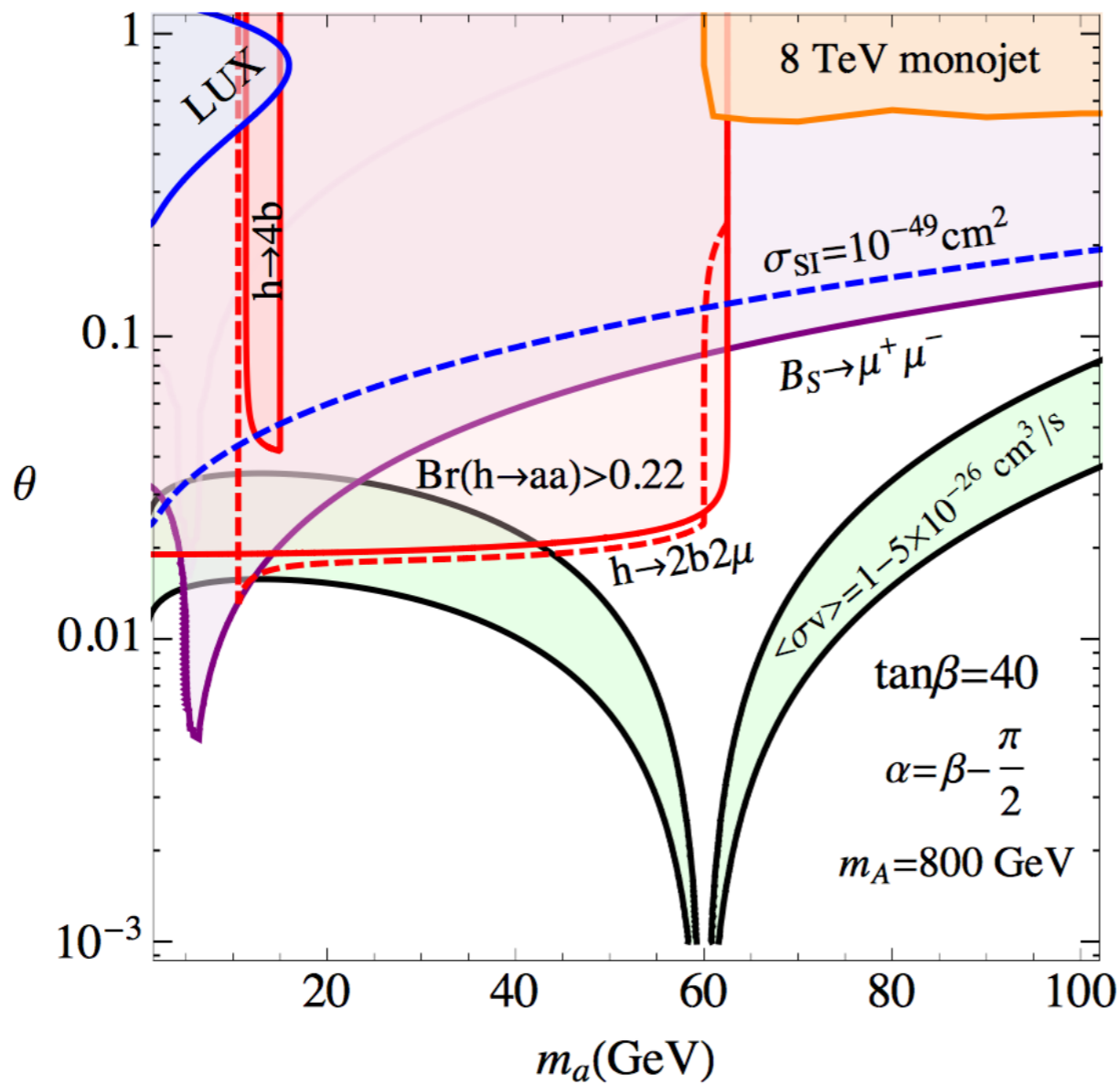
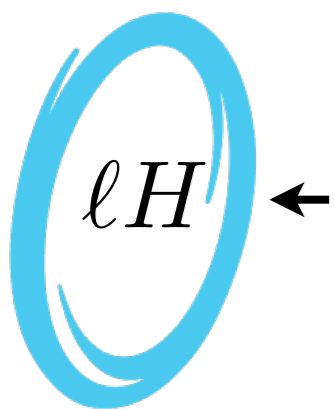
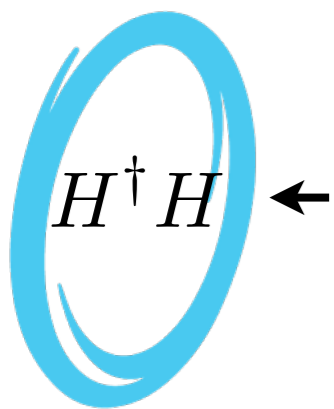
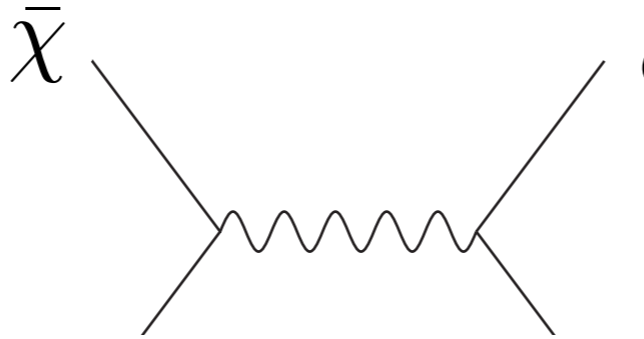
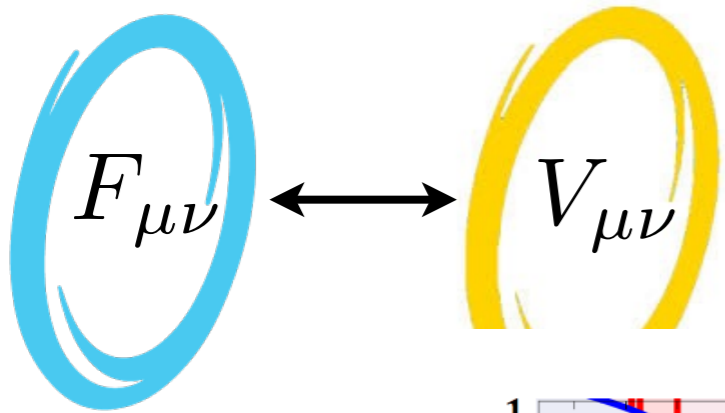
Higgs Portal



Neutrino Portal



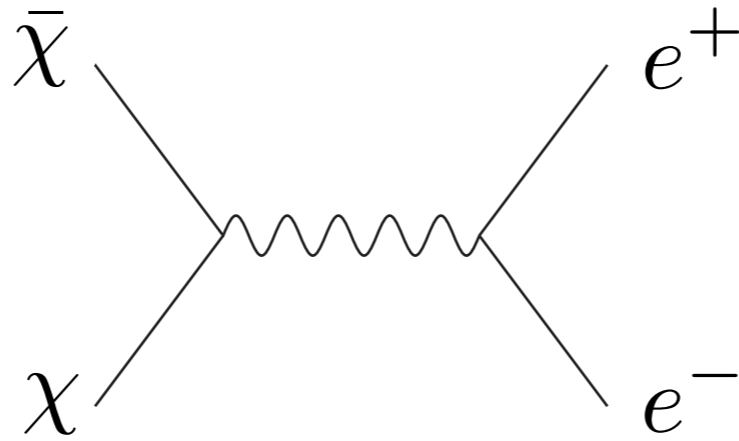
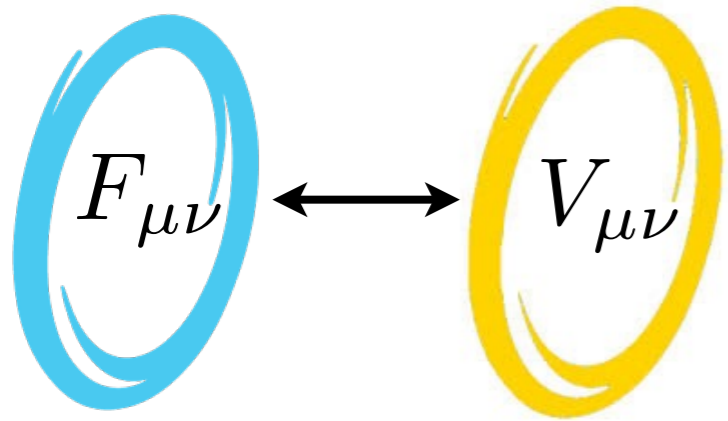
Neutrino Portal



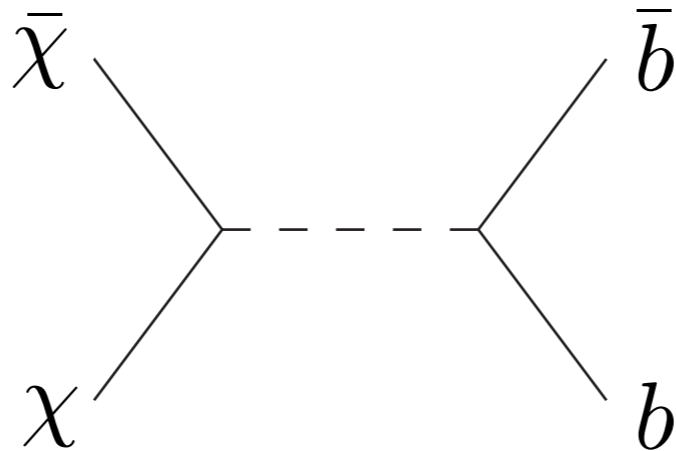
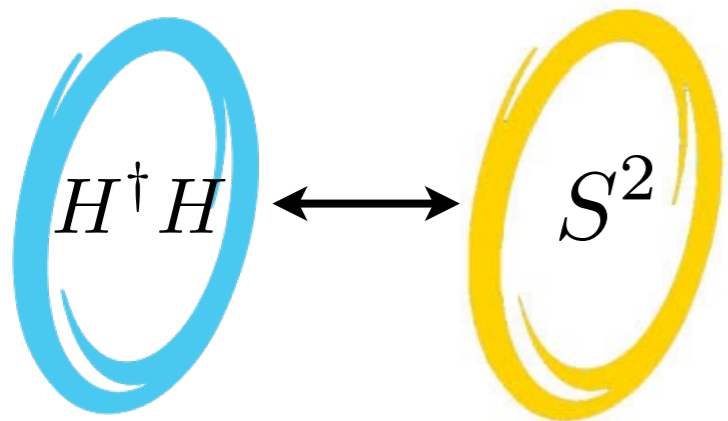
$\bar{\nu}$

Neutrino Portal

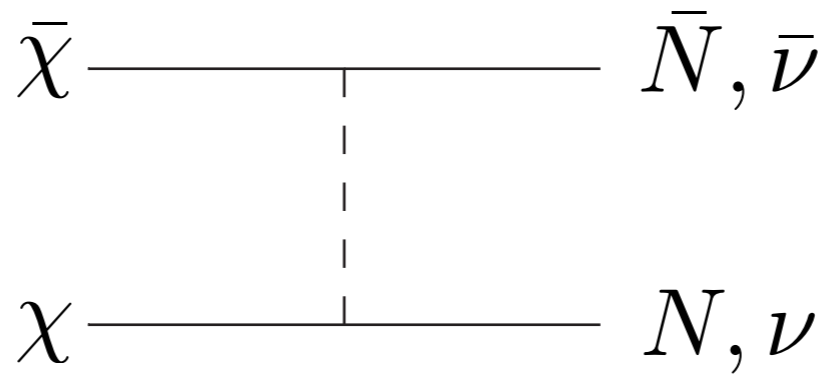
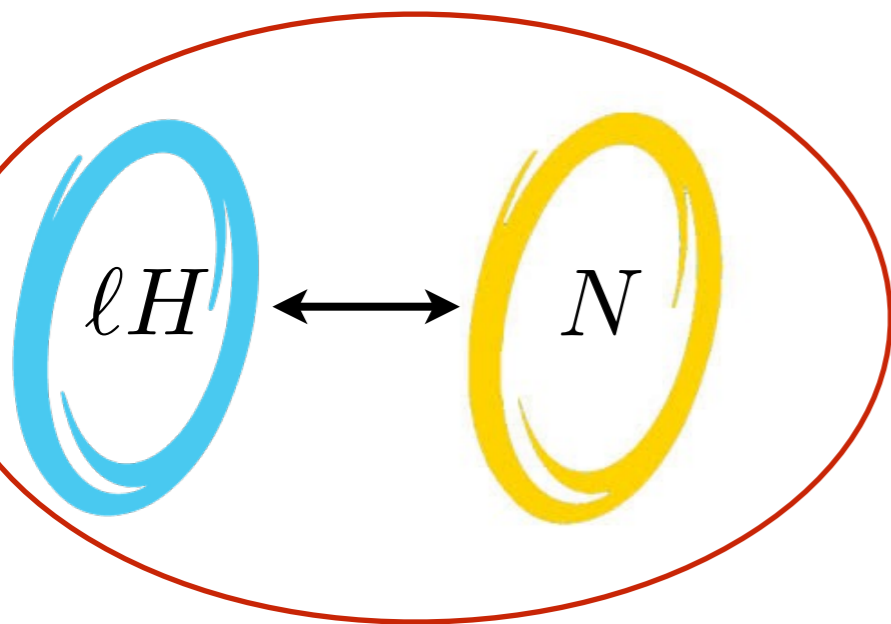
$\nu$



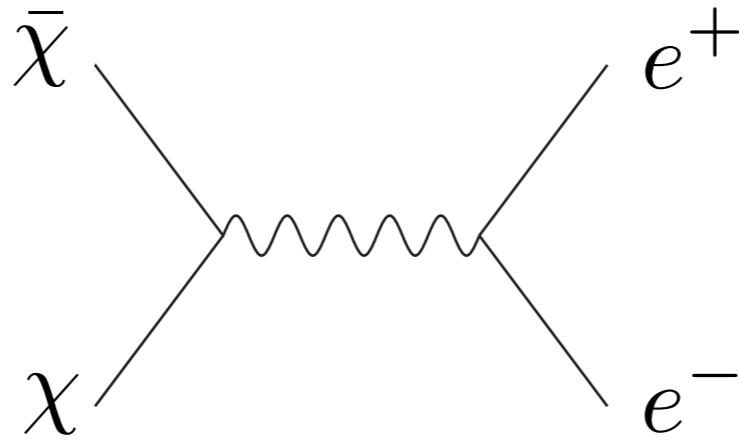
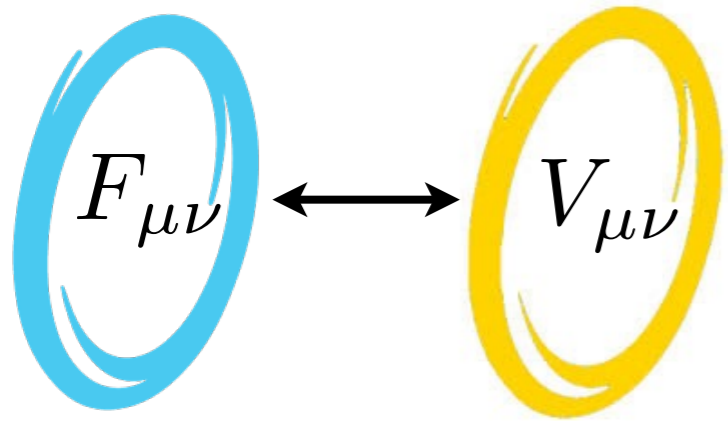
Vector Portal



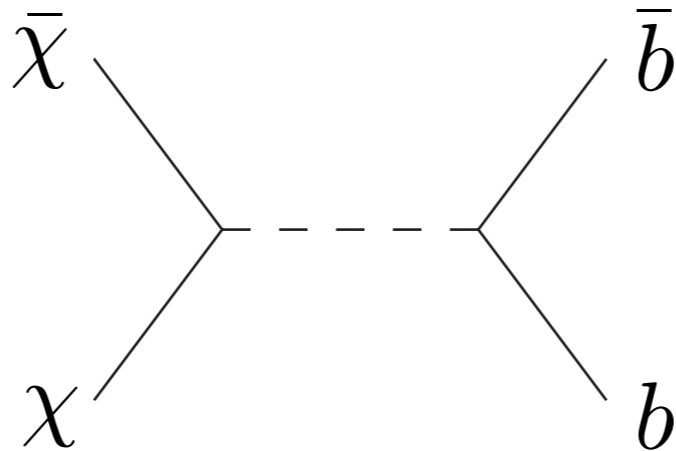
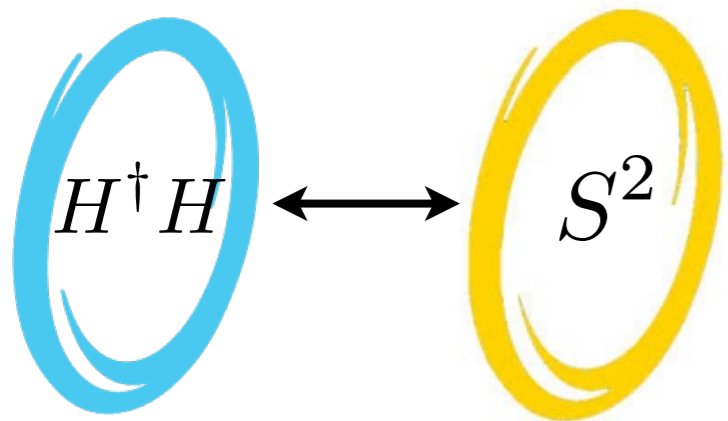
Higgs Portal



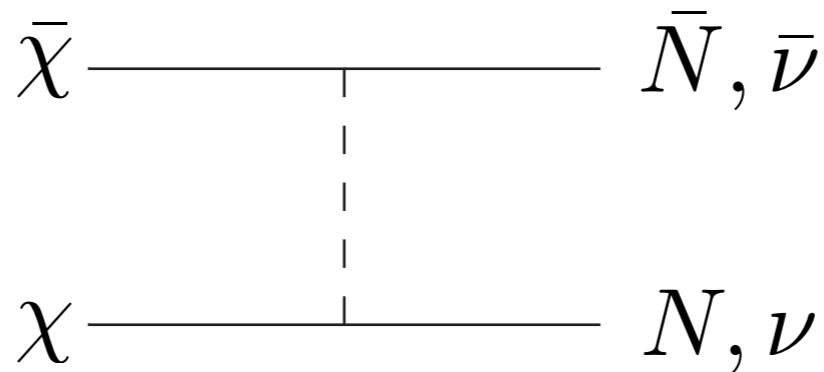
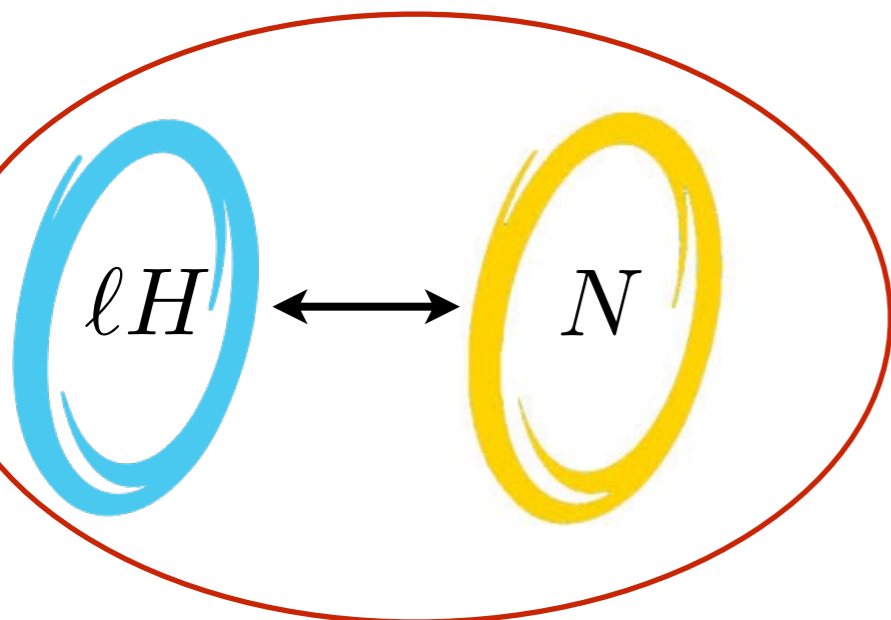
Neutrino Portal



Vector Portal



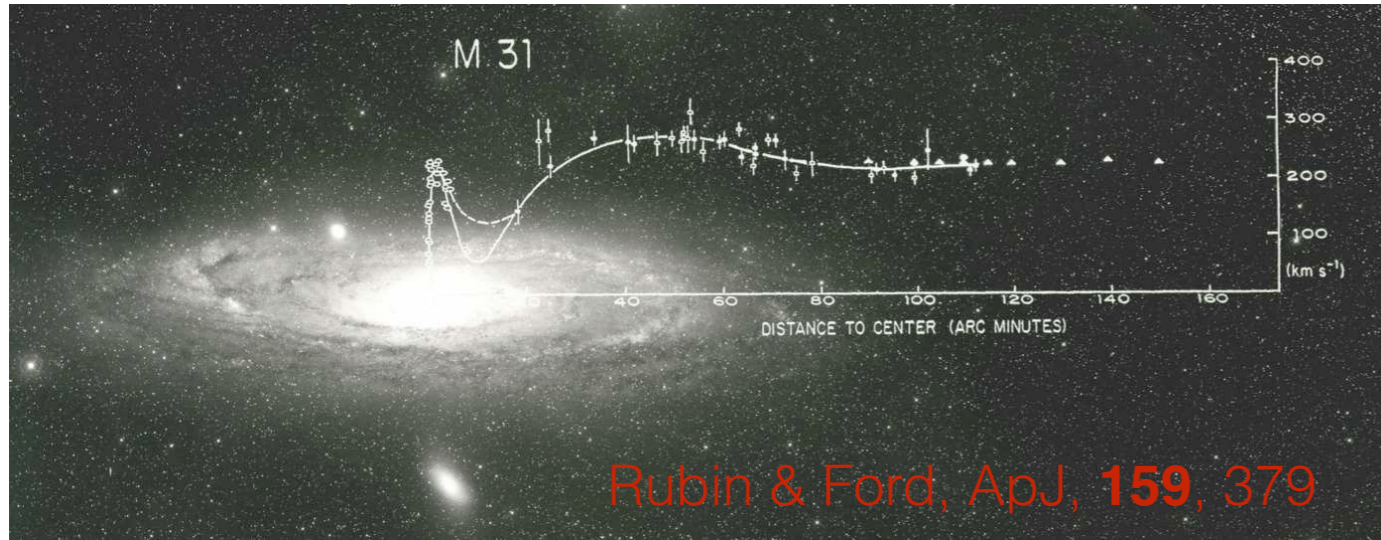
Higgs Portal



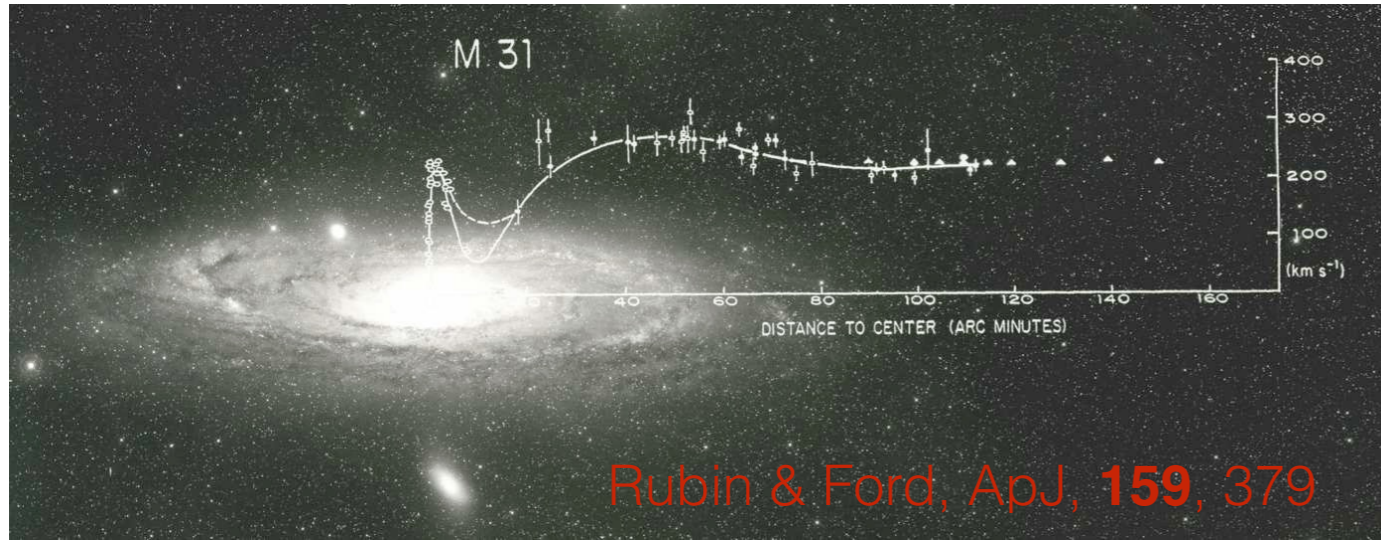
Neutrino Portal

(Maybe already used...)

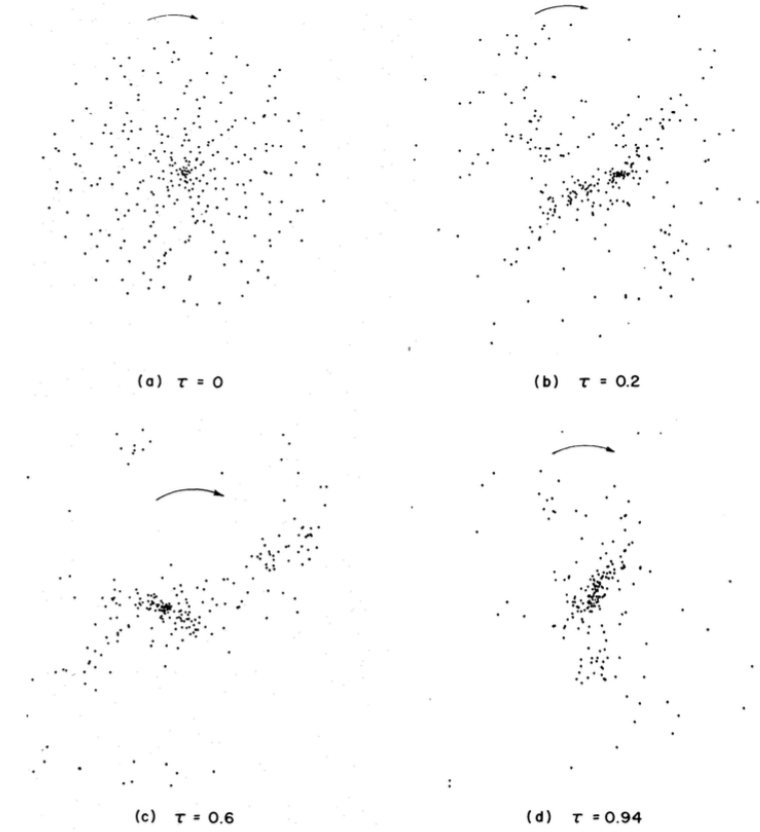
Why bother with a  
portal?





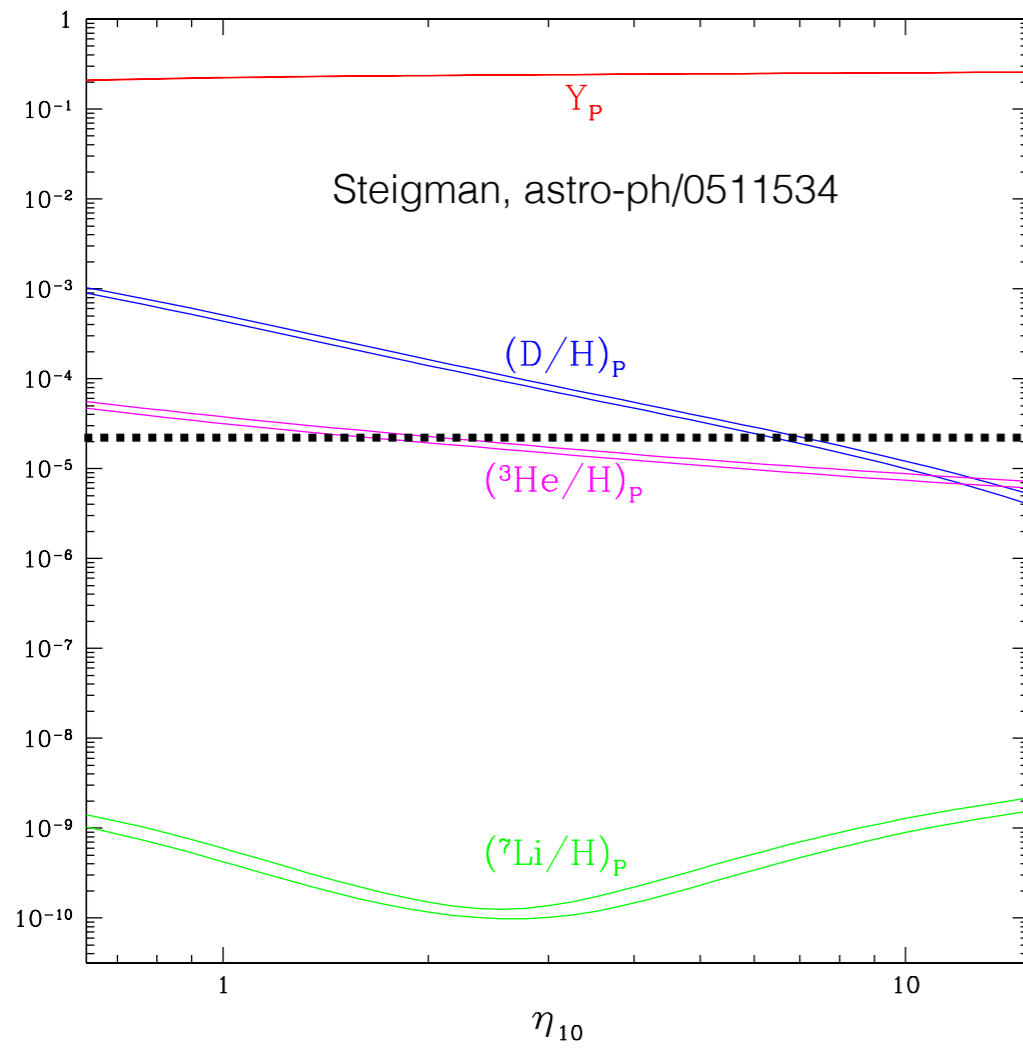
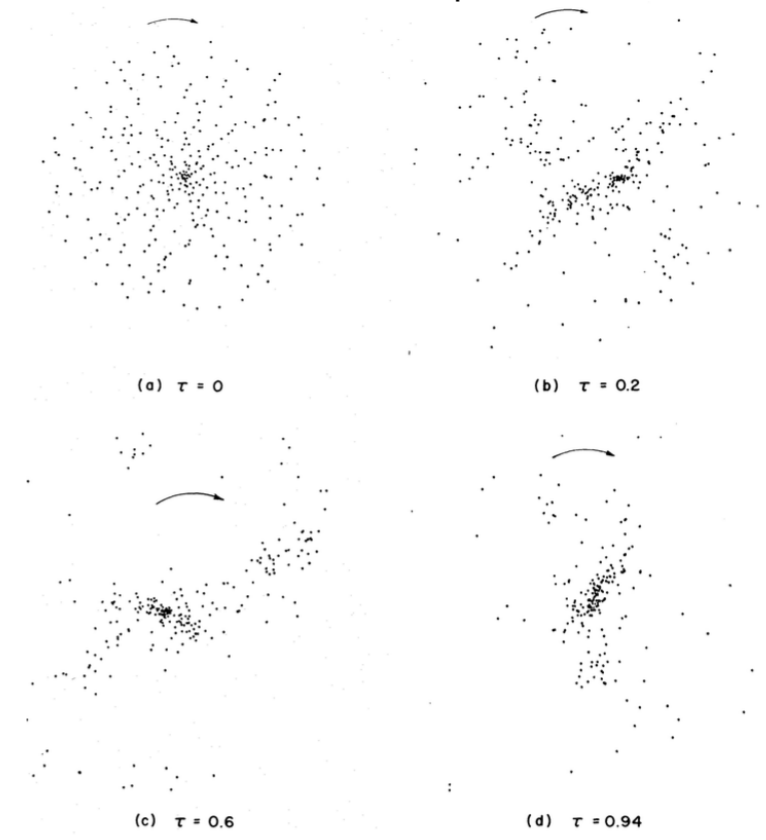
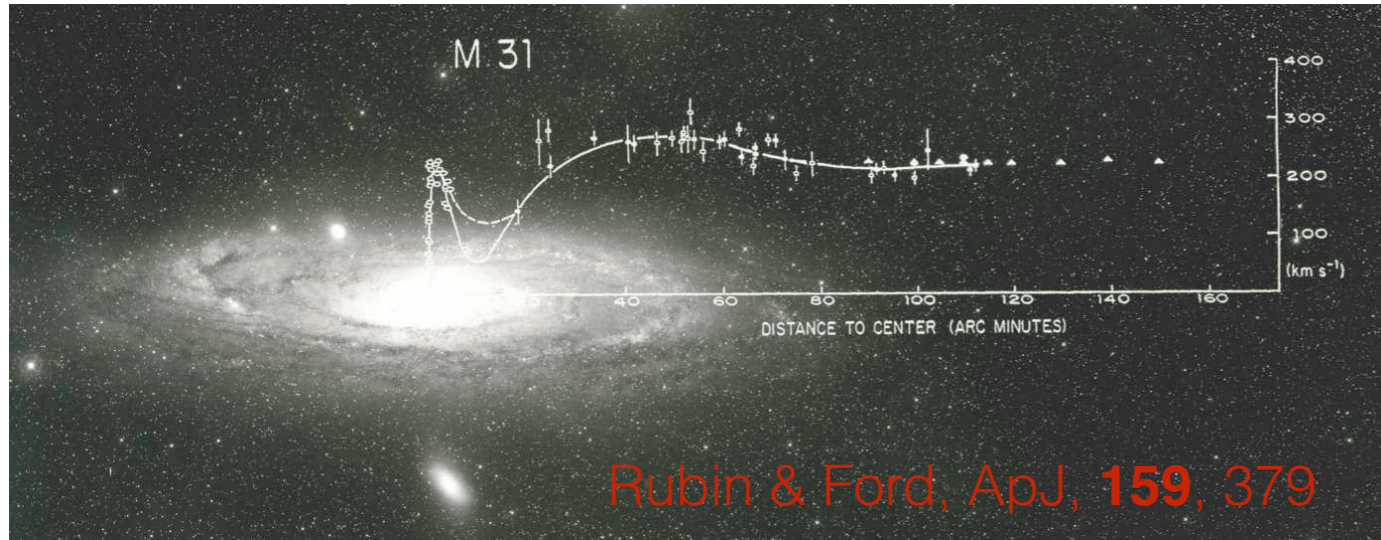


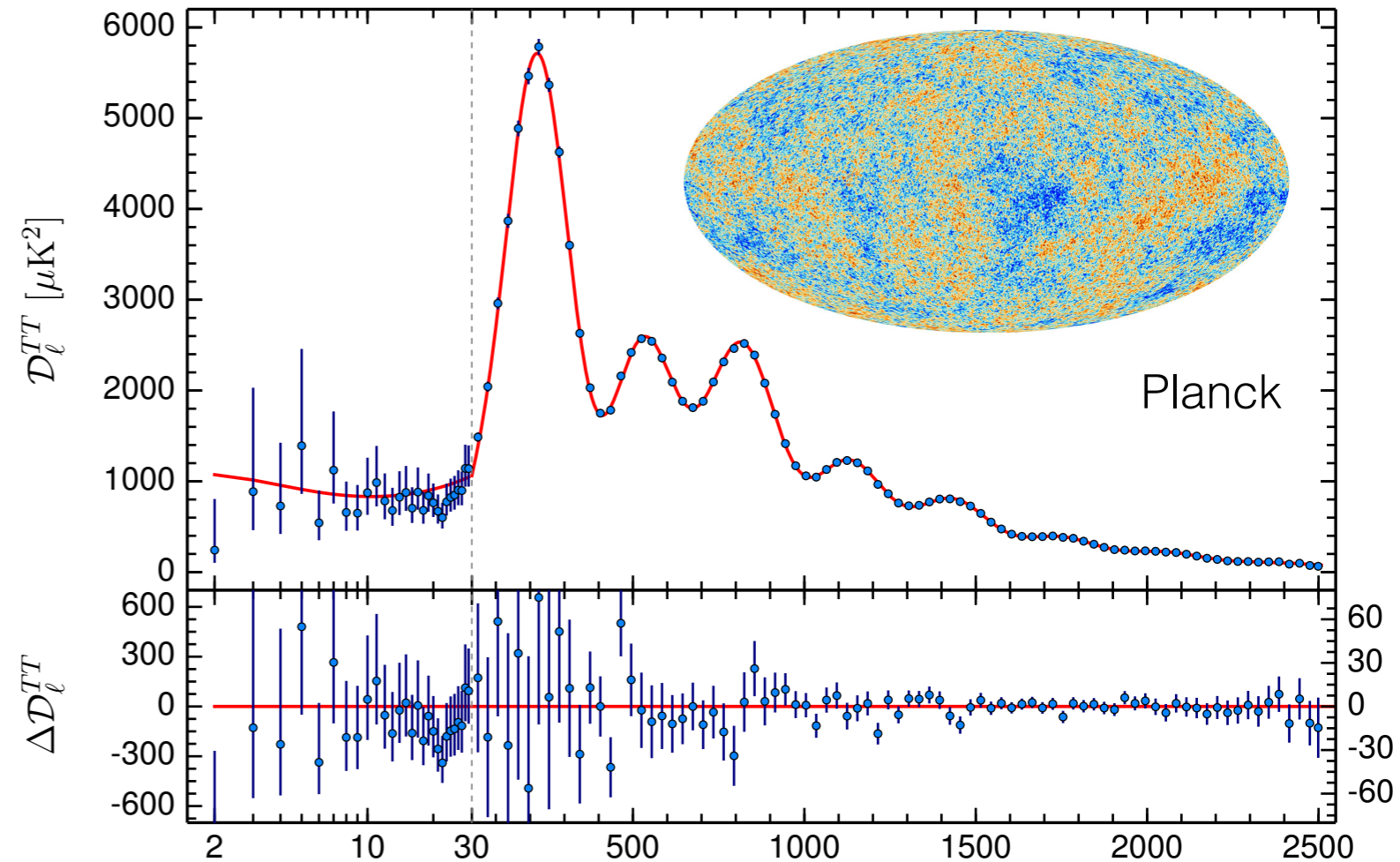
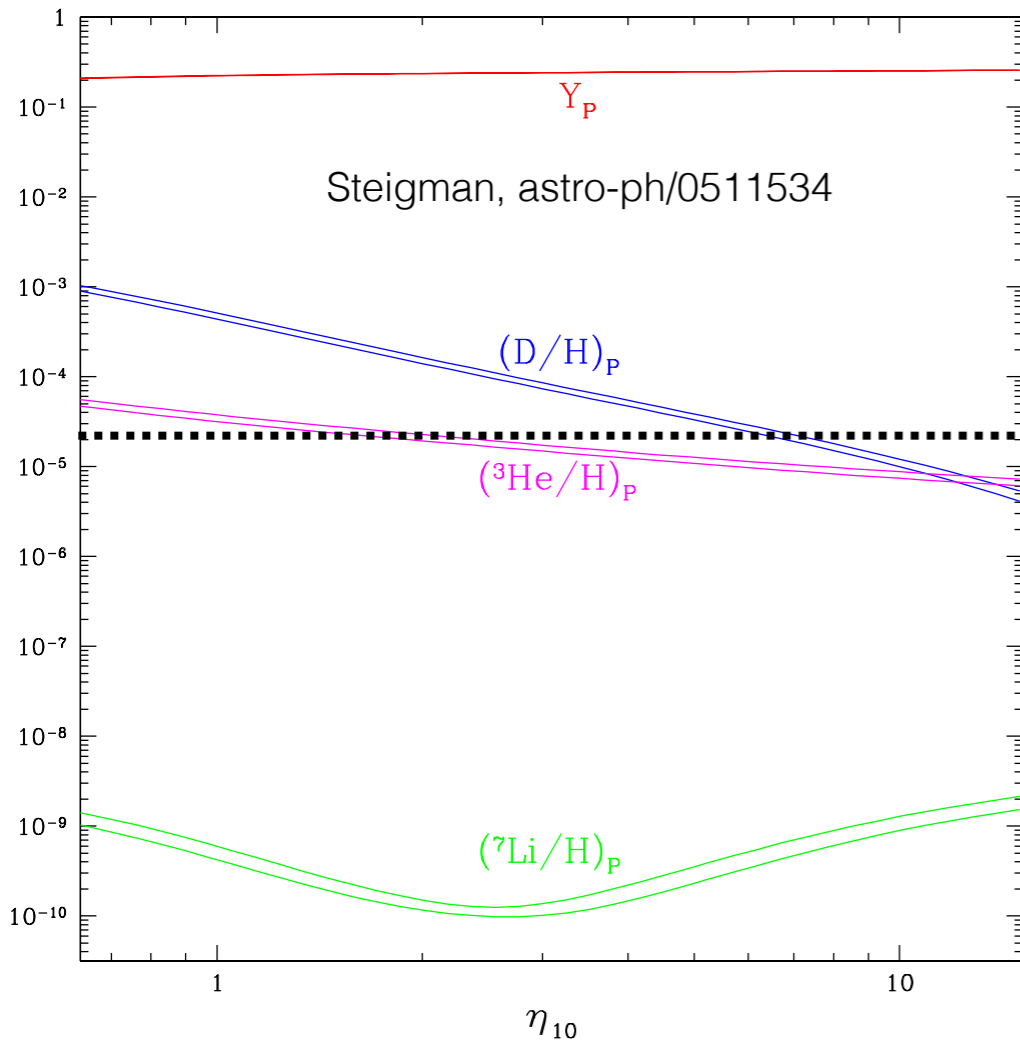
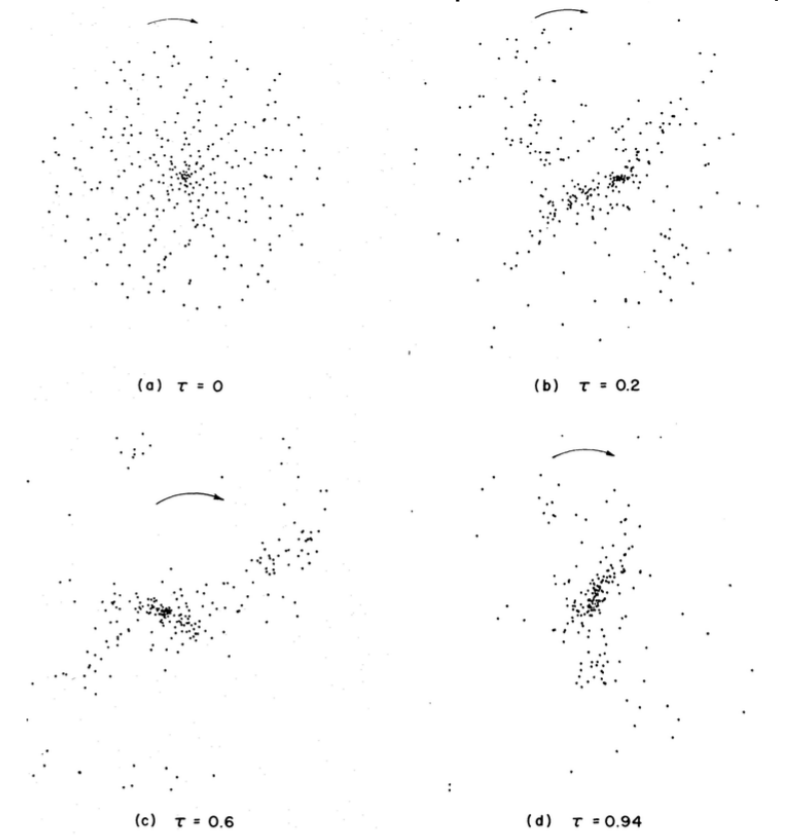
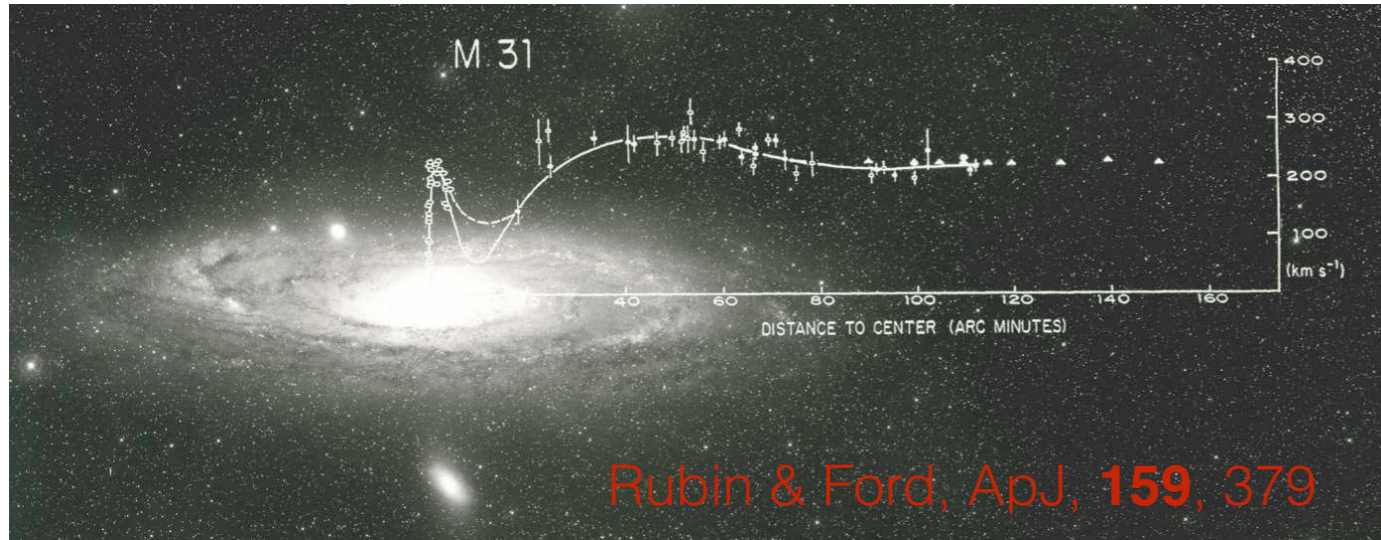
Ostriker & Peebles, ApJ, **186**, 467 ('73)



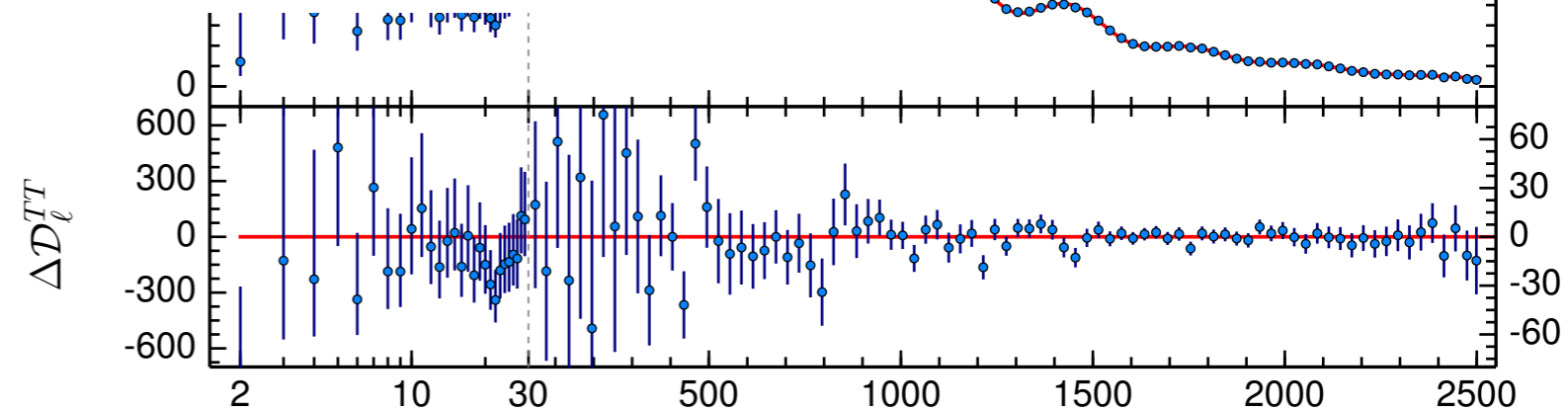
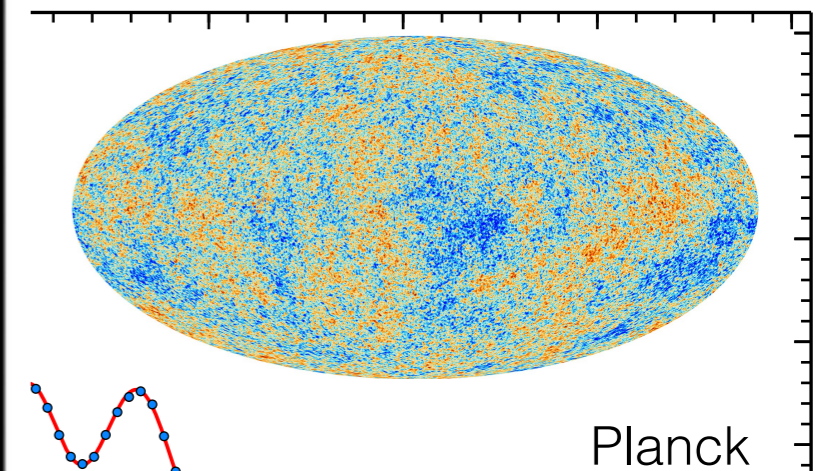
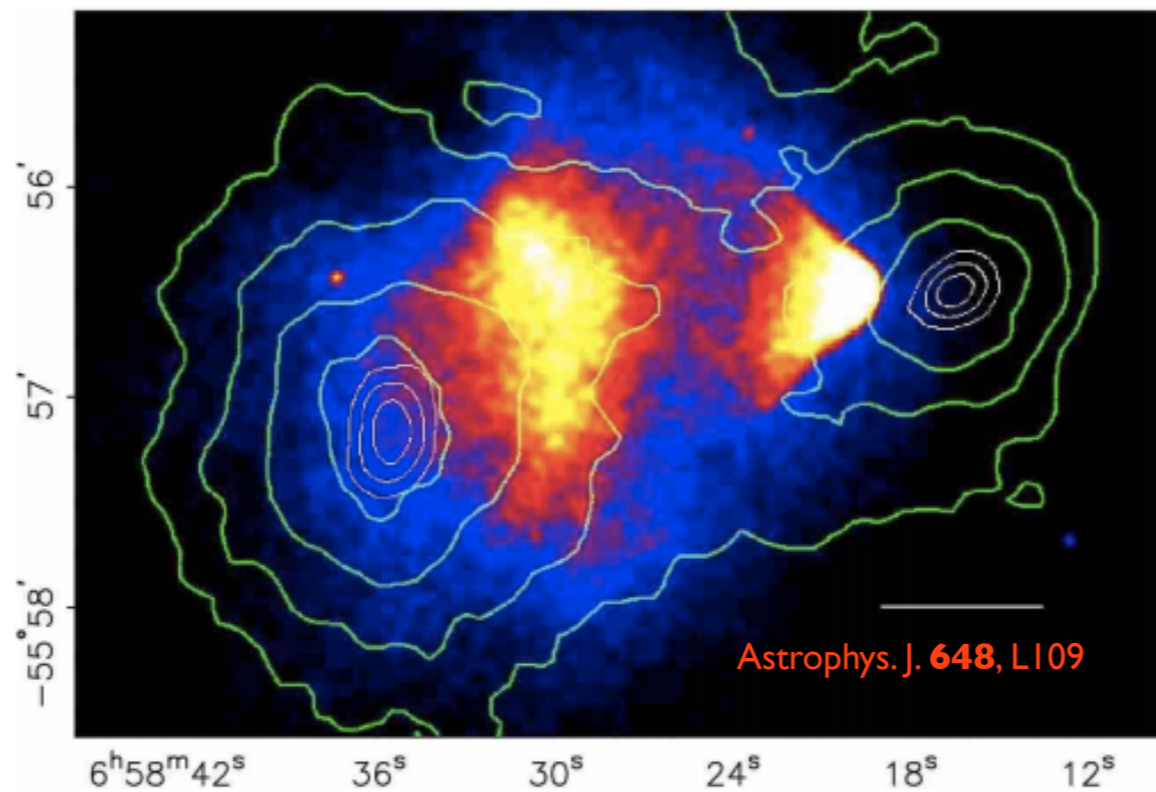
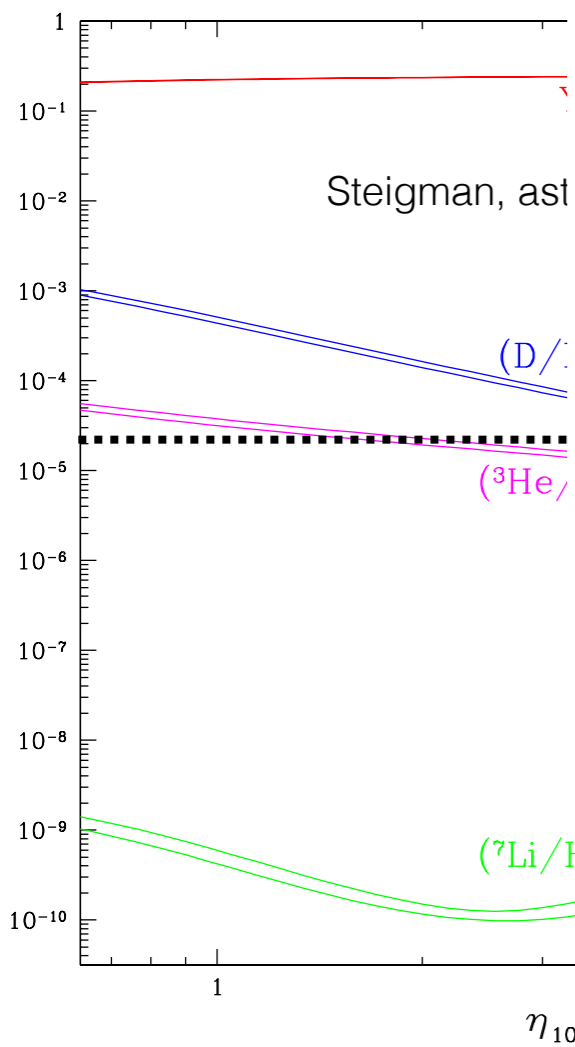
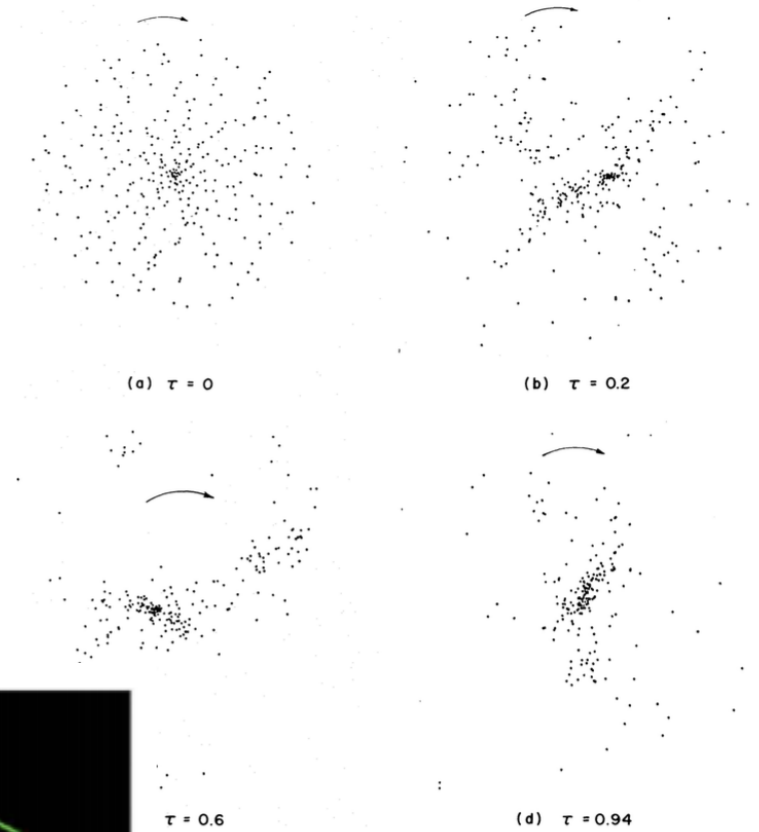
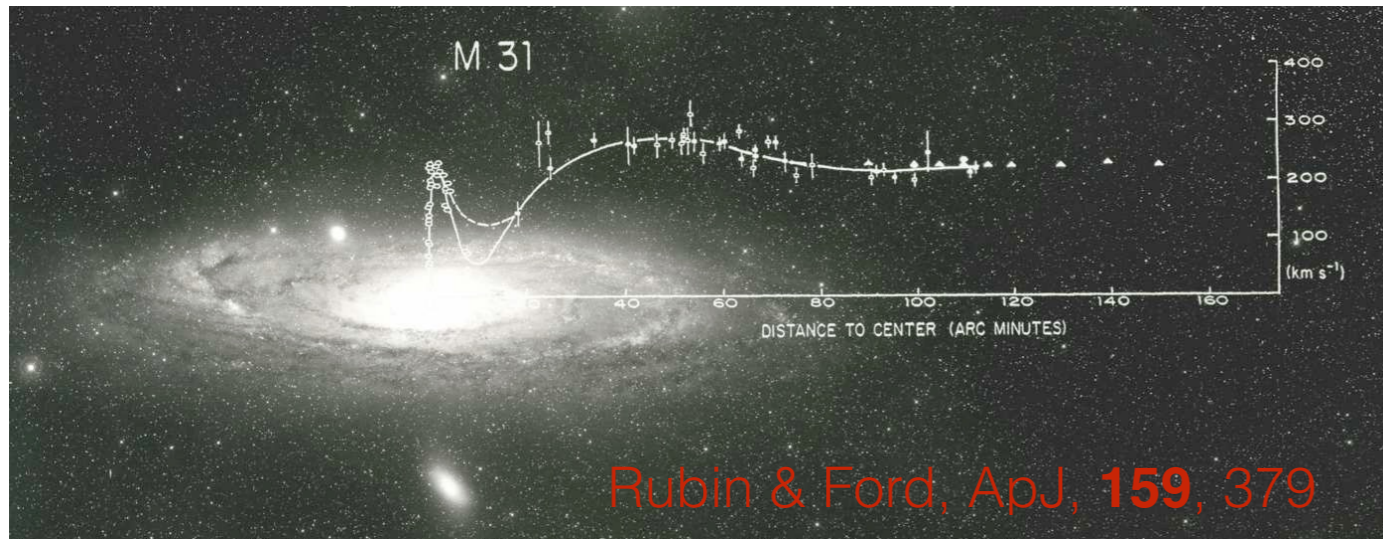


Ostriker & Peebles, ApJ, **186**, 467 ('73)





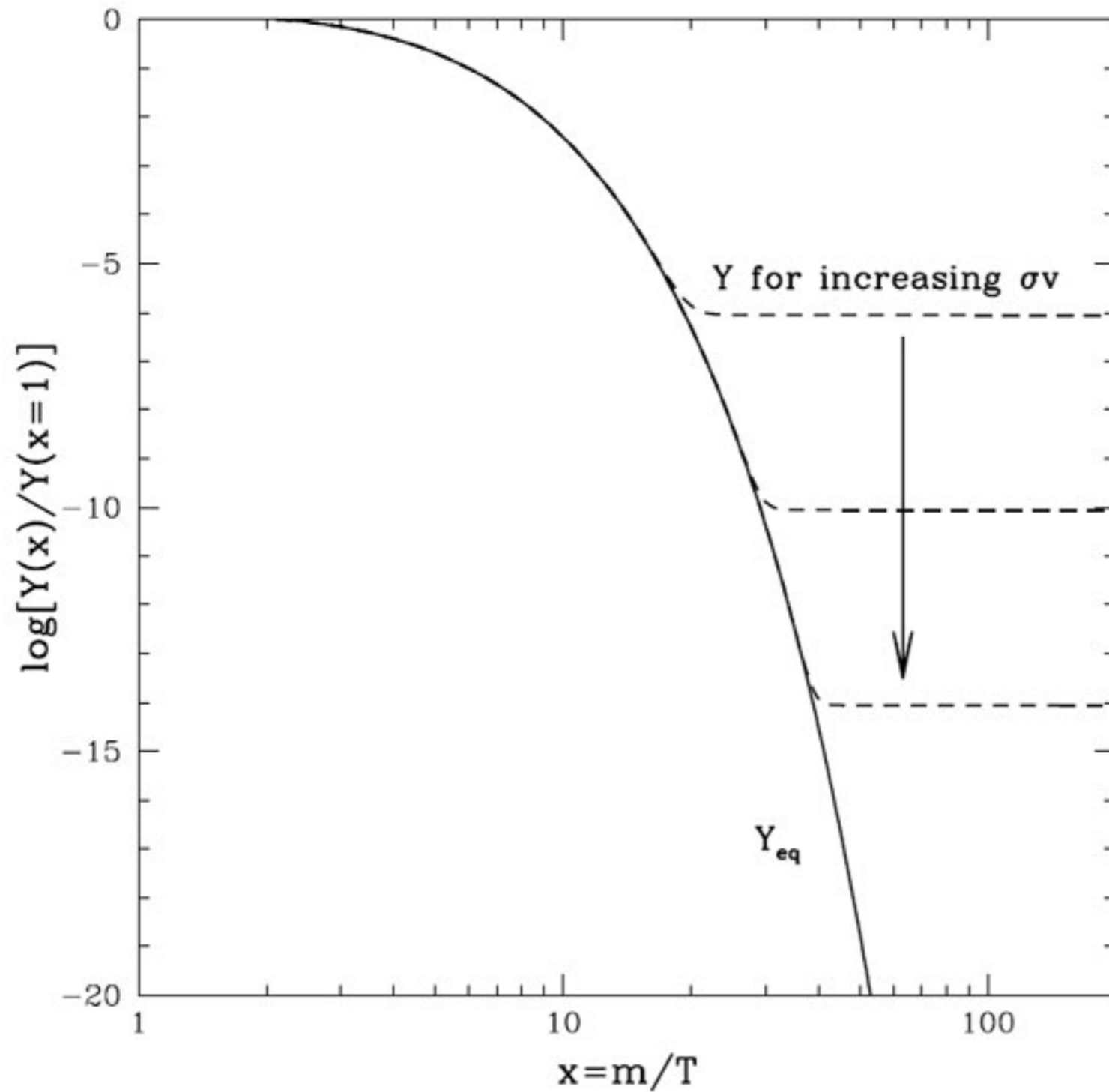




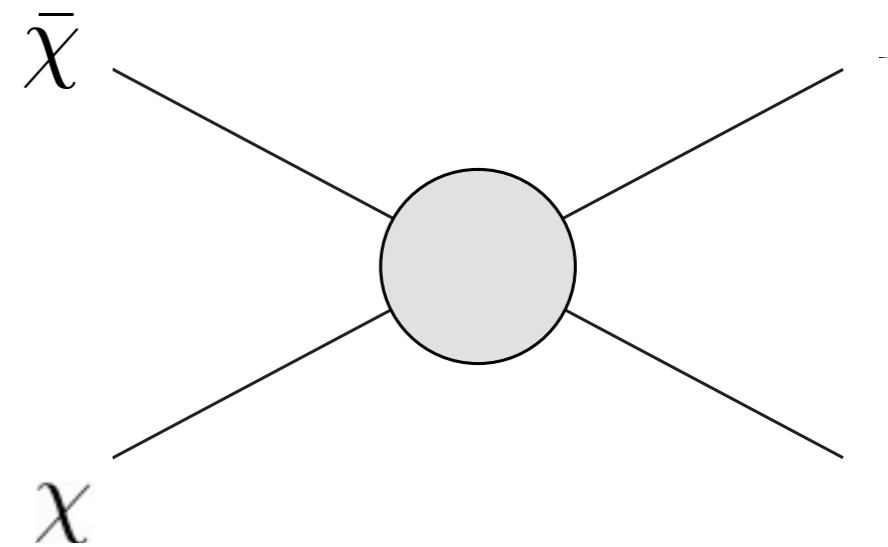
General consensus on energy budget of Universe



Consider thermal relic DM



Needs a coupling to light stuff



How do you generate  
that coupling

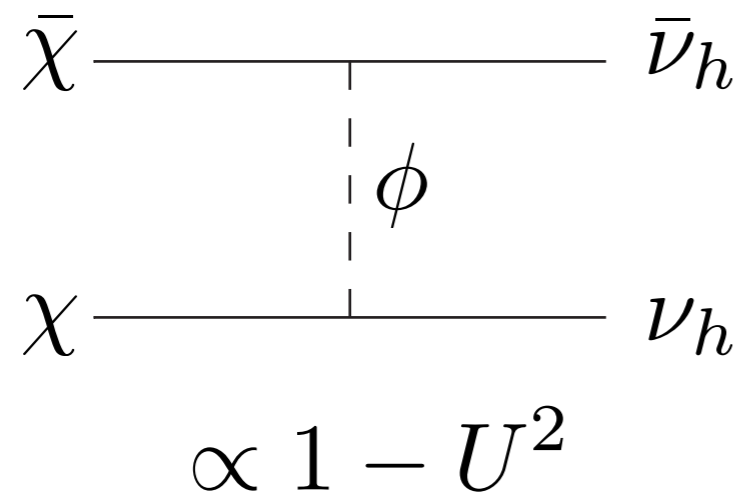
# Basic Idea

$$\mathcal{L} \supset -\lambda \bar{L} H N - y \bar{N} \chi \phi + \text{h.c.} \rightarrow -\lambda v \bar{\nu} N - y \bar{N} \chi \phi + \text{h.c.}$$

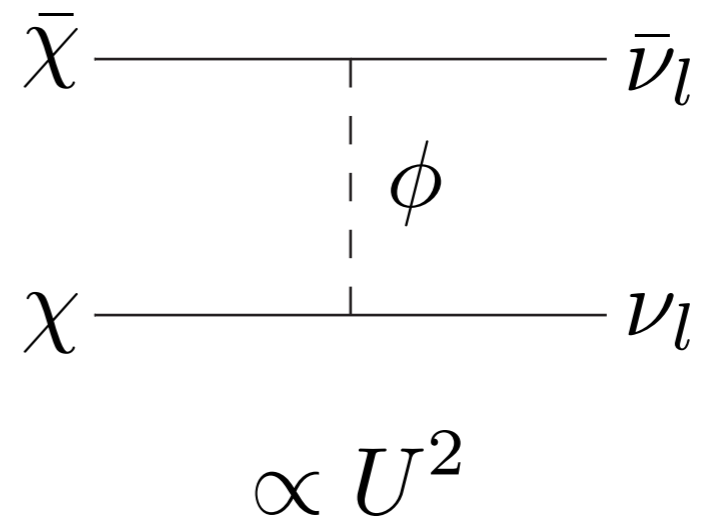
$$\nu_l = \sqrt{1 - U^2} \nu + U N$$

$$\nu_h = -U \nu + \sqrt{1 - U^2} N$$

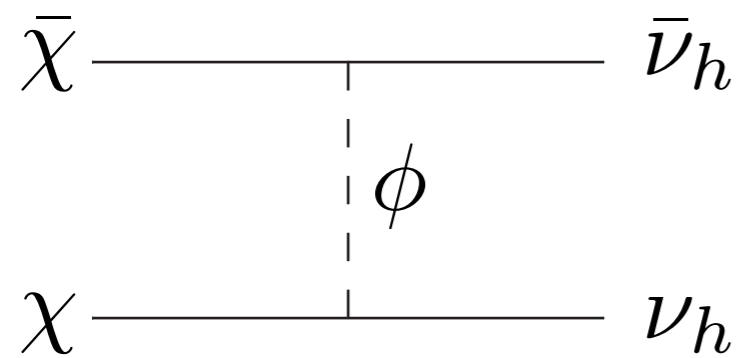
Case I:  $m_\chi > M$



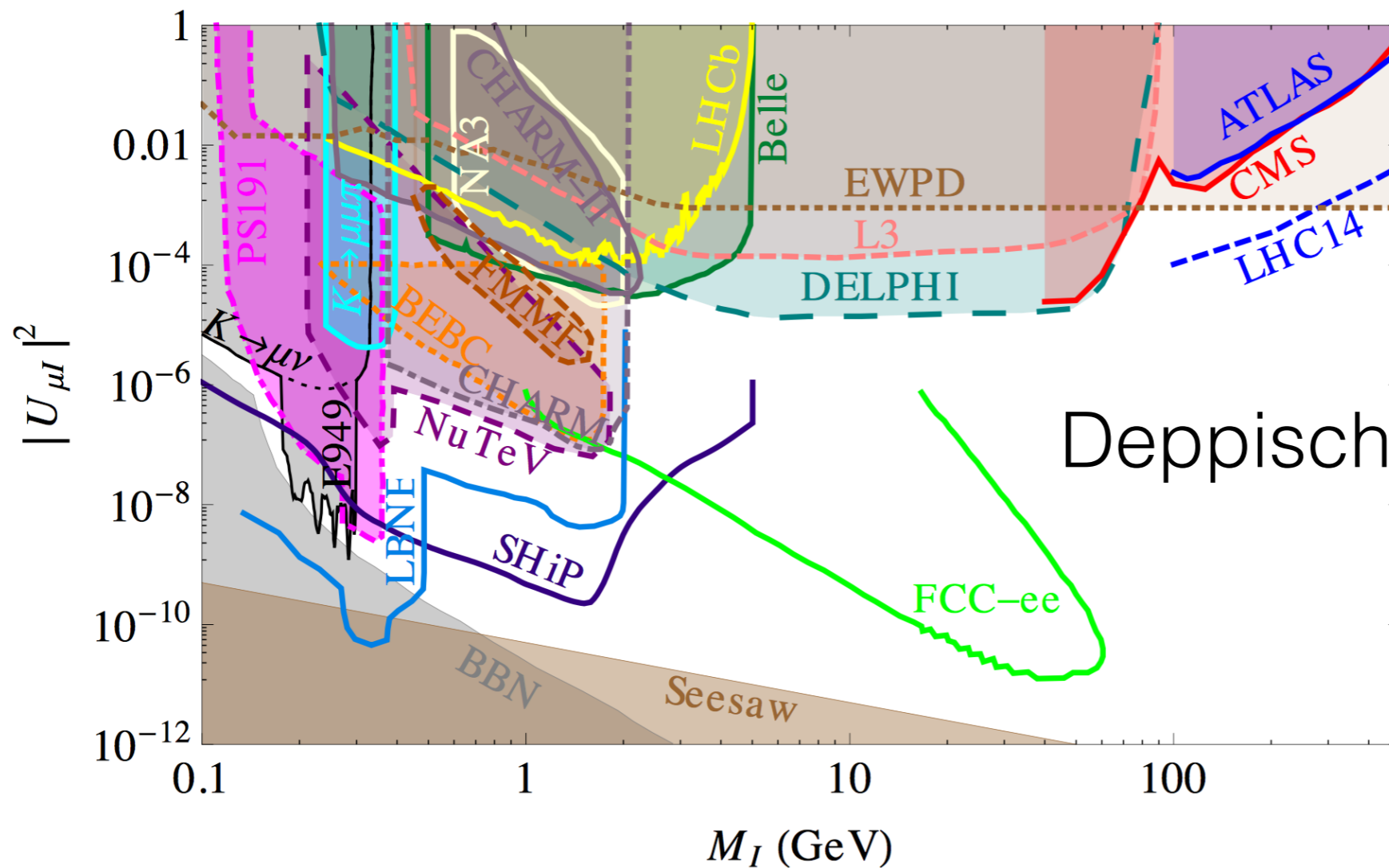
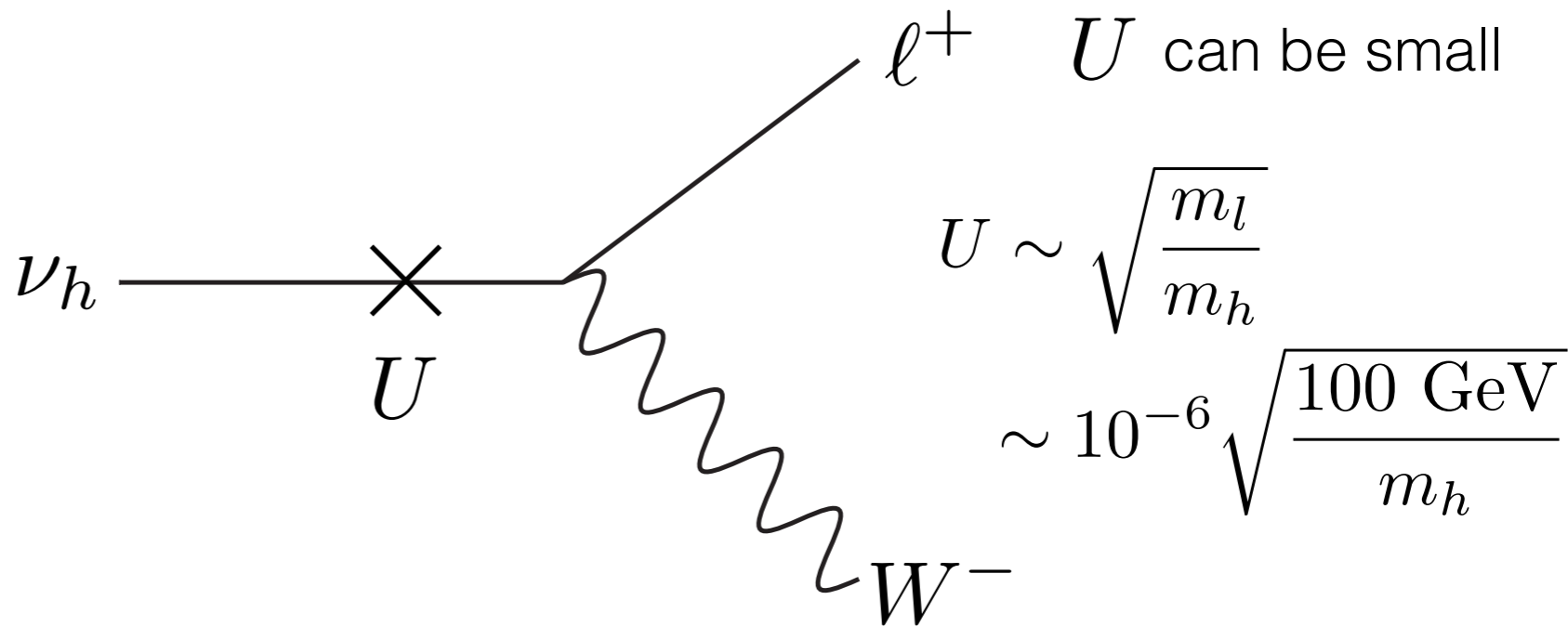
Case II:  $m_\chi < M$



Case I:  $m_\chi > M$



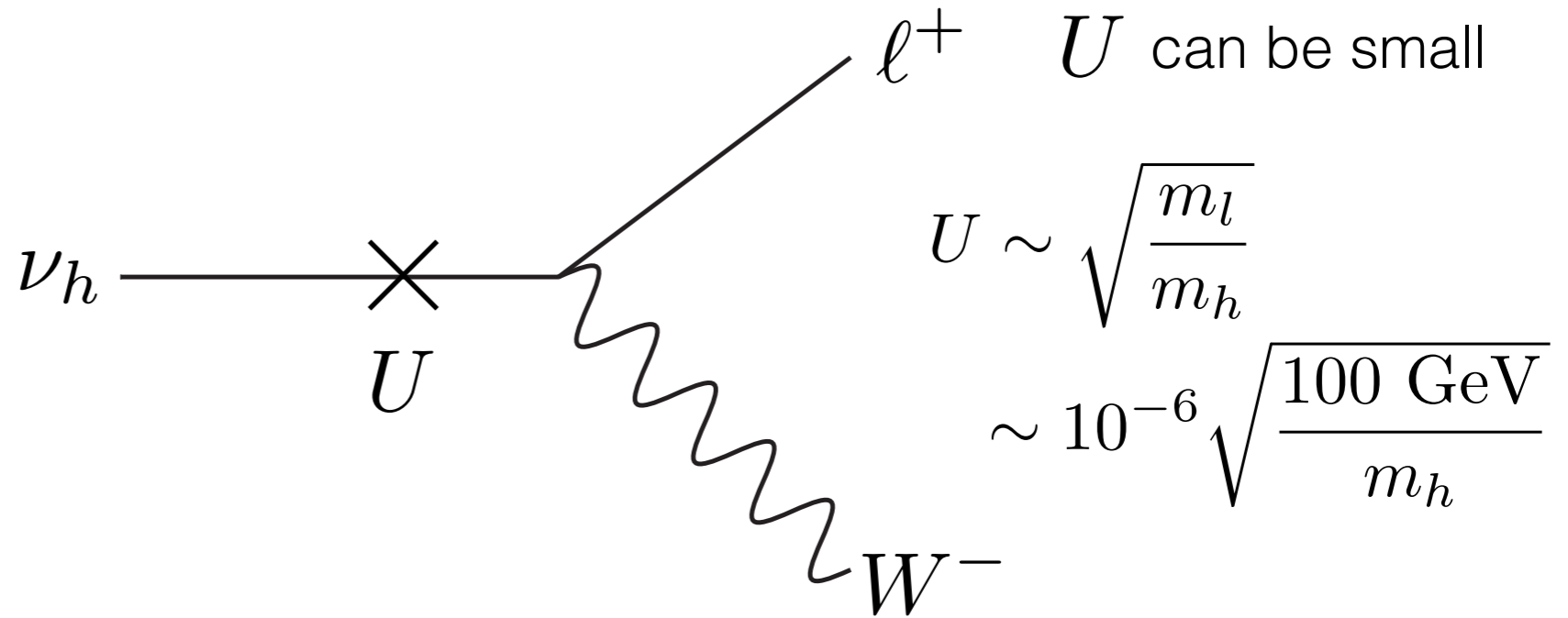
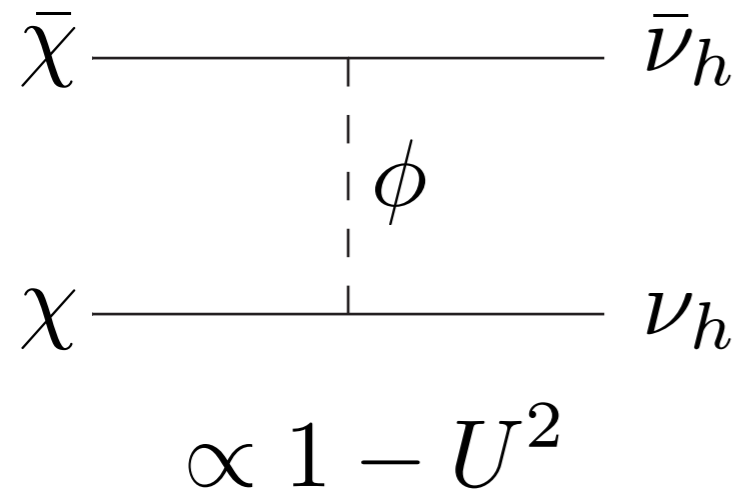
$$\propto 1 - U^2$$



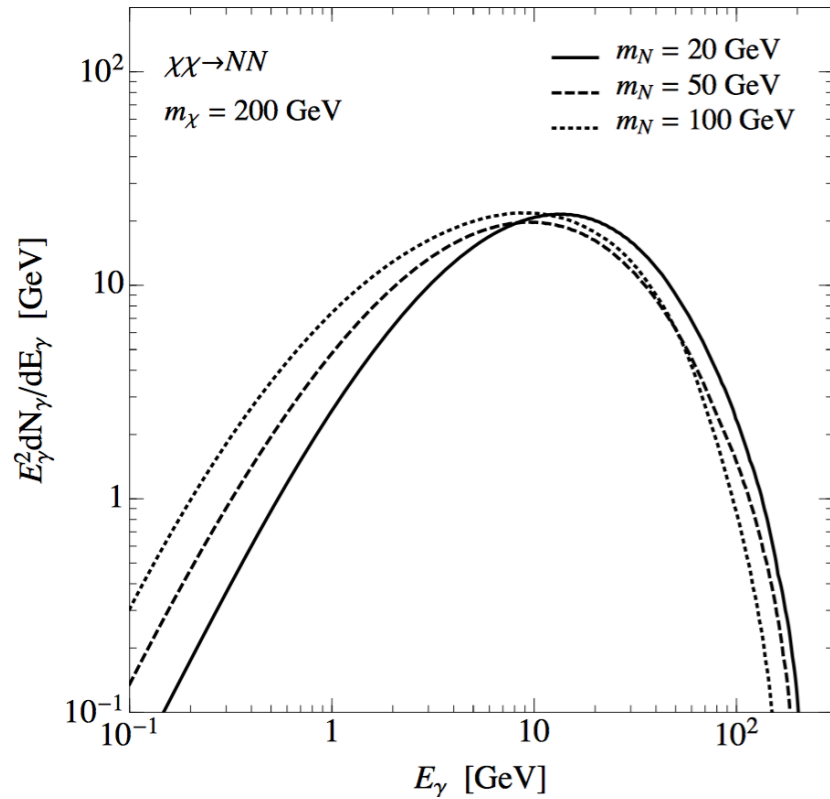
Deppisch, Dev, Pilaftsis



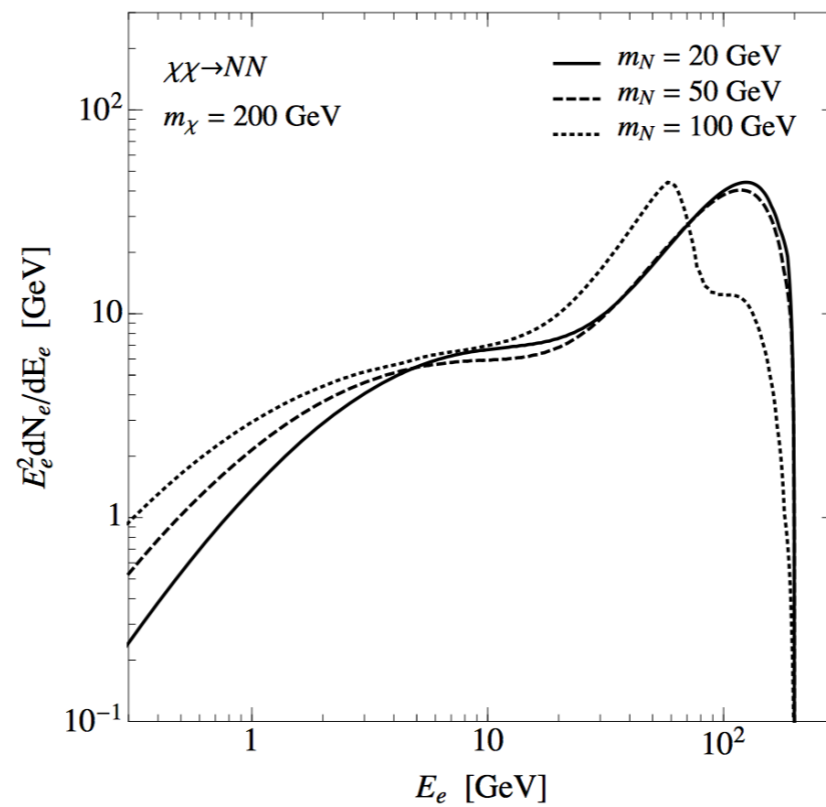
Case I:  $m_\chi > M$



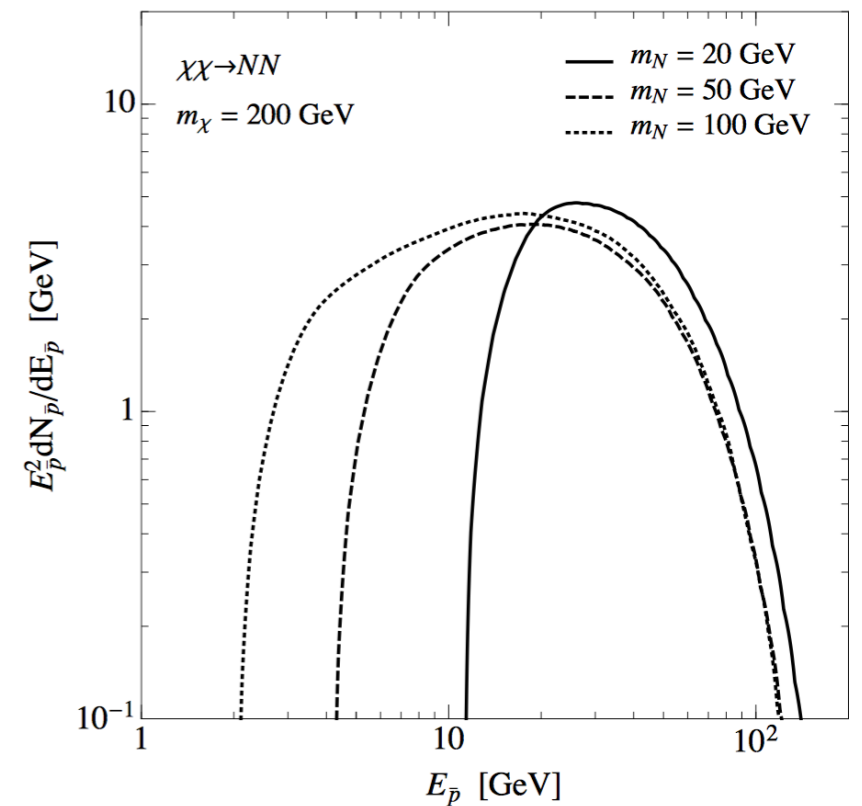
$\gamma$  Primary Spectra



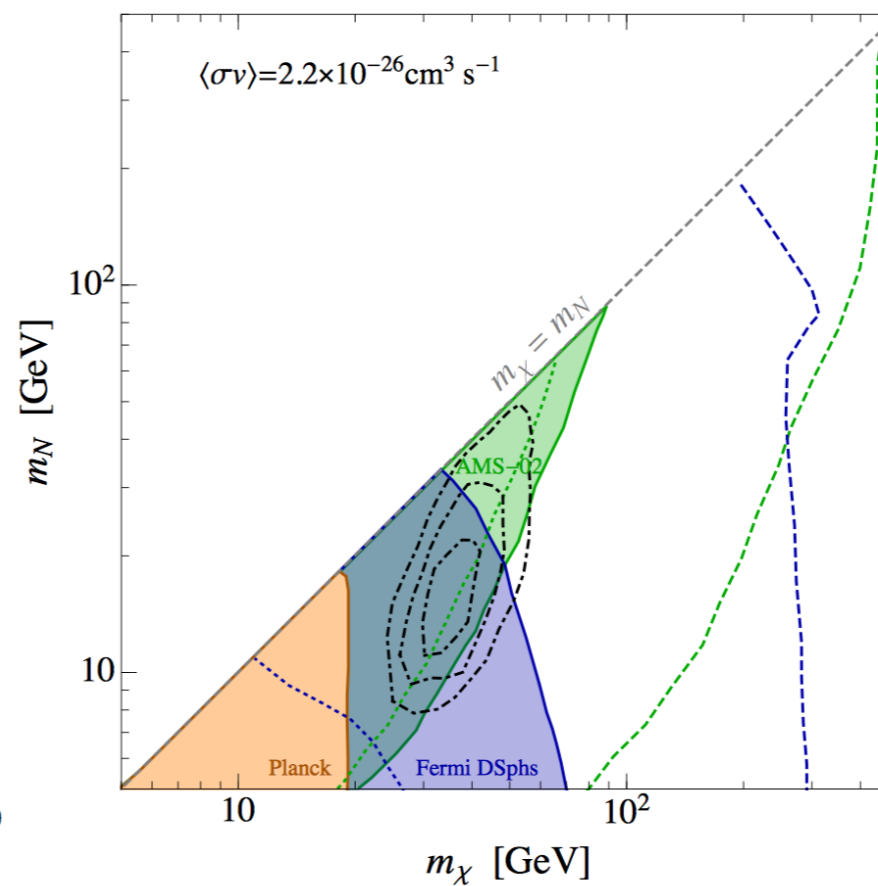
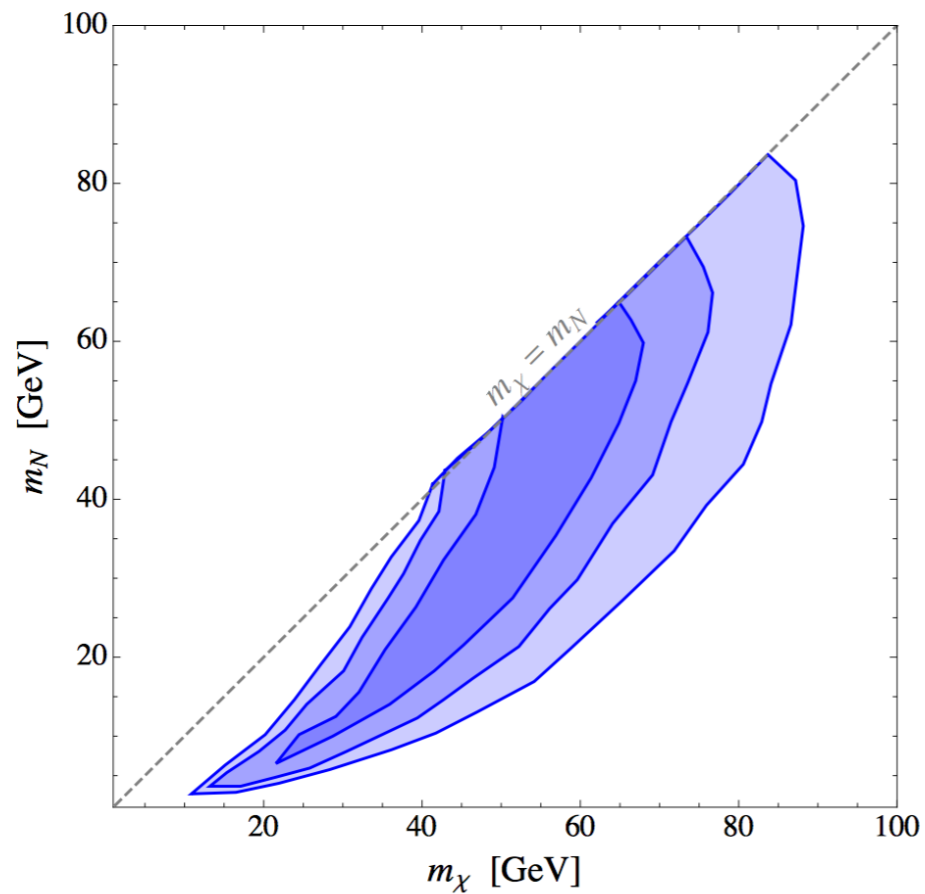
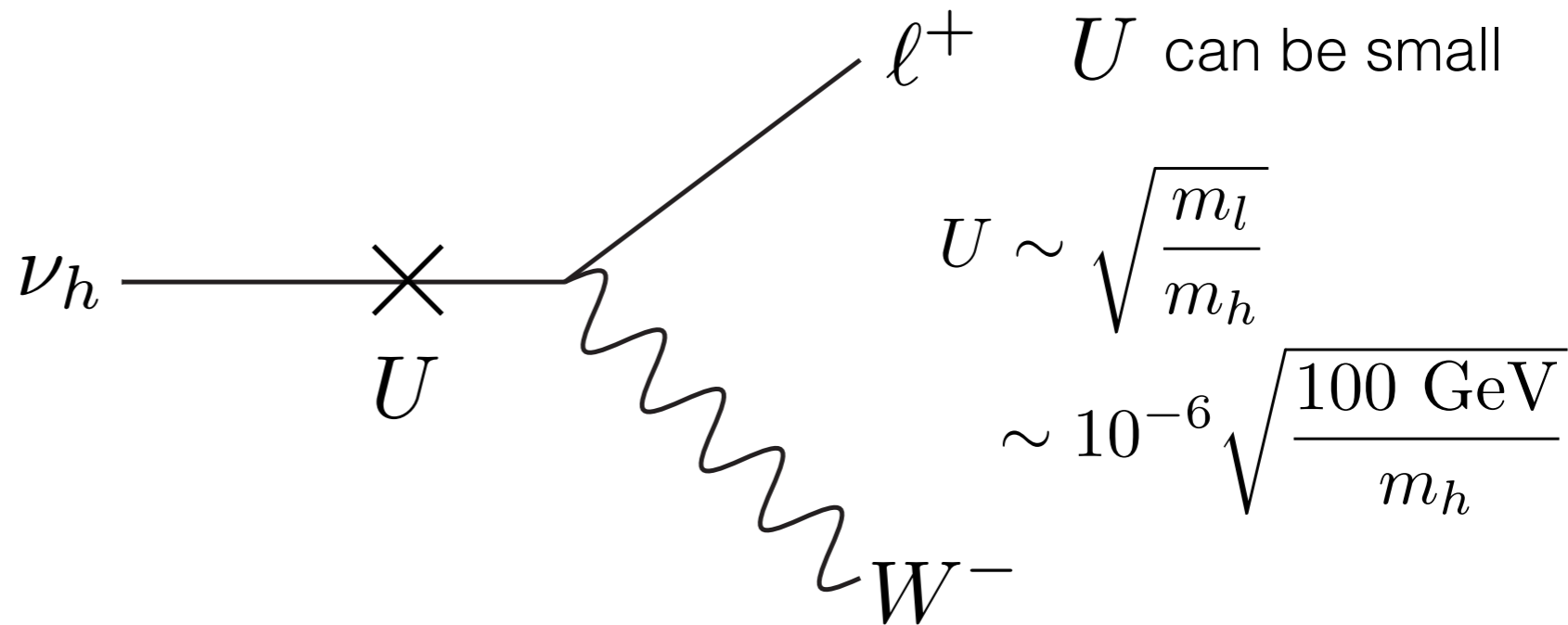
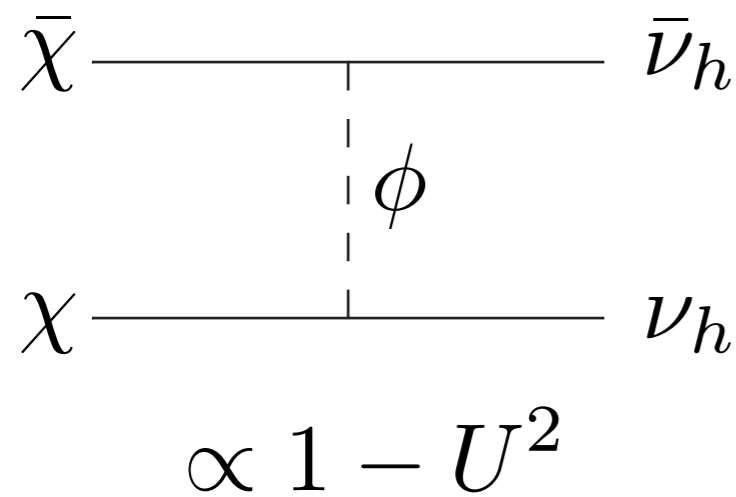
$e^-$  Primary Spectra



$\bar{p}$  Primary Spectra



Case I:  $m_\chi > M$



Batell, Han, & Shams Es Haghi 1704.08708

Case II:  $m_\chi < M$

$\bar{\chi}$  —————  $\bar{\nu}_l$

|  
ϕ  
|

$\chi$  —————  $\nu_l$

$\propto U^2$

$$\langle \sigma v \rangle = \frac{y^4 U^4}{32\pi} \frac{m_\chi^2}{m_\phi^4} \left( 1 + \frac{m_\chi^2}{m_\phi^2} \right)^{-2}$$

$$\simeq 3 \times 10^{-26} \frac{\text{cm}^3}{\text{s}} \left( \frac{yU}{0.2} \right)^4 \left( \frac{m_\chi}{10 \text{ GeV}} \right)^{-2} \left( \frac{m_\phi/m_\chi}{3} \right)^{-4}$$

Requires large mixing angle!

## Minimal Model

$$\mathcal{L} \supset -\lambda_i \bar{L}_i H N_R - M_N \bar{N}_L N_R - \phi \bar{\chi} (y_L N_L + y_R N_R) + \text{h.c.}$$

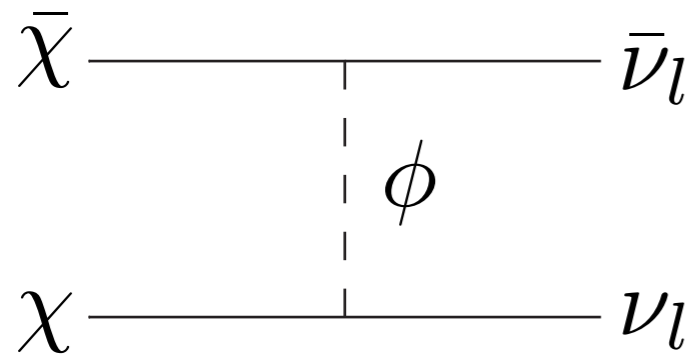
$$\rightarrow -\lambda_i \nu \bar{\nu}_{iL} N_R - M_N \bar{N}_L N_R - \phi \bar{\chi} (y_L N_L + y_R N_R) + \text{h.c.}$$

lepton number conserved (for small  $\nu$  masses & large mixing)

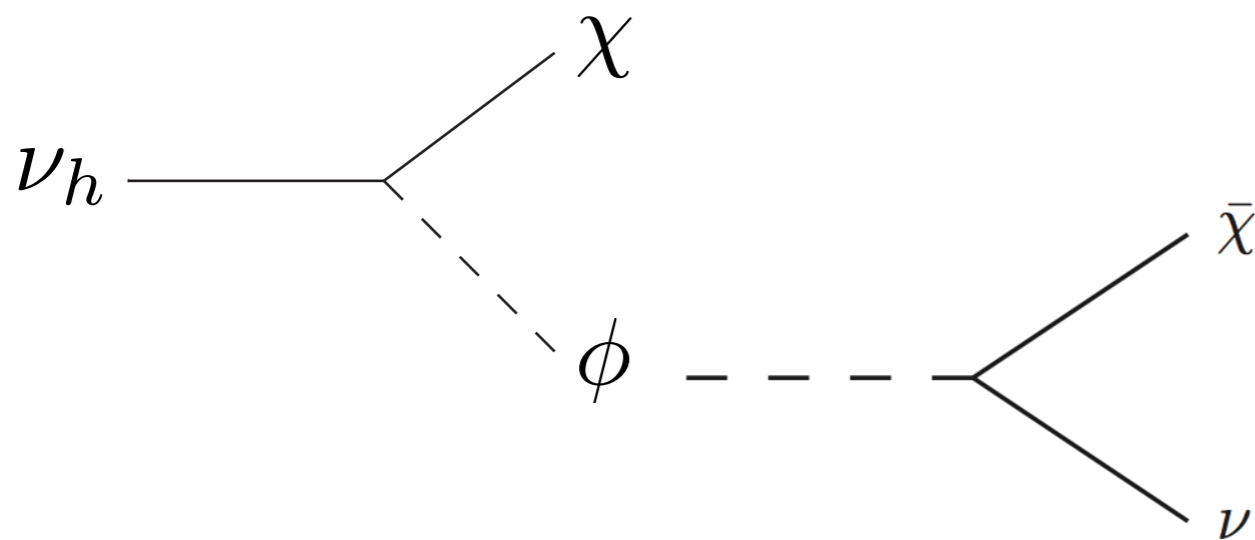
Bertoni, Ipek, DM, & Nelson 1412.3113

Batell, Han, DM, & Shams Es Haghi *in prep.*

Case II:  $m_\chi < M$



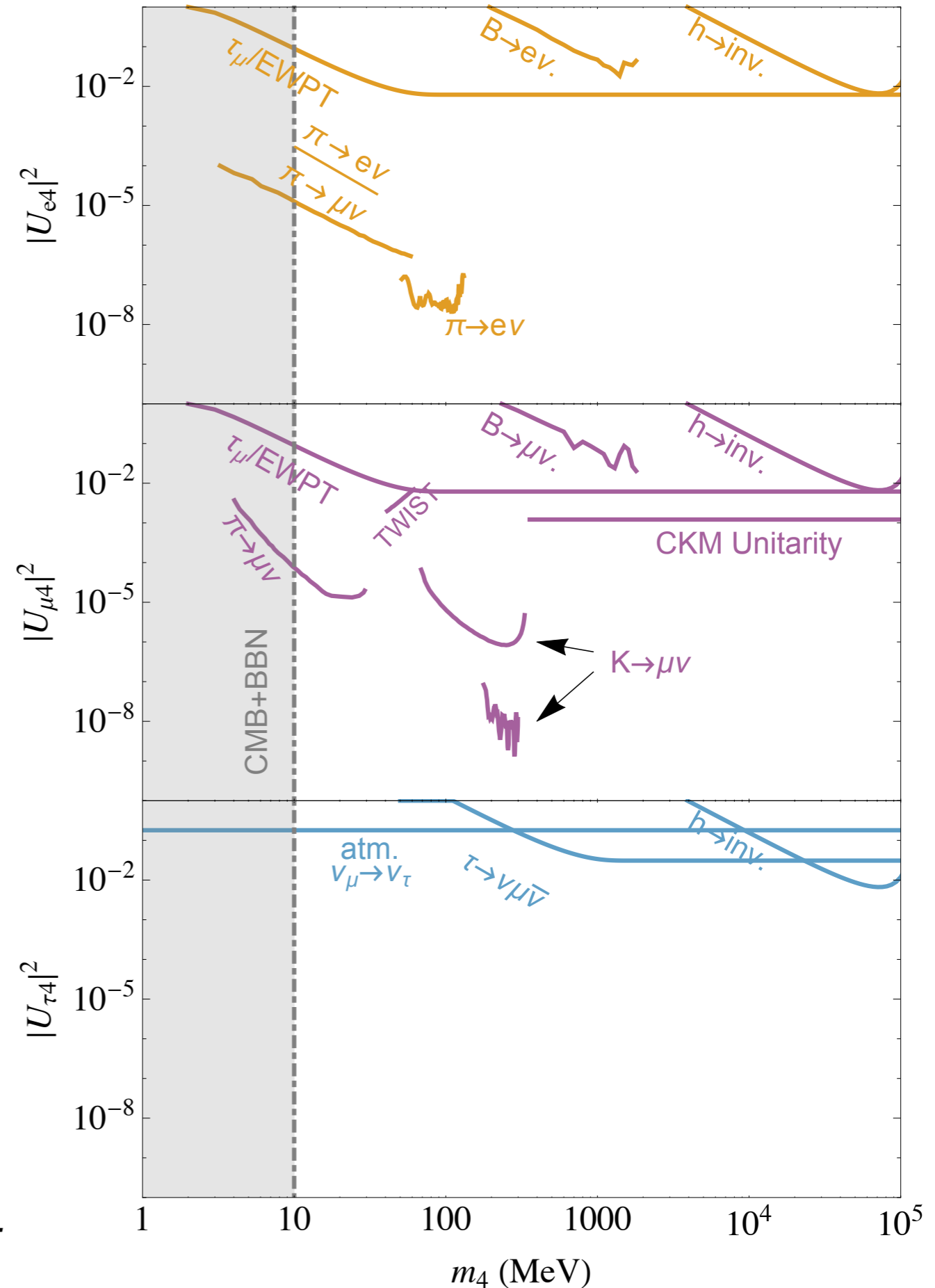
$\propto U^2$



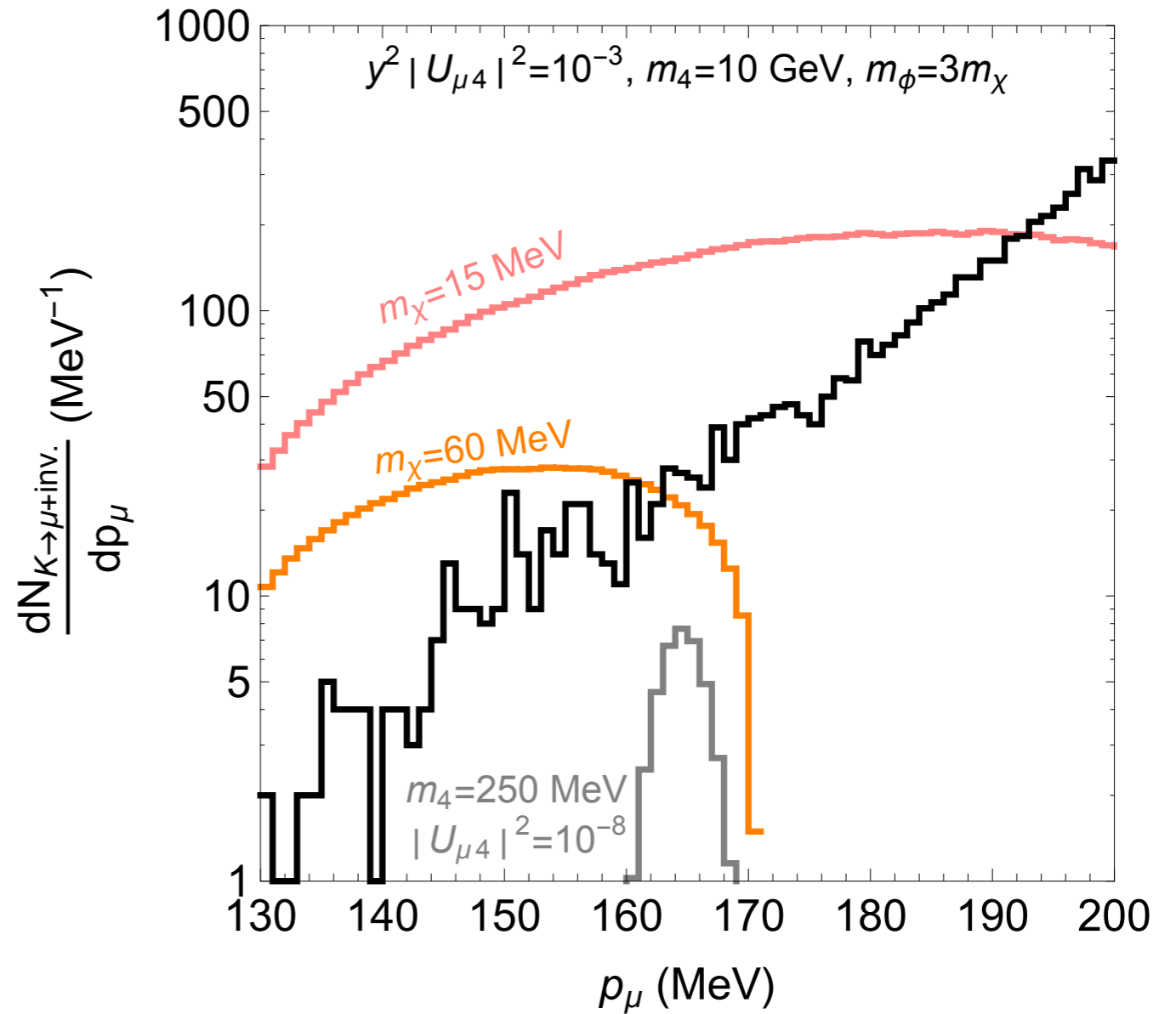
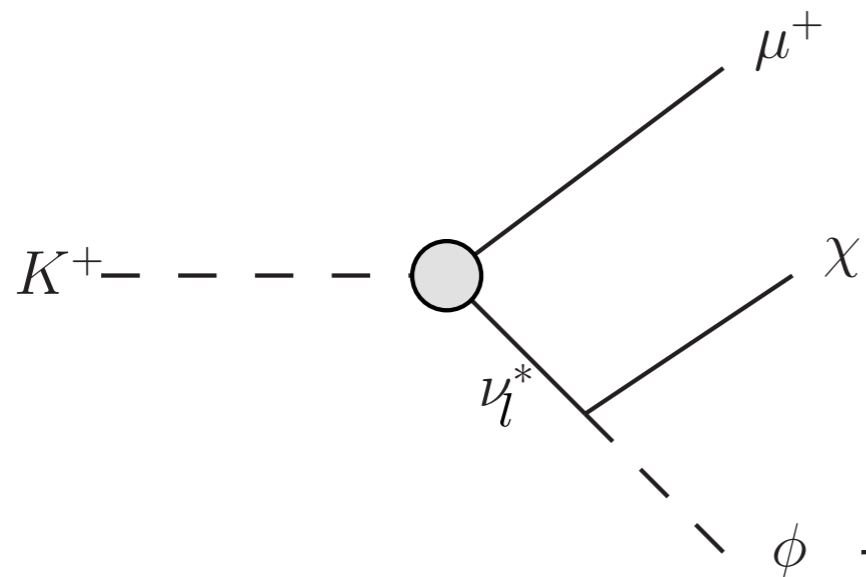
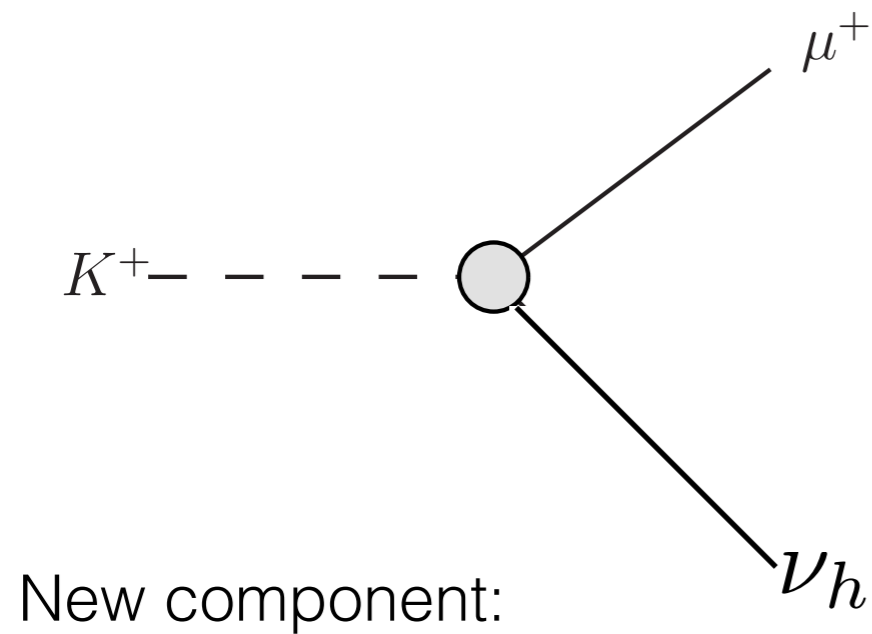
heavy neutrino  
decays invisibly!

Bertoni, Ipek, DM, & Nelson 1412.3113

Batell, Han, DM, & Shams Es Haghi *in prep.*



# Meson decays



E949:  $10^{12}$  kaons

NA62 increase by  
~order of mag.

# Atmospheric Neutrino Oscillations

$\nu_\mu, \nu_{\tau N}$  Hamiltonian:

$$H = \left( \frac{\Delta m^2}{4E} \right) \begin{pmatrix} -\cos 2\theta & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} + \begin{pmatrix} V_\mu & 0 \\ 0 & V_{\tau N} \end{pmatrix}$$

$$V_\mu = -\frac{G_F}{\sqrt{2}} n_n \sim \frac{1}{4000 \text{ km}}$$

$$V_{\tau N} = -\frac{G_F}{\sqrt{2}} n_n \cos \theta_\tau$$

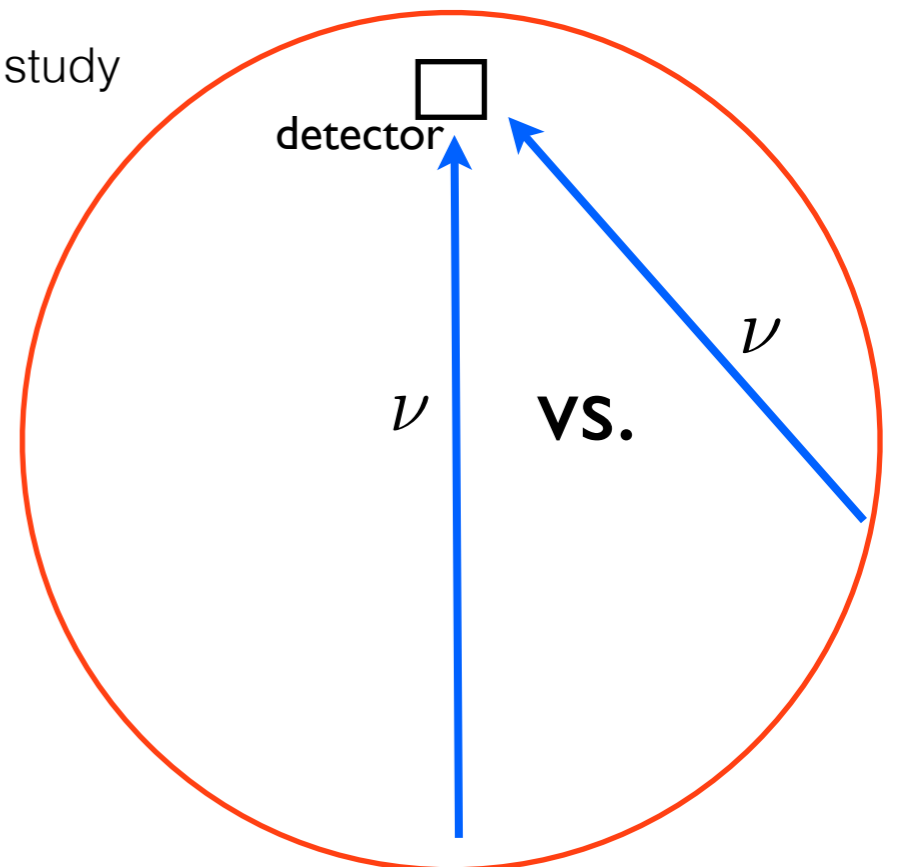
$$\left[ \begin{array}{l} \text{Non-standard int.} \\ \epsilon_{\tau\tau} = \frac{1}{6} \left( \frac{V_{\tau N}}{V_{nc}} - 1 \right) = \frac{\sin^2 \theta_\tau}{6} \end{array} \right]$$

see de Gouvea for DUNE study

Oscillation pattern depends on amount of matter traversed

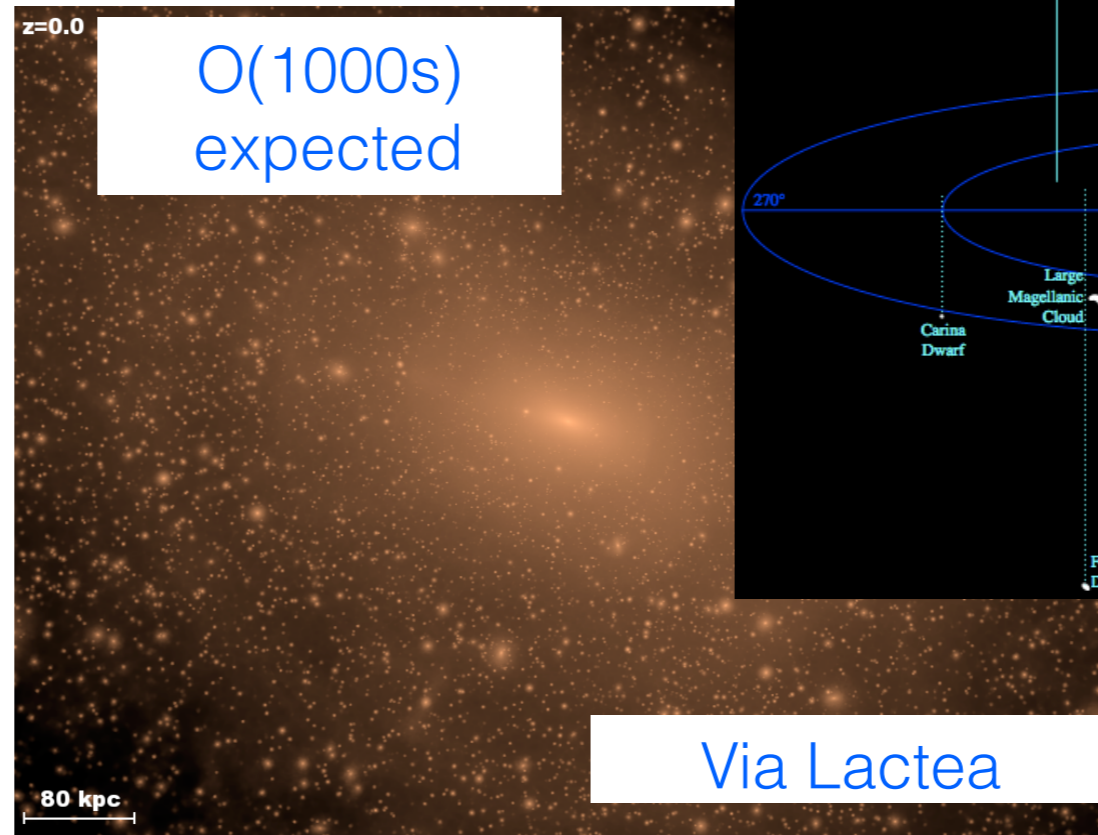
Super-K, arXiv:1410.2008

$\sin \theta_\tau < 0.42$  (stat. limited!)

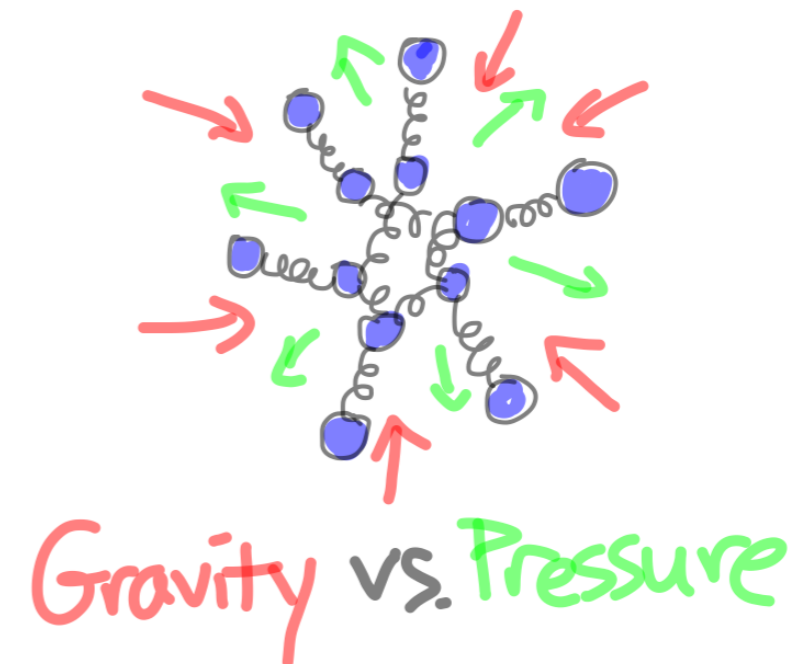
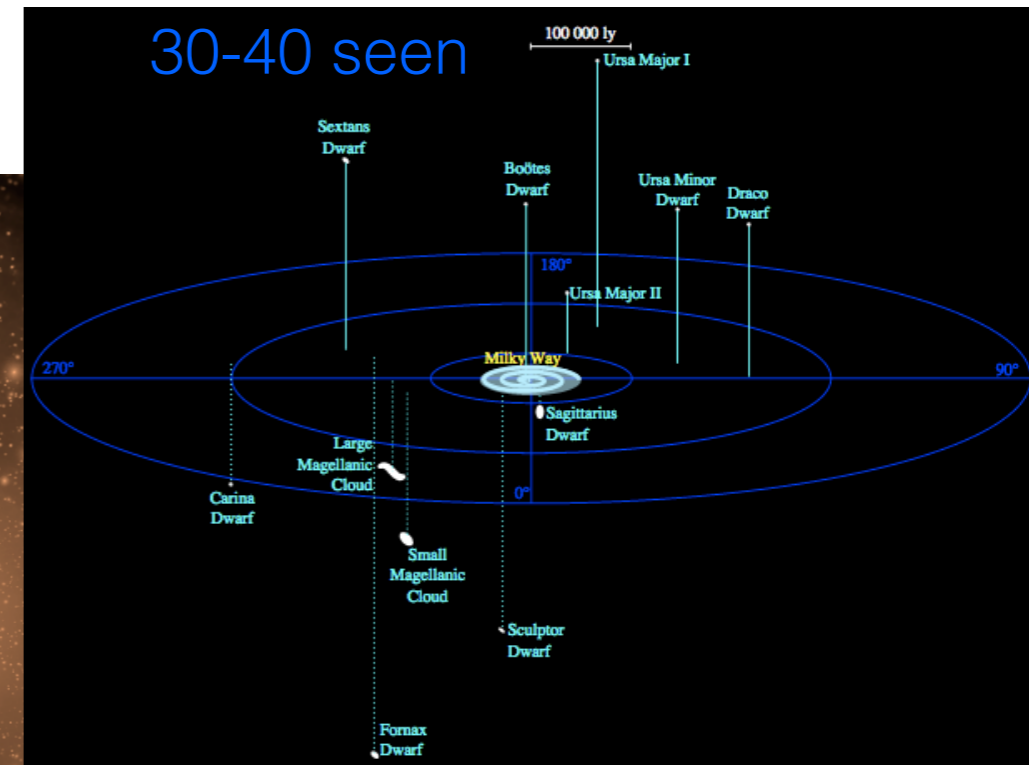
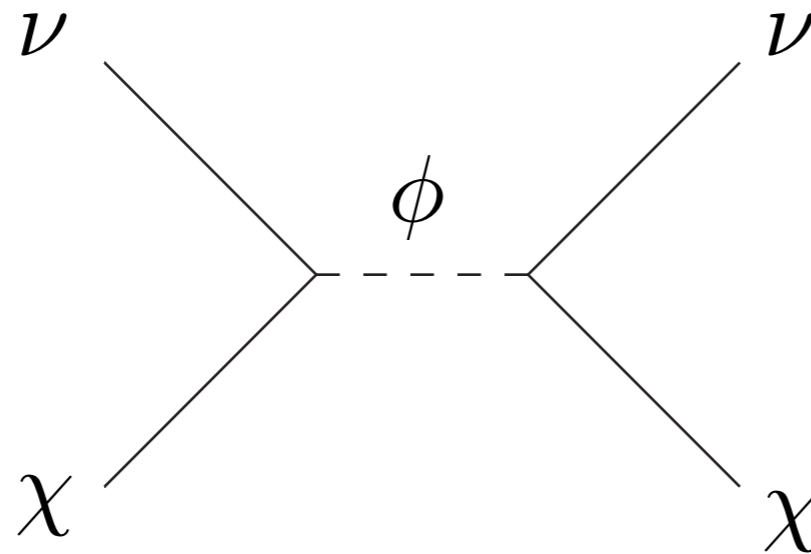


Are there hints for DM-neutrino interactions?

Count satellites of Milky Way galaxy:



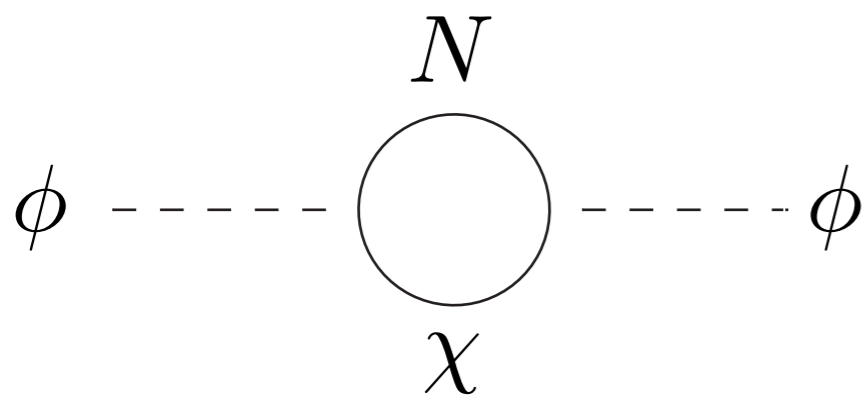
DM-scattering on neutrinos



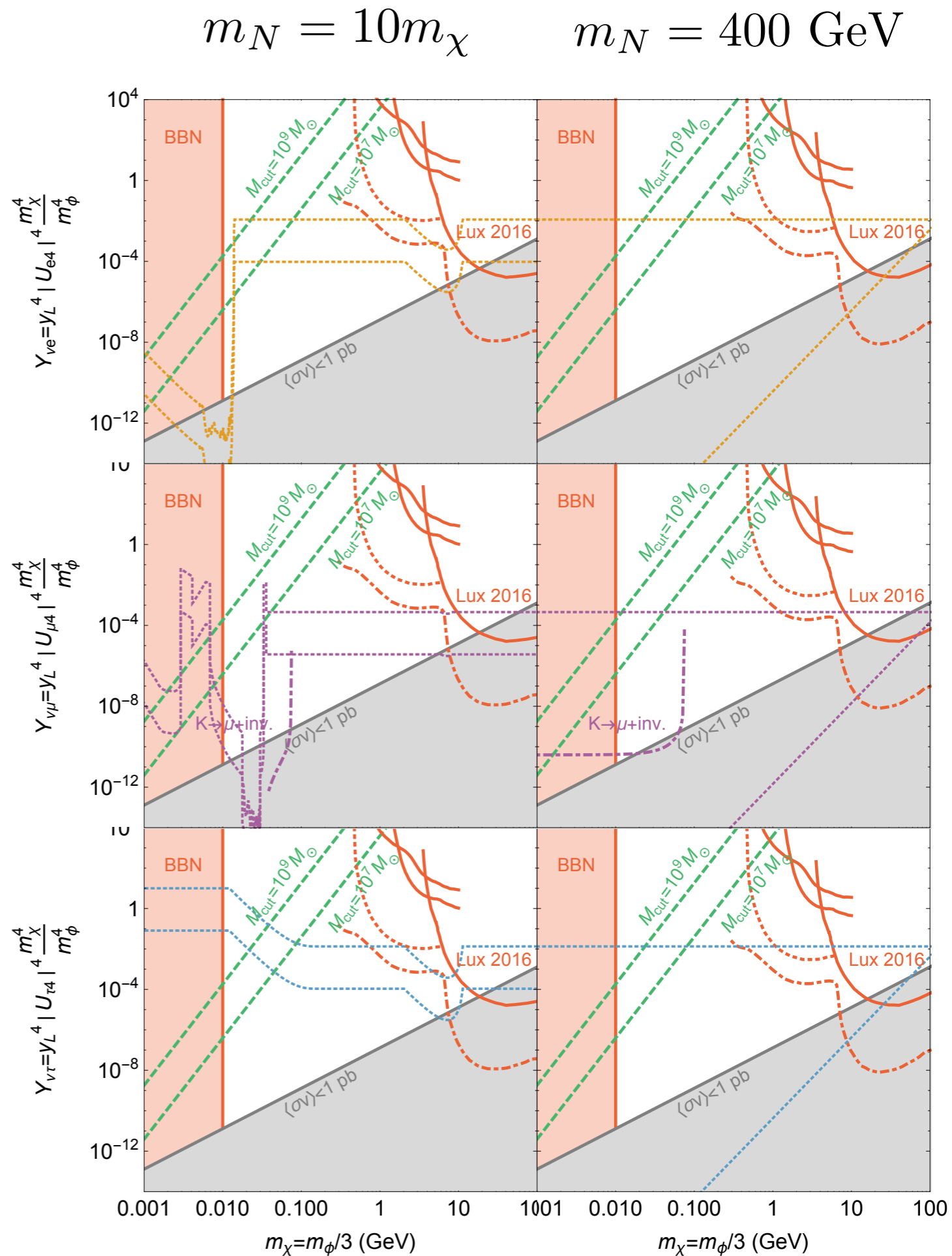
- Boehm et al.
- van den Aarssen et al.
- Shoemaker et al.
- Bertoni, Ipek, DM, & Nelson
- Hooper, Kaplinghat, Strigari, & Zurek

Can summarize using particular combination of couplings:

$$Y \equiv y_L^4 \left( \sum_i |U_{i4}|^2 \right)^2 \frac{m_\chi^4}{m_\phi^4} = 32\pi m_\chi^2 \langle \sigma v \rangle$$



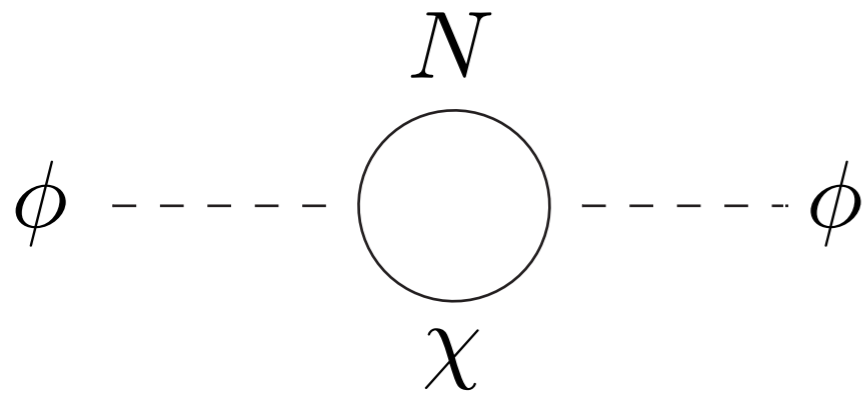
$$\delta m_\phi^2 \sim \frac{y^2}{16\pi^2} m_N^2$$



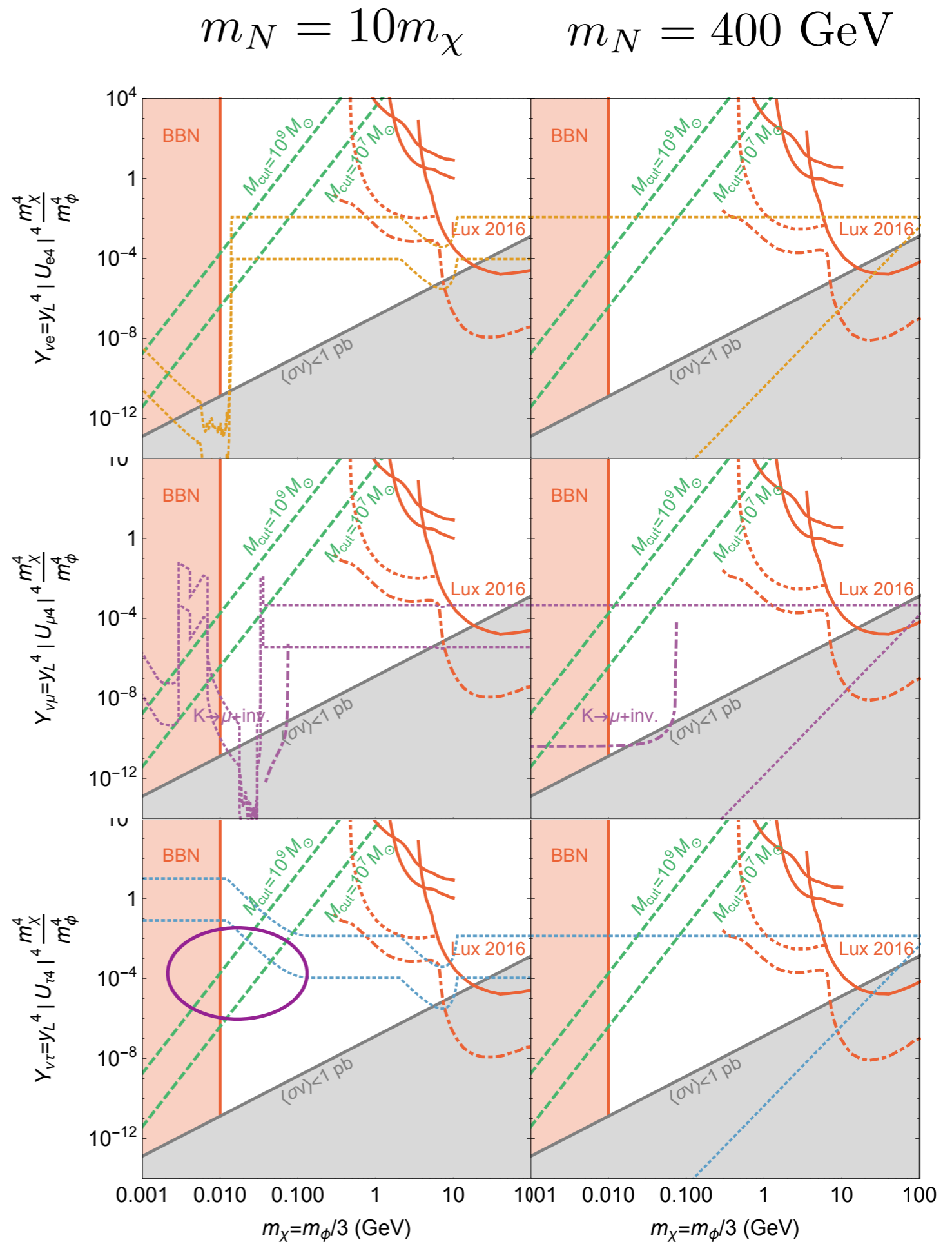


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$$\delta m_\phi^2 \sim \frac{y^2}{16\pi^2} m_N^2$$



# Neutrino Oscillations when large tau mixing

Assume mixing is dominantly with  $\tau$ , just 1 more mixing angle in addition to the usual 3, and just 1 more (large) mass splitting

$$U = \begin{pmatrix} U_{e1}^{3 \times 3} & U_{e2}^{3 \times 3} & U_{e3}^{3 \times 3} & 0 \\ U_{\mu 1}^{3 \times 3} & U_{\mu 2}^{3 \times 3} & U_{\mu 3}^{3 \times 3} & 0 \\ c_{\theta} U_{\tau 1}^{3 \times 3} & c_{\theta} U_{\tau 2}^{3 \times 3} & c_{\theta} U_{\tau 3}^{3 \times 3} & s_{\theta} \\ -s_{\theta} U_{\tau 1}^{3 \times 3} & -s_{\theta} U_{\tau 2}^{3 \times 3} & -s_{\theta} U_{\tau 3}^{3 \times 3} & c_{\theta} \end{pmatrix}$$

$$|U_{e2}|^2, |U_{\mu 2}|^2 + |U_{\tau 2}|^2 \quad \text{solar neutrinos}$$

$\Rightarrow$  Solar neutrinos  
potentially  
sensitive

$$|U_{e1}|^2 |U_{e2}|^2 \quad \text{KamLAND}$$

$$|U_{\mu 3}|^2 (1 - |U_{\mu 3}|^2) \quad \text{atmospheric/accelerator}$$

Uncertainty on  
flux ( $^8\text{B}$ )  $\sim 15\%$

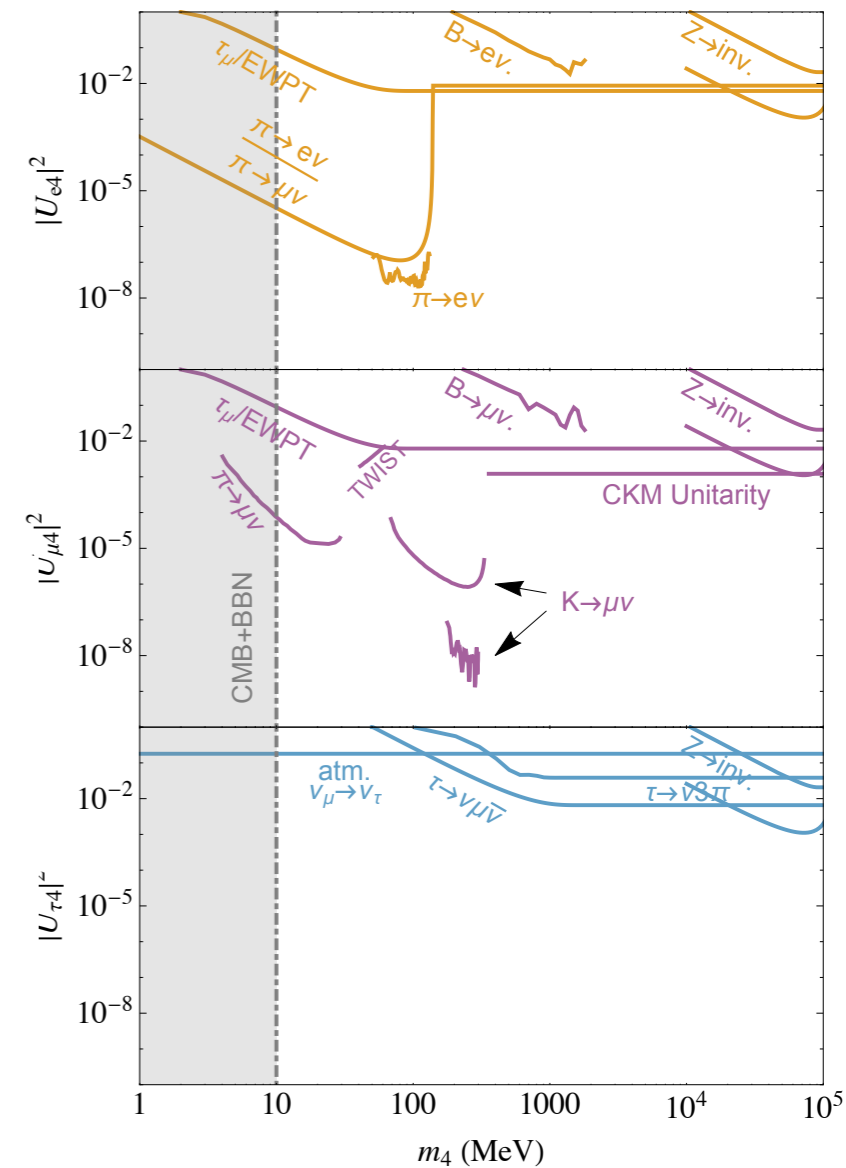
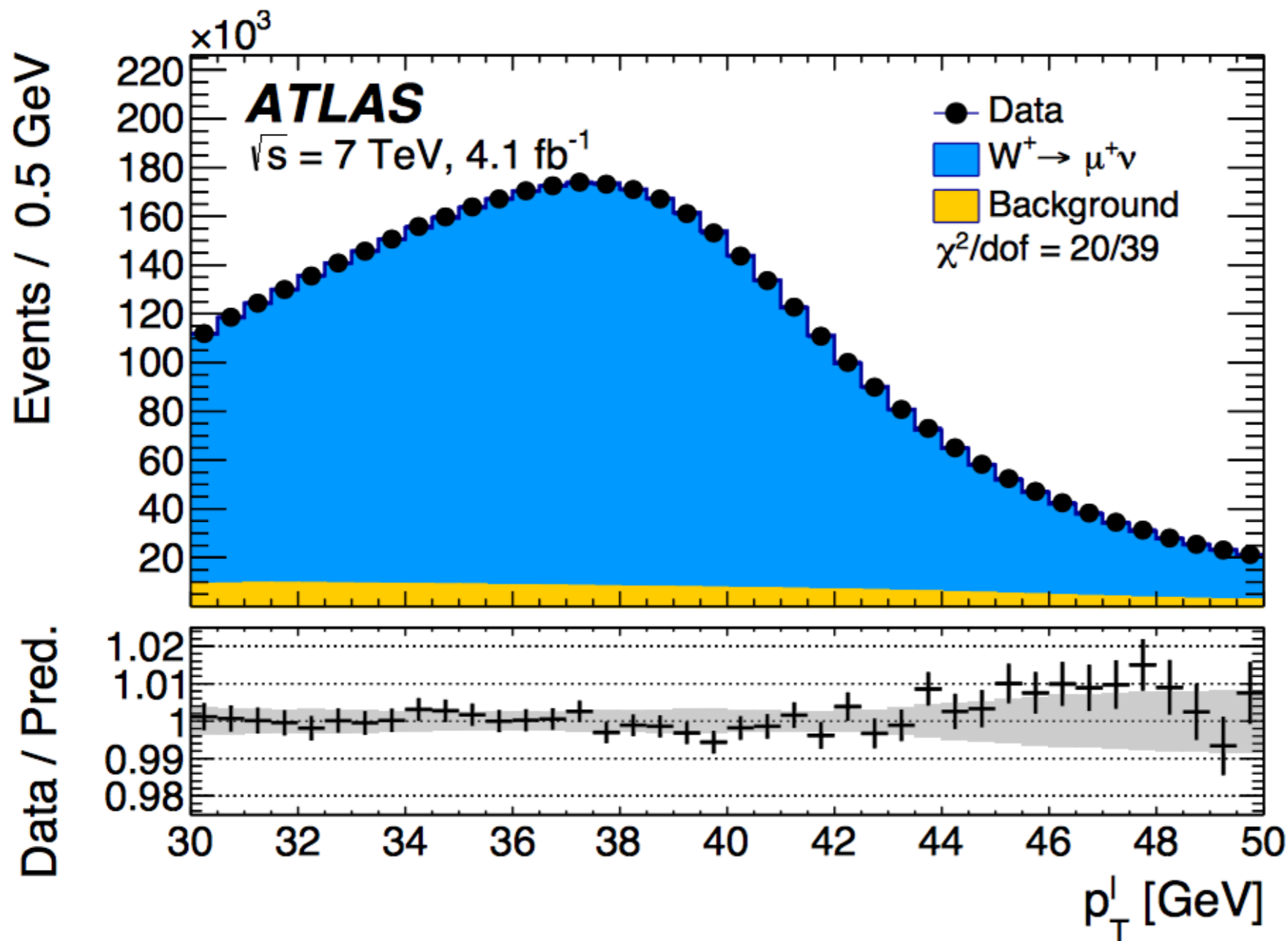
$$|U_{e3}|^2 (1 - |U_{e3}|^2) \quad \text{short baseline reactors}$$

$$\sin \theta_{\tau} < 0.6$$

$$|U_{e3}|^2 |U_{\mu 3}|^2 \quad \text{long baseline accelerator}$$

In progress...

# ATLAS W mass measurement 1701.07240

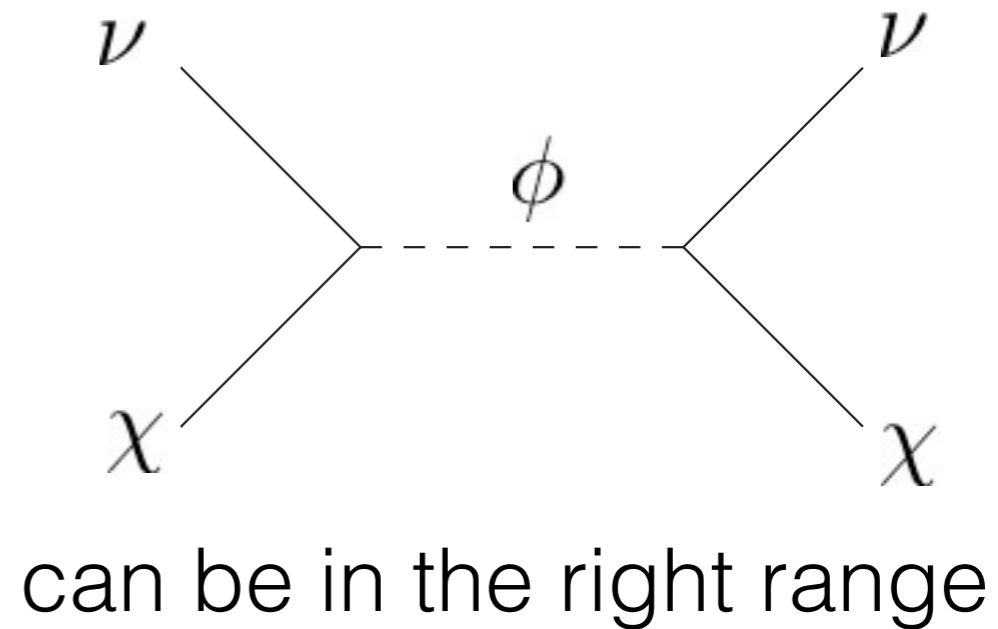


$\updownarrow \sim 2U^2$

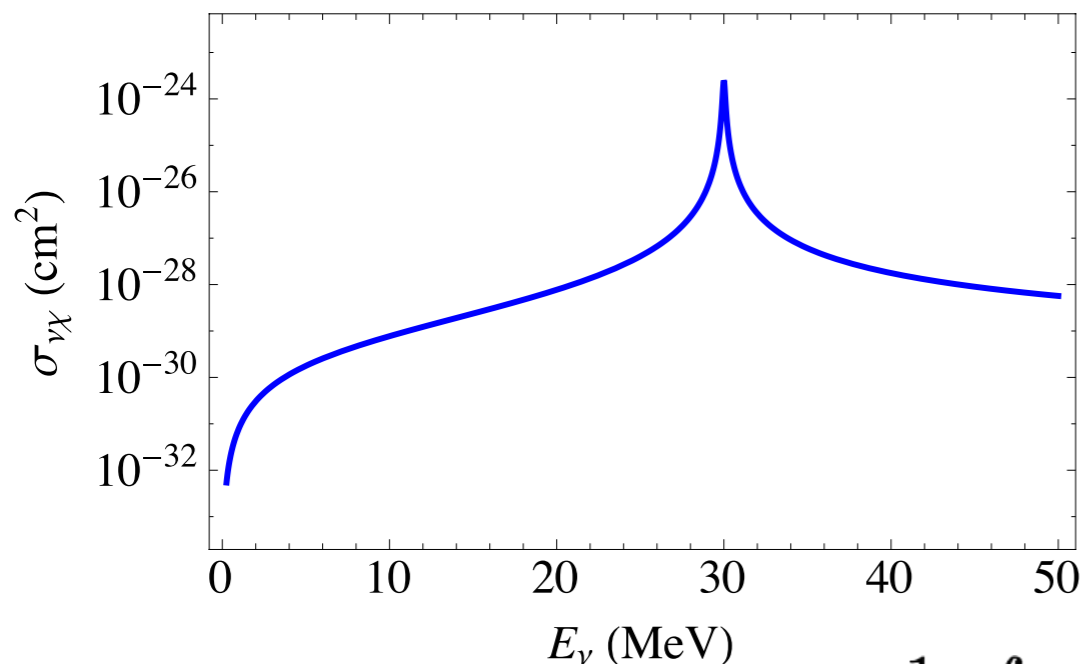
# Neutrinos from Supernovae

MeV energy neutrinos  
from SN scatter on DM

Resonance at  $E_\nu = \frac{m_\phi^2 - m_\chi^2}{2m_\chi}$

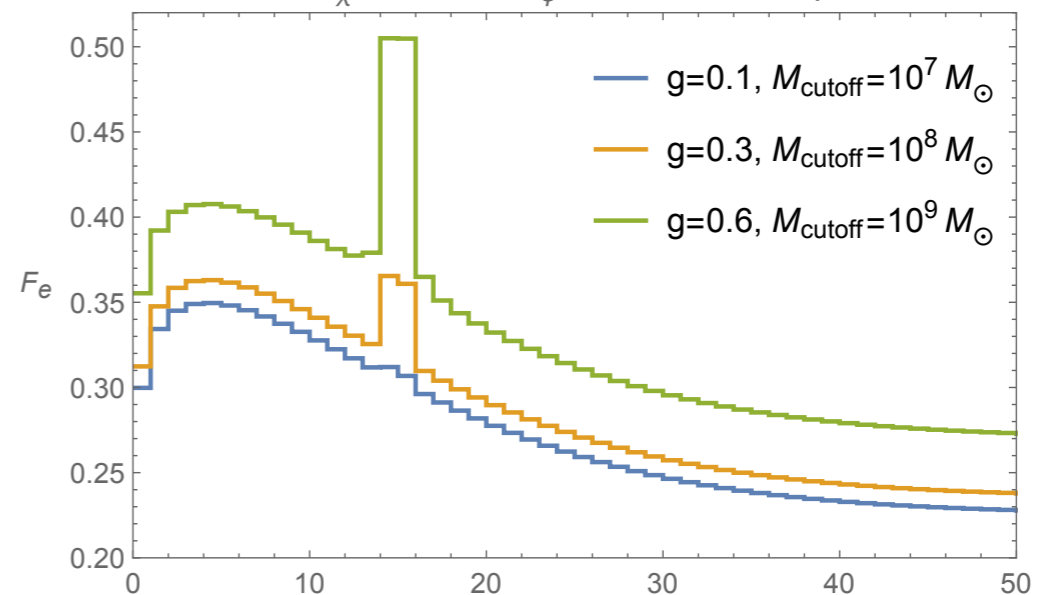


can be in the right range



$$\text{Flux}_i \propto e^{-\Gamma_i d} \quad \Gamma = \sigma_{\nu\chi} \times \frac{1}{d} \int dx n_\chi$$

Electron neutrino fraction (SN1987A)  
 $m_\chi=10$  MeV,  $m_\phi=20$  MeV,  $l=51$  Kpc



$$\frac{1}{\Gamma_1} \simeq \frac{6}{\Gamma}, \quad \frac{1}{\Gamma_2} \simeq \frac{3}{\Gamma}, \quad \frac{1}{\Gamma_3} \simeq \frac{2}{\Gamma}$$

# Supernovae Limits

Neutrinos produced in SN at  $T \sim 30$  MeV

Initial neutronization burst of  $\nu_e$  followed by cooling

DM light enough to be produced but doesn't contribute to cooling, thermal dist. with neutrinos to large radii

Neutrinos free stream when density is low,  $T \sim 5$  MeV: DM production suppressed, similar to strong  $\nu$  self-interactions

Fayet, Hooper, & Sigl, hep-ph/0602169 find  $m_\chi > 10$  MeV

Mangano et al., hep-ph/0606190 & Boehm et al., 1303.6270:

$$\sigma_{\hat{\nu}_i \chi} \lesssim 10^{-25} \text{ cm}^2 \left( \frac{m_\chi}{\text{MeV}} \right) \quad 33$$

# Supernovae Limits

Large fraction of DM gravitationally bound:  $v_{\text{esc}} \sim 0.5 c$

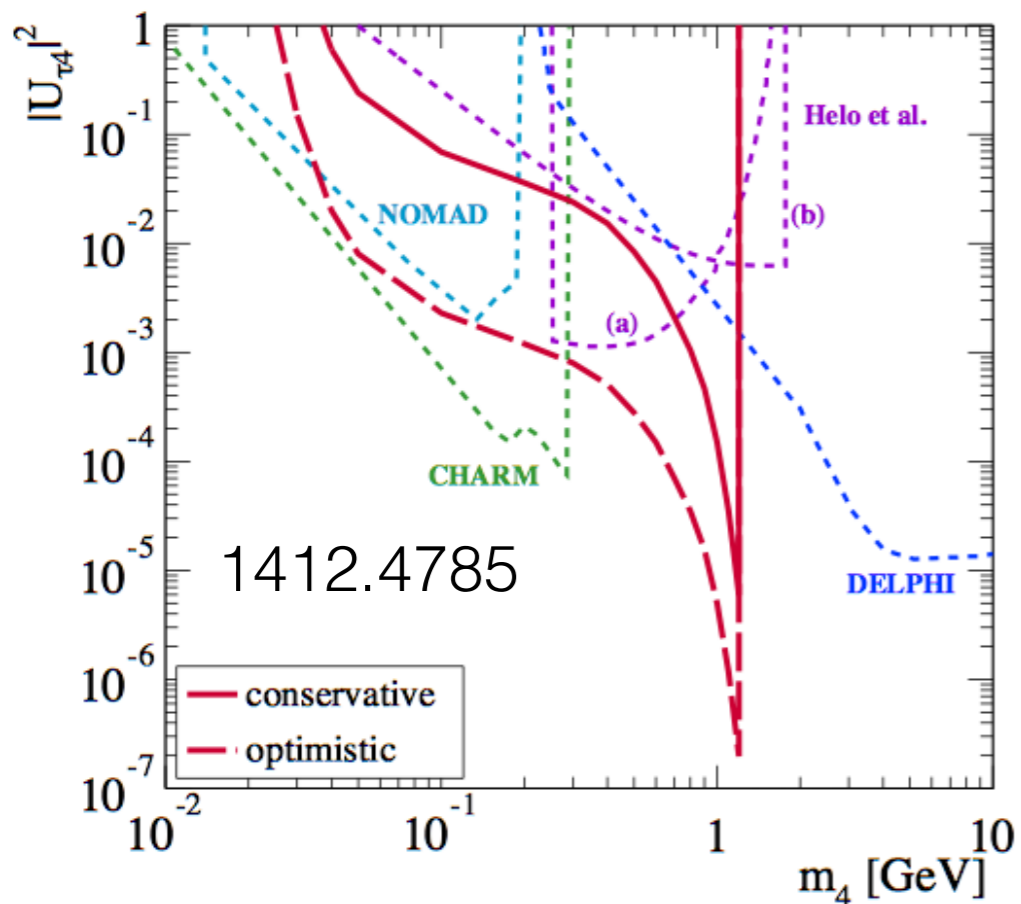
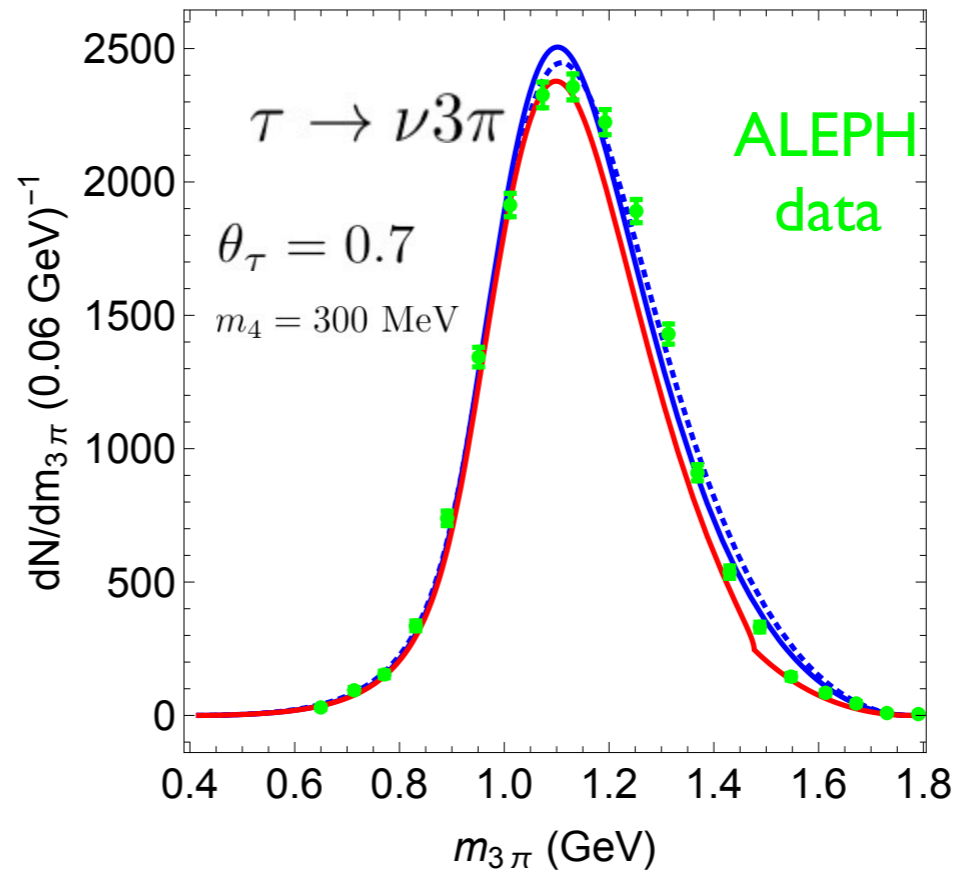
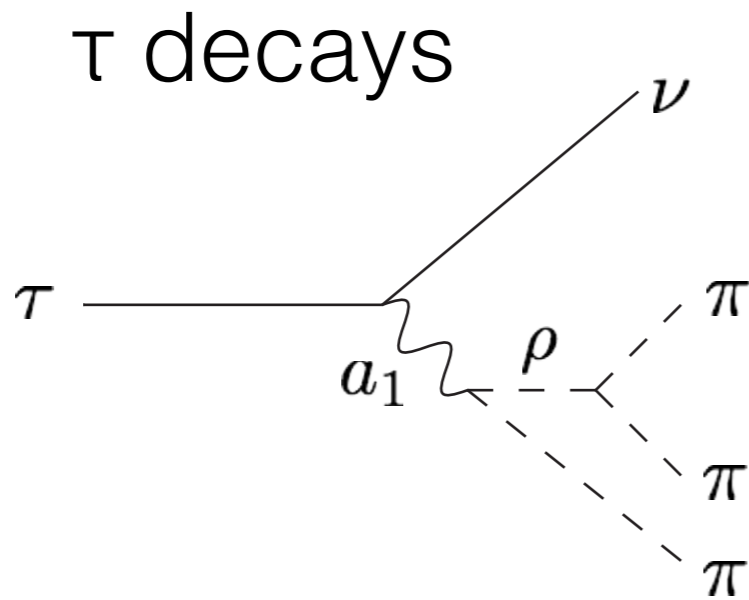
Is location (temperature) of  $v$ -sphere changed?

What are effects of flavor?

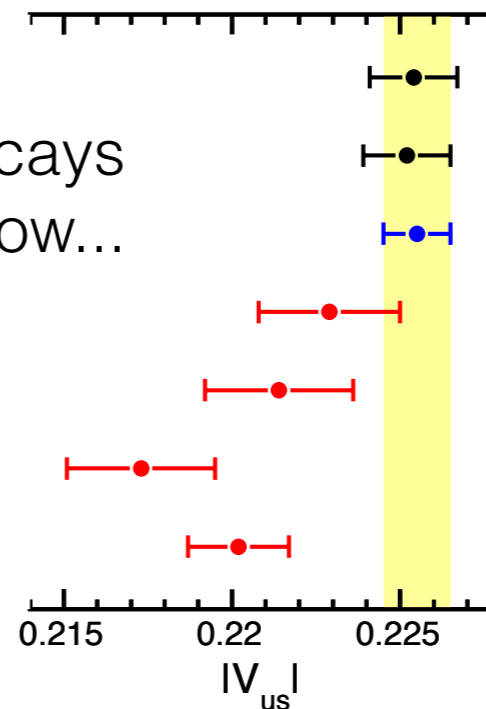
Could  $v$  “dwell” time be increased?

Very complicated...

# Future tests

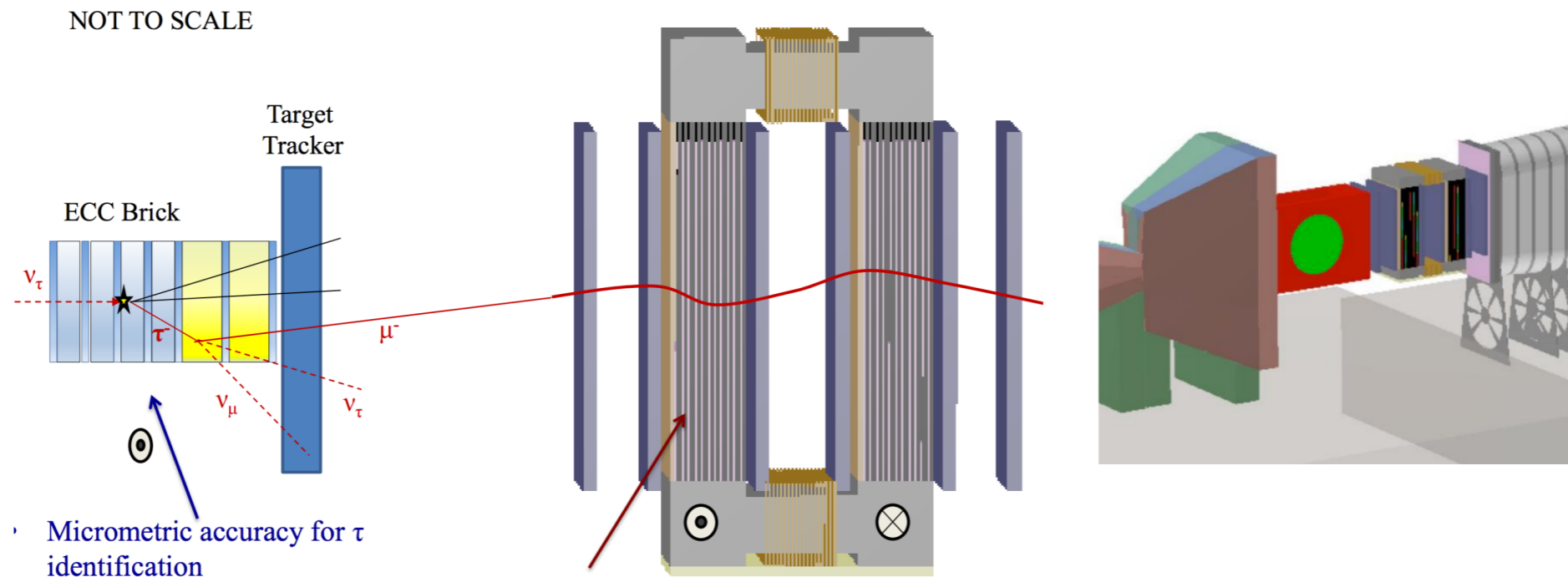


$\tau \rightarrow K$  decays  
 slightly low...



- $K_{13}$  decays, FlaviaNet 2010  
 $0.2254 \pm 0.0013$
- $K_{12}$  decays, FlaviaNet 2010  
 $0.2252 \pm 0.0013$
- CKM unitarity  
 $0.2255 \pm 0.0010$
- $\tau \rightarrow K\nu / \tau \rightarrow \pi\nu$ , HFAG 2012  
 $0.2229 \pm 0.0021$
- $\tau \rightarrow K\nu$ , HFAG 2012  
 $0.2214 \pm 0.0022$
- $\tau \rightarrow s$  inclusive, HFAG 2012  
 $0.2173 \pm 0.0022$
- $\tau$  average, HFAG 2012  
 $0.2202 \pm 0.0015$

Can an  $O(3-4k)$   $v_\tau$  sample at SHiP impact a scenario like this?



(see talk by N. Serra)



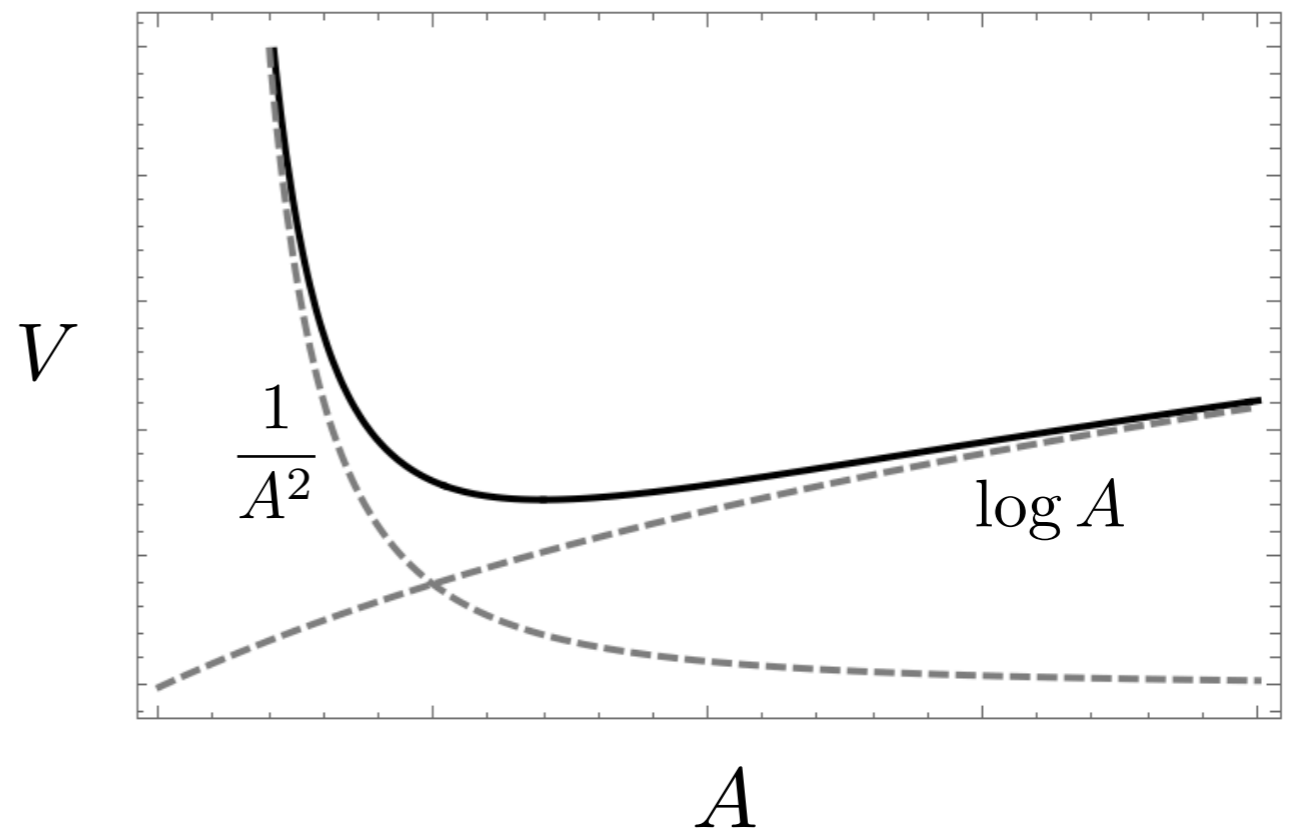
# Sterile neutrino portal to a light scalar

Consider  $\mathcal{L}_{\text{mass}} = -m_D \nu N - m_N N N + \text{h.c.}$

with  $m_N(A) = m_0 + \kappa A$   $V_0 = \Lambda^4 \log \left( 1 + \left| \frac{A}{\sigma} \right| \right)$

$$V(A, T) = \Lambda^4 \log \left( 1 + \left| \frac{A}{\sigma} \right| \right) + \frac{m^2(A) T^2}{24} \quad m \propto \frac{1}{A}$$

Temperature-  
dependent potential



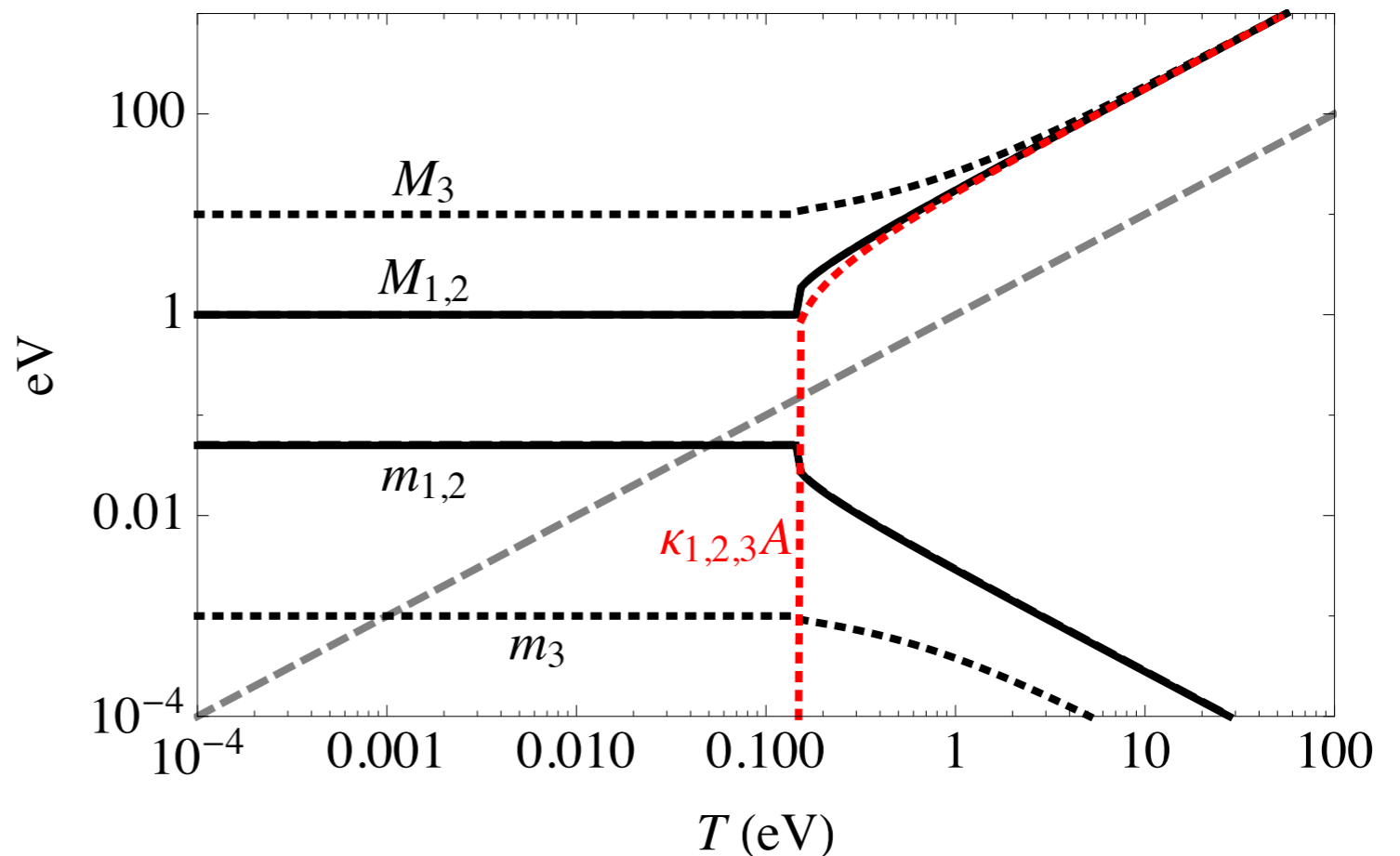
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$$V(A, T) = \Lambda^4 \log \left( 1 + \left| \frac{A}{\sigma} \right| \right) + \frac{m^2(A) T^2}{24}$$

Temperature-  
dependent masses



# Sterile neutrino portal to a light scalar

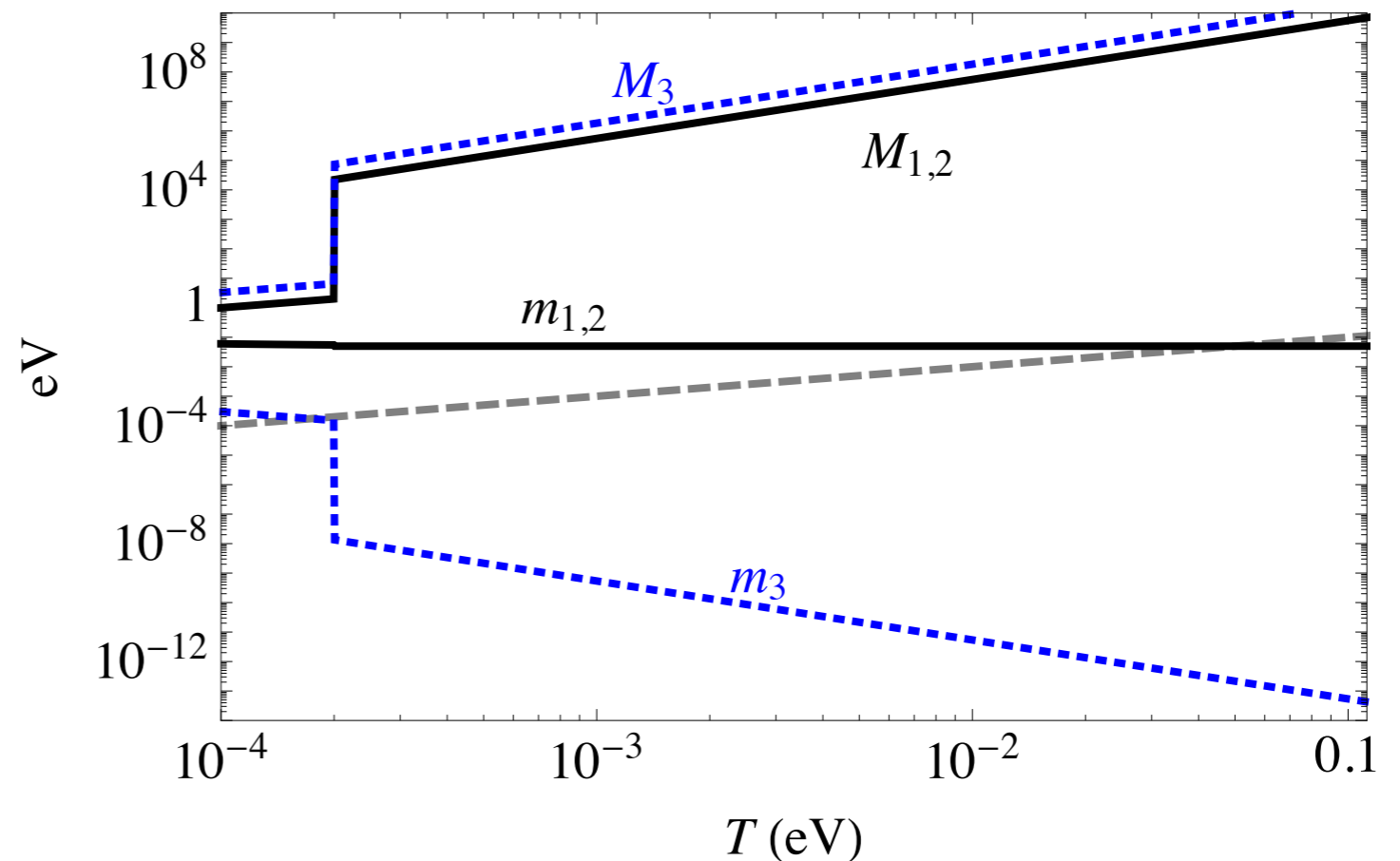
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$$V(A, T) = \Lambda^4 \log \left( 1 + \left| \frac{A}{\sigma} \right| \right) + \frac{m^2(A) T^2}{24}$$

Add small active Majorana mass for dark energy

$$\mathcal{L} \supset -\mu \nu \nu$$



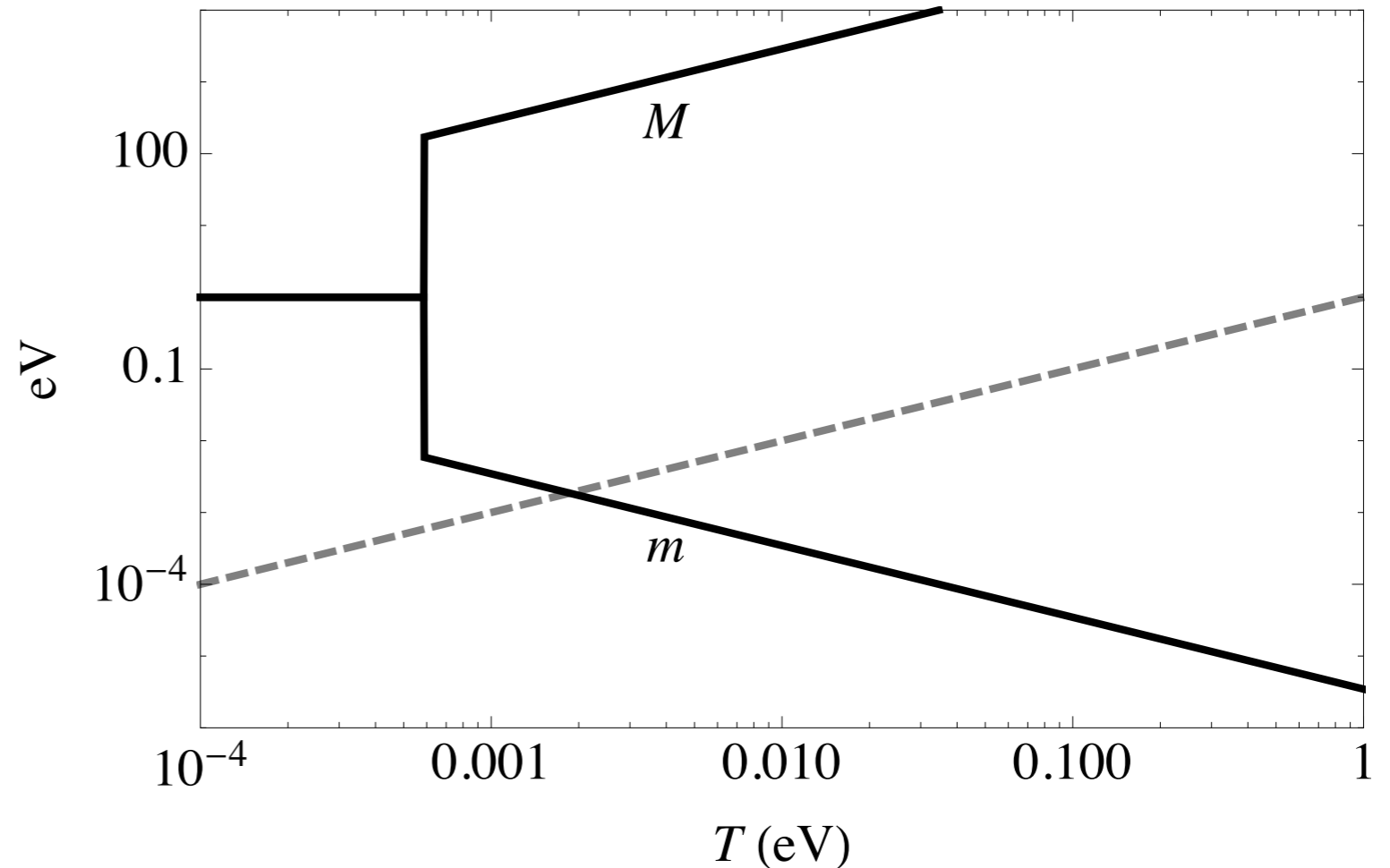
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$$V(A, T) = \Lambda^4 \log \left( 1 + \left| \frac{A}{\sigma} \right| \right) + \frac{m^2(A) T^2}{24}$$

Or get active  
neutrinos at an eV  
  
~massless till late  
times, then  
“nuggets” form



# Wrap up

Neutrino portal is a viable, less well studied way to couple to dark sector

Leads to a rich phenomenology

Can help with some problems in dark matter

Interesting new probes—lots of connections across fields!