

EDMs & P-Conserving T-Violation



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Outline

- I. Intro & Motivation*
- II. TVPC Interactions: Background*
- III. EFT Illustration: m_ν & μ_ν*
- IV. TVPC Interactions & EDMS*
- V. Summary*

Questions for Theory

C and P Symmetries (assuming CPT)

D. Mack, MENU

C	<p>C, P, CP Strong, EM</p> <p>Big SM "background" in any search for new forces</p>	<p>C, P, CP Weak (loop-level)</p> <p>Small SM "background". New sources of P, CP constrained by EDM searches</p>
G	<p>G, P, CP Weak (loop-level)</p> <p>Small SM "background". New sources of P, CP less constrained by EDM searches</p>	<p>G, P, CP Weak</p> <p>Big SM "background" in any search for new forces</p> <p><i>New sources of PV also constrained by amplitude- sensitive PV asymmetry measurements</i></p>

C Violation Basics

D. Mack, MENU

The charge conjugation operator C reverses all generalized charges, effectively replacing a particle by its anti-particle.

C violation is known only in

- 1. Weak interactions at tree level which violate P (hence conserving CP)*
- 2. Weak interactions at loop level which violate CP*

Both C - and CP -violation are among the Sakharov criteria for baryogenesis.

Everybody knows strong and EM forces conserve C but direct bounds on C violation in these amplitudes are only $\sim 0.5\%$. How to improve this?

It is surprisingly hard:

- Only a few neutral particles are states of good C and thus suitable for tests (γ , π^0 , η , J/ψ , or a self-conjugate system like e^+e^-).*
- Most of the particles of good C appropriate for initial states aren't easy to make in large quantities (and with sufficiently low backgrounds).*

η Decays Testing C Violation

Why η 's?

- The η full width is only 1.3 keV. It cannot decay by the isospin conserving strong interaction. This means that achievable BR's of 10^{-6} to 10^{-7} probe the weak scale.
- η decays are flavor-conserving, a sector less thoroughly studied than $\Delta S = 1$, etc.
- Theory calculations predict large mass enhancements, hence relatively crude η decay BR upper limits place tighter constraints than more precise π^0 decay BR upper limits.
- The η has a significant s -bar content, unlike the π^0 or nucleon.

→	3γ	$< 1.6 \cdot 10^{-5}$	3
	" $\pi^0\gamma$ "	$< 9 \cdot 10^{-5}$	
→	$2\pi^0\gamma$	$< 5 \cdot 10^{-4}$	5
	$3\gamma\pi^0$	Nothing published	
	$3\pi^0\gamma$	$< 6 \cdot 10^{-5}$	7
	$3\gamma 2\pi^0$	Nothing published	

PDG
2012

Considerations of acceptance and phase space have focused us on $\eta \rightarrow 3\gamma$ and $\eta \rightarrow 2\pi^0\gamma$.

Most C test channels are all-neutral except for $\eta \rightarrow \pi^0\gamma^* \rightarrow \pi^0\ell\ell$.

D. Mack, MENU

Theory Issues for C Violation

Placing the tightest direct limits on C violation sounds interesting to experimentalists, but what about theorists?

- *Little literature on C violation with P conservation.*

(appropriate models for this would be non-renormalizable - Herczeg)

- *Some literature on T violation with P conservation*

(under CPT, equivalent to C violation with P conservation).

- *By contrast, tremendous literature on CP violation and EDM's.*

- *C violation without P violation is apparently not on the radar of those working with SUSY, leptoquarks.*

- *C violation does arise in discussions of violation of Lorentz invariance, but the predicted C violating η decay BR's are effectively zero for any experiment, ever.*

We'd like theorists studying T violation with P conservation to know that η decays can place tight limits in an isospin-violating sector.

TVPC Interactions: Background

TVPC Interactions

- *Herczeg: No renormalizable TVPC boson-exchange interactions involving only SM fields [Hyperfine Int, 75 (1992) 127]*
- *Low-energy ($k \ll \Lambda_{EW}$) four fermion interactions first arise at $d=7$:*

$$\mathcal{O}_7^{ff'} = C_7^{ff'} \bar{\psi}_f \overleftrightarrow{D}_\mu \gamma_5 \psi_f \bar{\psi}_{f'} \gamma^\mu \gamma_5 \psi_{f'}$$

*Khriplovich '91
Conti & Khriplovich '92
Engel, Frampton, Springer '96*

TVPC Interactions, cont'd

- *Additional low-energy ($k \ll \Lambda_{EW}$) $d=7$ interactions:*

$$\mathcal{O}_7^{\gamma g} = C_7^{\gamma g} \bar{\psi} \sigma_{\mu\nu} \lambda^a \psi F^{\mu\lambda} G_\lambda^{a\nu}$$

$$\mathcal{O}_7^{\gamma Z} = C_7^{\gamma Z} \bar{\psi} \sigma_{\mu\nu} \psi F^{\mu\lambda} Z_\lambda^\nu$$

+ ...

MR-M '99

Kurylov, McLaughlin, MR-M '01

TPVC Observables

- “D Coefficient” in β -decay:

$$\frac{d^3\Gamma}{dE_e d\Omega_e d\Omega_\nu} = \frac{1}{(2\pi)^5} p_e E_e (E_0 - E_e)^2 \xi \left\{ 1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \langle \frac{\vec{J}}{J} \rangle \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] + \dots \right\}$$

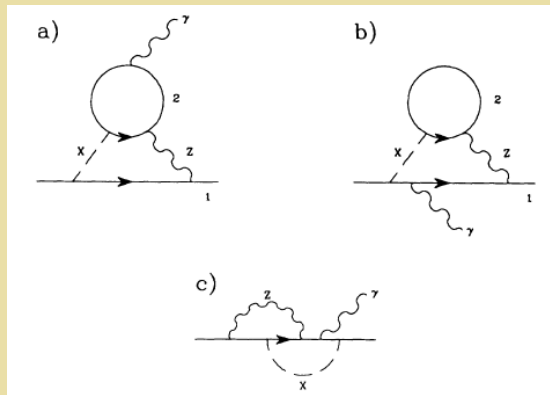
- Correlations in $\vec{n} + A$ scattering:

$$\vec{\sigma}_n \cdot \vec{k}_n \times \vec{J} \left(\vec{k}_n \cdot \vec{J} \right)$$

- $\eta \rightarrow 3\gamma, \eta \rightarrow 2\pi^0\gamma, \dots$:

TPVC Interactions & EDMs

- *Conti & Khriplovich: TVPC interactions + SM radiative corrections (PV) induce non-vanishing EDMs*



$$\mathcal{O}_5 = -\frac{i}{2} C_5^f \bar{\psi} \sigma_{\mu\nu} \gamma_5 \psi F^{\mu\nu}$$

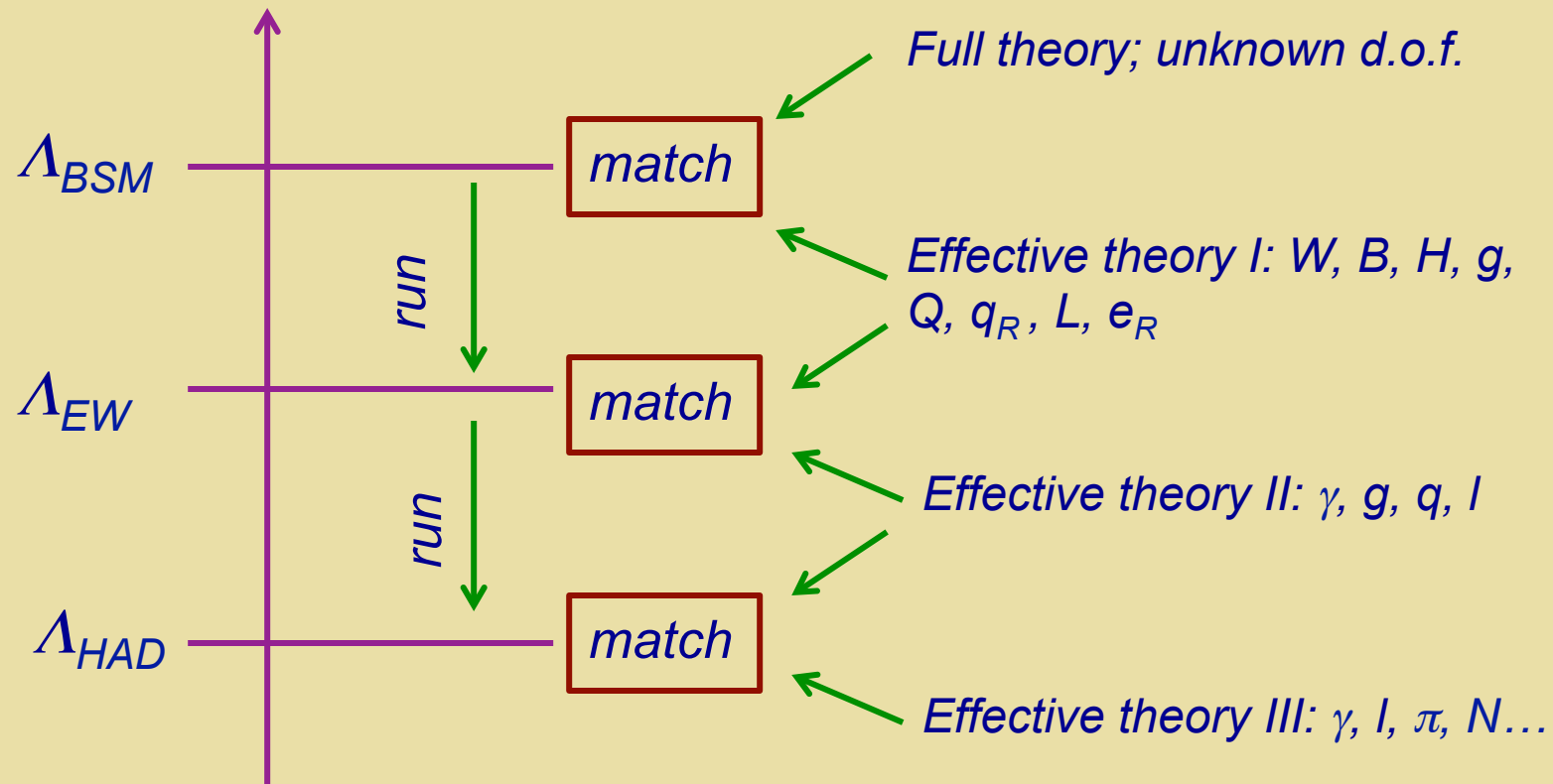
- *EDM limits imply vanishingly small effects from TVPC interactions*

$$C_7^{ff'} \lesssim \left(\frac{\Lambda_{TVPC}}{1 \text{ TeV}} \right) \times 10^{-3}$$

How Robust Is Bound?

- *Non-renormalizable interactions: EFT, running, matching & “naturalness”*
- *Illustration with neutrino magnetic moments*
- *Application to TVPC interactions*

Non-Renormalizable Interactions & EFT



Effective Theory I

$$\mathcal{L}_{\text{CPV}} = \mathcal{L}_{\text{CKM}} + \mathcal{L}_{\bar{\theta}} + \mathcal{L}_{\text{BSM}}^{\text{eff}}$$

$$\mathcal{L}_{\text{BSM}}^{\text{eff}} = \frac{1}{\Lambda^2} \sum_i \alpha_i^{(n)} O_i^{(6)} + \dots$$

*Effective theory I: W, B, H, g,
Q, q_R, L, e_R*

Effective Theory II

$$\mathcal{L}_{\text{new}} = \mathcal{L}_4 + \frac{1}{\Lambda_{\text{TVPC}}} \mathcal{L}_5 + \frac{1}{\Lambda_{\text{TVPC}}^2} \mathcal{L}_6$$
$$+ \frac{1}{\Lambda_{\text{TVPC}}^3} \mathcal{L}_7 + \dots,$$

+...

Effective theory II: γ, g, q, l

Effective Theory II

$$\mathcal{L}_{\text{new}} = \mathcal{L}_4 + \frac{1}{\Lambda_{\text{TVPC}}} \mathcal{L}_5 + \frac{1}{\Lambda_{\text{TVPC}}^2} \mathcal{L}_6$$
$$+ \frac{1}{\Lambda_{\text{TVPC}}^3} \mathcal{L}_7 + \dots,$$

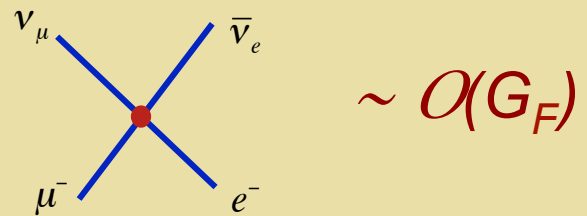
+ ...

Effective theory II: γ, g, q, l

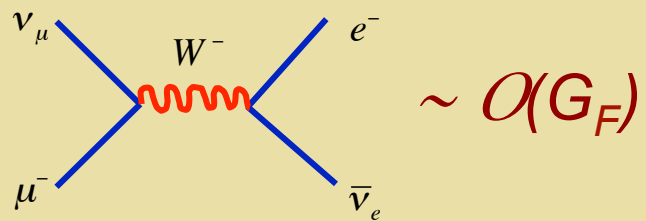
Matching I & II: compute in II with massive W,Z

Matching

Fermi Effective Theory

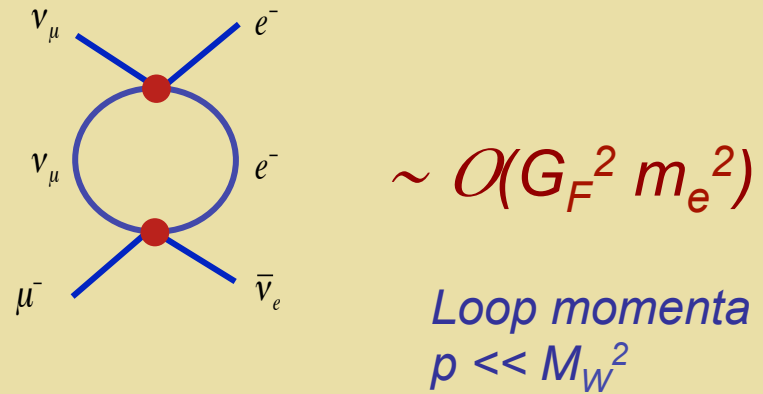
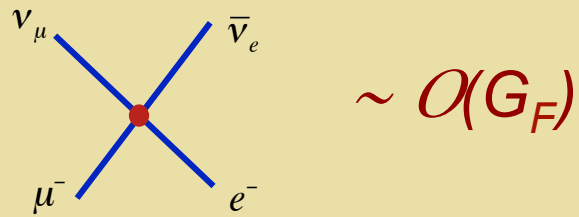


Standard Model

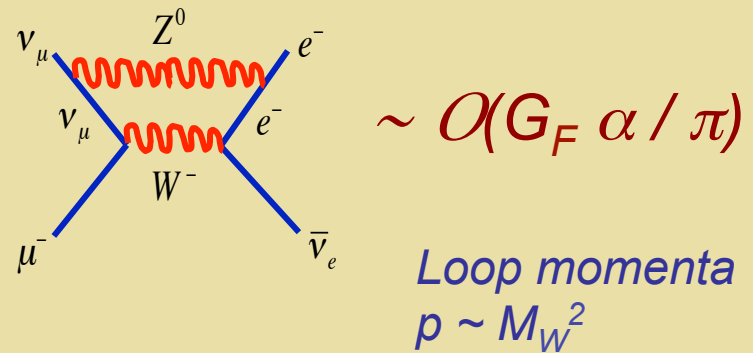
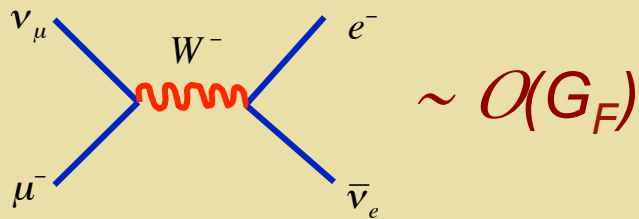


Matching

Fermi Effective Theory

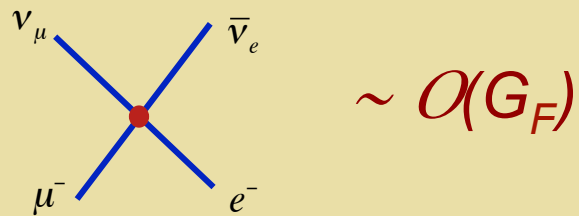


Standard Model

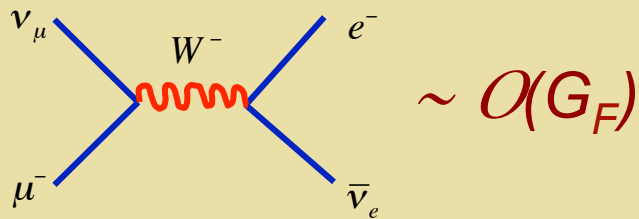


Matching

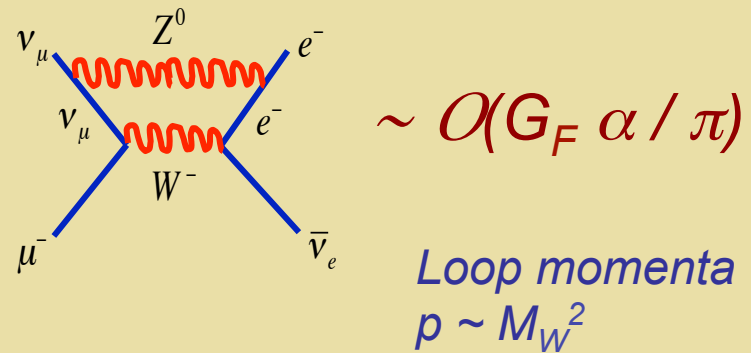
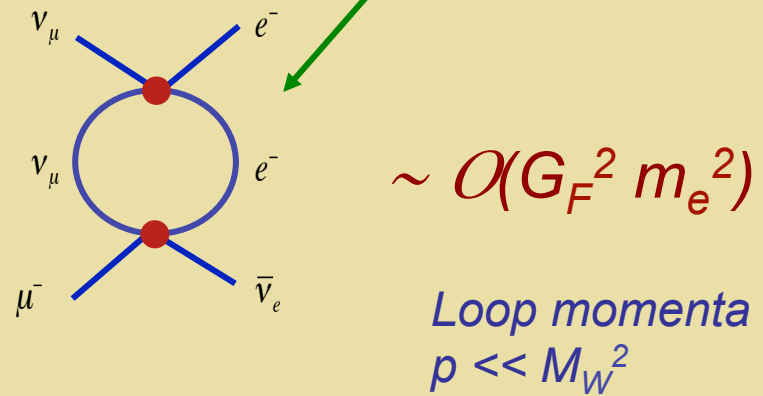
Fermi Effective Theory



Standard Model



Match EFT onto full theory
by considering $p \sim \Lambda$ (NDA)



EFT Illustration: m_ν & μ_ν

Bell, Cirigliano, R-M, Vogel, Wise '05

Also

Bell, Gorchtein, R-M, Vogel, Wang '06

Erwin, Kile, R-M, Wang '07

Kile, R-M '07

Neutrino Magnetic Moments

Can neutrinos have magnetic moments?

Dirac neutrinos

$$\nu_L \rightarrow e^{i\alpha_L} \nu_L \quad \nu_R \rightarrow e^{i\alpha_R} \nu_R$$

$$m_\nu \bar{\nu}\nu = m_\nu \{ \bar{\nu}_L \nu_R + \bar{\nu}_R \nu_L \}$$

\mathcal{L} not invariant if $m_\nu \neq 0$ and $\alpha_L \neq \alpha_R$

Magnetic moment operator forbidden

$$\mu_\nu \bar{\nu} \sigma_{\alpha\beta} \nu F^{\alpha\beta} = \mu_\nu \{ \bar{\nu}_L \sigma_{\alpha\beta} \nu_R + \bar{\nu}_R \sigma_{\alpha\beta} \nu_L \} F^{\alpha\beta}$$

The Scale of m_ν and μ_ν

Minimal extension of the Standard Model with ν_R and non-vanishing m_ν gives

$$\mu_\nu \approx 3 \times 10^{-19} \mu_B [m_\nu / 1 \text{ eV}]$$

Too small to be observed

What about new physics at scale $\Lambda > v$? NDA

$$\mu_\nu \sim eG/\Lambda \quad , \quad m_\nu \sim G\Lambda$$

$$m_\nu \sim \frac{\Lambda^2}{2m_e} \frac{\mu_\nu}{\mu_B} \sim \frac{\mu_\nu}{10^{-18} \mu_B} \left(\frac{\Lambda}{1 \text{ TeV}} \right)^2 \text{ eV}$$

Evading the NDA Estimates

NDA

$$m_\nu \sim \frac{\Lambda^2}{2m_e} \frac{\mu_\nu}{\mu_B} \sim \frac{\mu_\nu}{10^{-18}\mu_B} \left(\frac{\Lambda}{1 \text{ TeV}} \right)^2 \text{ eV}$$

The “Voloshin” mechanism: a loophole

SU(2)_ν symmetry:

(ν, ν^c) transf as doublet

m_ν term transforms as triplet forbidden

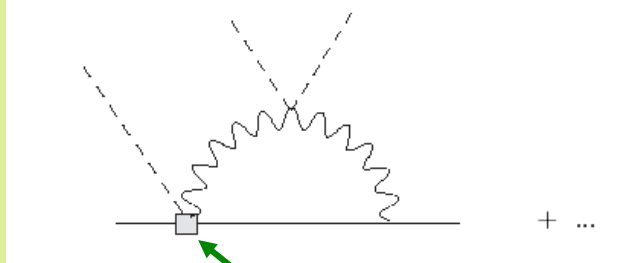
μ_ν transforms as singlet allowed

Evading the NDA Estimates

NDA

$$m_\nu \sim \frac{\Lambda^2}{2m_e} \frac{\mu_\nu}{\mu_B} \sim \frac{\mu_\nu}{10^{-18}\mu_B} \left(\frac{\Lambda}{1 \text{ TeV}} \right)^2 \text{ eV}$$

The “Voloshin” mechanism: a loophole



electroweak μ_ν operators

Radiatively-induced
neutrino mass

Voloshin sym & generalizations broken by SM
gauge & Yukawa interactions: m_ν bounds on μ_ν

Dirac Neutrinos

$$\delta m_\nu = -C_3^6(\nu) \frac{\nu^3}{2\sqrt{2}\Lambda^2} \quad \frac{\mu_\nu}{\mu_B} = -4\sqrt{2} \left(\frac{m_{e\nu}}{\Lambda^2} \right) C_+(\nu) \quad C_+ = C_1^6 + C_2^6$$

Operator Basis:

$$O_M^{(4)} = \bar{L}\tilde{\phi}\nu_R \quad \tilde{\phi} = i\tau_2\phi^*$$

$$O_1^{(6)} = g_1(\bar{L}\sigma^{\mu\nu}\tilde{\phi})\ell_R B_{\mu\nu}$$

$$O_2^{(6)} = g_2(\bar{L}\sigma^{\mu\nu}\tau^a\tilde{\phi})\nu_R W_{\mu\nu}^a$$

$$O_3^{(6)} = (\bar{L}\tilde{\phi}\nu_R) (\phi^\dagger\phi)$$

Dirac Neutrinos

$$\delta m_\nu = -C_3^6(\nu) \frac{\nu^3}{2\sqrt{2}\Lambda^2} \quad \frac{\mu_\nu}{\mu_B} = -4\sqrt{2} \left(\frac{m_{e\nu}}{\Lambda^2}\right) C_+(\nu) \quad C_+ = C_1^6 + C_2^6$$

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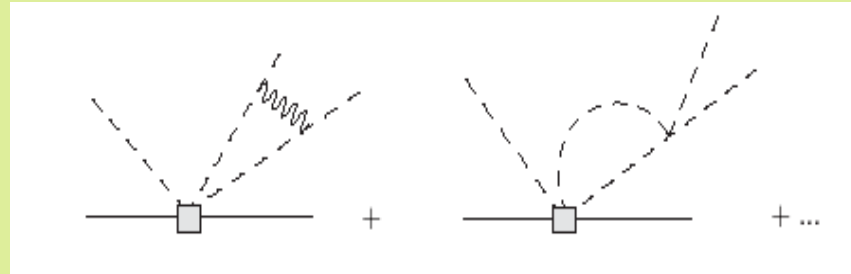
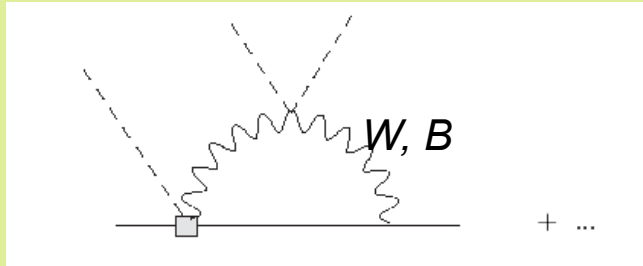
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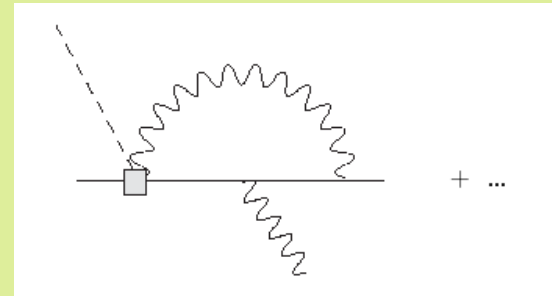
*Close under
renormalization*

Dirac Neutrinos: Mixing



Operator Basis:

$$O_M^{(4)} = \bar{L}\tilde{\phi}\nu_R \quad \tilde{\phi} = i\tau_2\phi^*$$



$$O_1^{(6)} = g_1(\bar{L}\sigma^{\mu\nu}\tilde{\phi})\ell_R B_{\mu\nu}$$

$$O_2^{(6)} = g_2(\bar{L}\sigma^{\mu\nu}\tau^a\tilde{\phi})\nu_R W_{\mu\nu}^a$$

$$O_3^{(6)} = (\bar{L}\tilde{\phi}\nu_R)(\phi^+\phi)$$

Close under renormalization

Dirac Neutrinos: Mixing & “Naturalness”

Renormalization Group: Leading Log

Solution with $C_3^6(\Lambda) = 0$: δm_ν generated entirely from radiative corrections

$$\frac{|\mu_\nu|}{\mu_B} = \frac{G_F m_e}{\sqrt{2}\pi\alpha} \left[\frac{\delta m_\nu}{\alpha \ln(\Lambda/\nu)} \right] \frac{32\pi \sin^4 \theta_W}{9|f|}$$

$$f = (1-r) - \frac{2}{3}r \tan^2 \theta_W - \frac{1}{3}(1+r) \tan^4 \theta_W \quad r = C_-/C_+$$

$$\frac{|\mu_\nu|}{\mu_B} \lesssim 8 \times 10^{-15} \times \left(\frac{\delta m_\nu}{1 \text{ eV}} \right) \frac{1}{|f|}$$

Dirac Neutrinos

$$\delta m_\nu = -C_3^6(\nu) \frac{\nu^3}{2\sqrt{2}\Lambda^2} \quad \frac{\mu_\nu}{\mu_B} = -4\sqrt{2} \left(\frac{m_{e\nu}}{\Lambda^2}\right) C_+(\nu) \quad C_+ = C_1^6 + C_2^6$$

Operator Basis:

$$O_M^{(4)} = \bar{L}\tilde{\phi}\nu_R$$

$$\tilde{\phi} = i\tau_2\phi^*$$

*Matching at
scale Λ*

$$O_1^{(6)} = g_1(\bar{L}\sigma^{\mu\nu}\tilde{\phi})\ell_R B_{\mu\nu}$$

$$O_2^{(6)} = g_2(\bar{L}\sigma^{\mu\nu}\tau^a\tilde{\phi})\nu_R W_{\mu\nu}^a$$

$$O_3^{(6)} = (\bar{L}\tilde{\phi}\nu_R)(\phi^+\phi)$$

*Close under
renormalization*

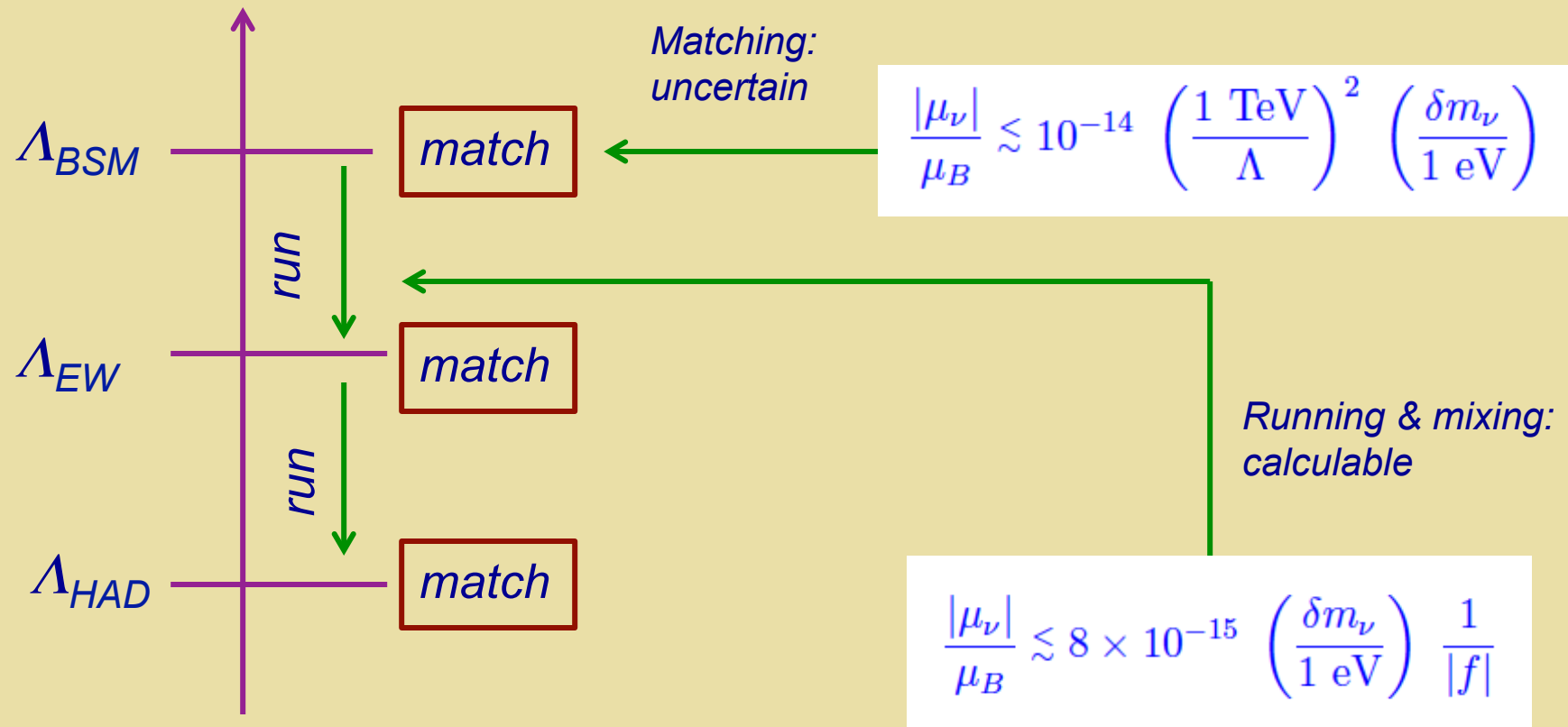
Dirac Neutrinos: Matching & “Naturalness”

Solution with $C_3^6(\Lambda) = 0$: δm_ν generated entirely from radiative corrections via $k_{loop} \sim \Lambda$, thereby inducing nonzero $C_M^4(\Lambda)$

$$\delta m_\nu \sim \frac{\alpha}{32\pi} \frac{\Lambda^2}{m_e} \frac{\mu_\nu}{\mu_B},$$

$$\frac{\mu_\nu}{\mu_B} \lesssim 10^{-14} \left(\frac{1 \text{ TeV}}{\Lambda} \right)^2 \left(\frac{\delta m_\nu}{1 \text{ eV}} \right)$$

Interpretation



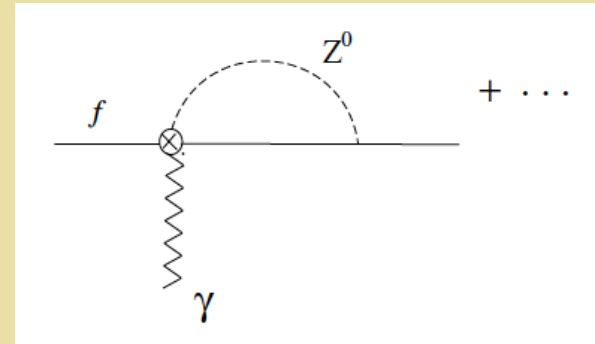
Applying to TVPC Interactions & EDMs

- *Khripolovich approach: compute in EFT II w/ cut-off regulator*
- *Khriplovich approach a la MR-M: compute in EFT II w/ dim reg*
- *Recast in EFT I framework*

Applying to TVPC Interactions & EDMs

$$\mathcal{O}_7^{\gamma g} = C_7^{\gamma g} \bar{\psi} \sigma_{\mu\nu} \lambda^a \psi F^{\mu\lambda} G_\lambda^{a\nu}$$

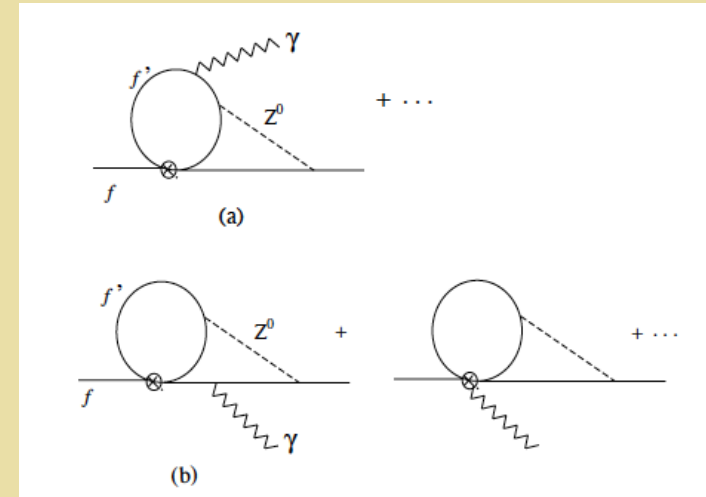
$$\mathcal{O}_7^{\gamma Z} = C_7^{\gamma Z} \bar{\psi} \sigma_{\mu\nu} \psi F^{\mu\lambda} Z_\lambda^\nu$$



$$C_5^f \sim e C_7^{\gamma Z} \left(\frac{M_Z}{\Lambda_{\text{TVPC}}} \right)^2 \left(\frac{1}{s_W c_W} \right) g_A^f \left(\frac{1}{96\pi^2} \right) \ln \frac{M_Z^2}{\mu^2}$$

Applying to TVPC Interactions & EDMs

$$\mathcal{O}_7^{ff'} = C_7^{ff'} \bar{\psi}_f \overleftrightarrow{D}_\mu \gamma_5 \psi_f \bar{\psi}_{f'} \gamma^\mu \gamma_5 \psi_{f'}$$



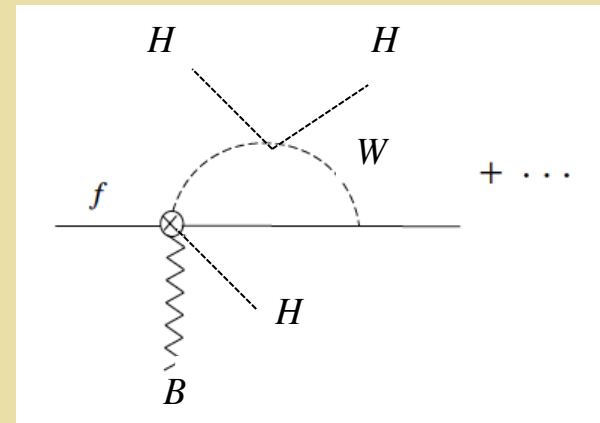
$$C_5^f \sim -e C_7^{ff'} \left(\frac{5}{12} \right) \left(\frac{M_Z}{\Lambda_{\text{TVPC}}} \right)^2 Q_f g_V g_A \left(\frac{G_F M_Z^2}{\sqrt{2}} \right) \times \left(\frac{1}{8\pi^2} \right)^2 \left(\ln \frac{M_Z^2}{\mu^2} \right)^2,$$

The EFT I Computation

$$\mathcal{O}_{\text{fWB}}^{(8)} = \bar{F} \sigma^{\mu\nu} \frac{\tau^a}{2} H f_R \widetilde{W}_{\mu\alpha}^a B_\nu^\alpha$$

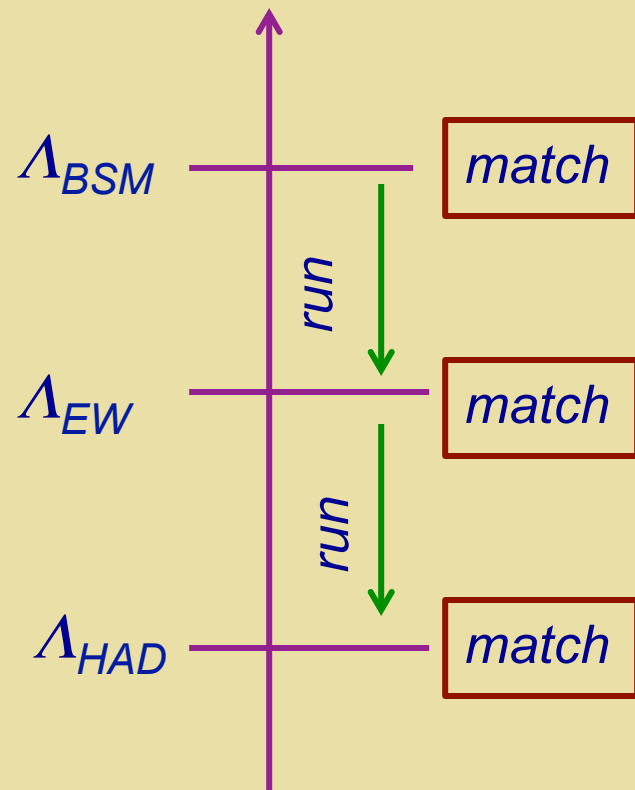
$$\mathcal{O}_{\text{fW}}^{(8)} = \bar{F} \sigma^{\mu\nu} \frac{\tau^a}{2} H f_R \widetilde{W}_{\mu\alpha}^a H^\dagger H$$

$$\mathcal{O}_{\text{fB}}^{(8)} = \bar{F} \sigma^{\mu\nu} H f_R \widetilde{B}_{\mu\alpha}^a H^\dagger H$$



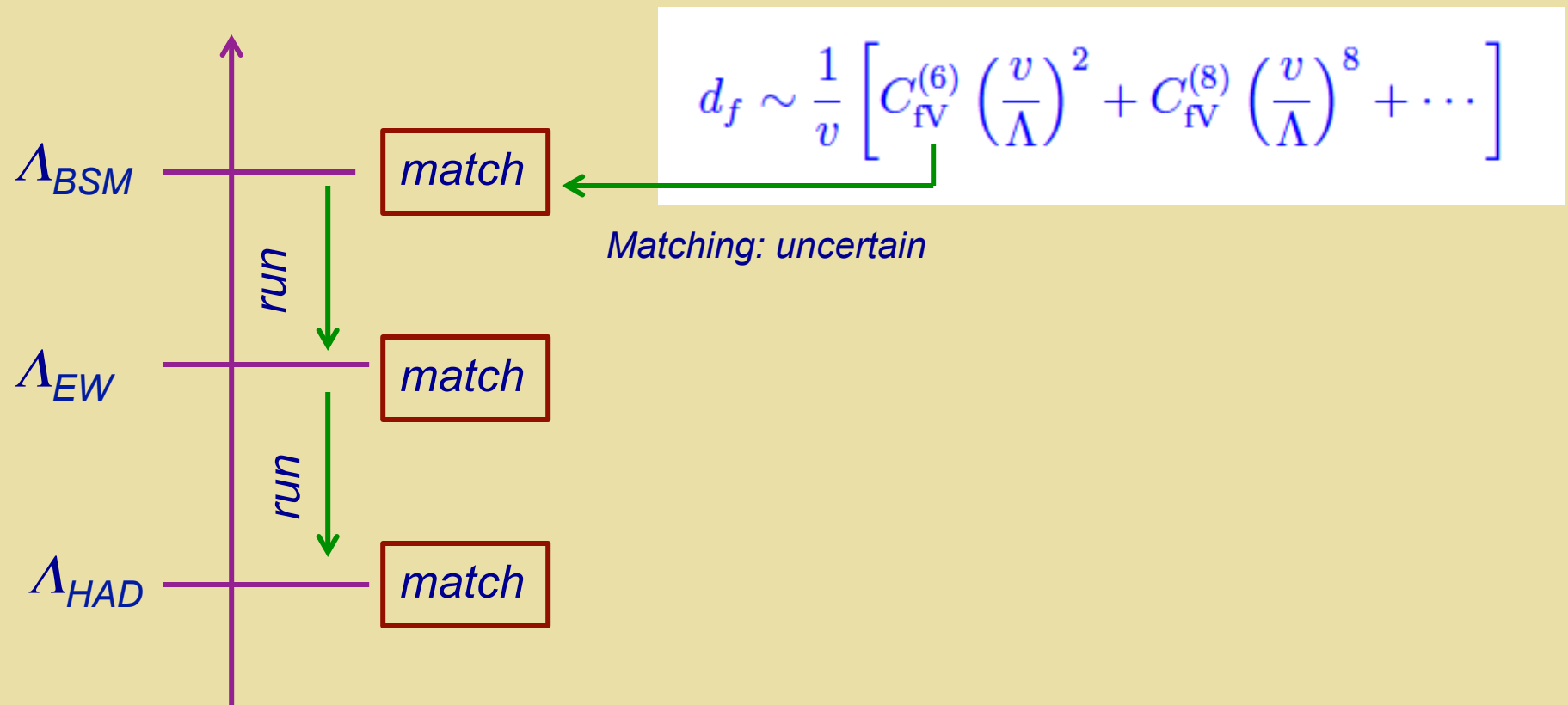
$$C_{\text{fV}}^{(8)} \sim \left(\frac{\alpha}{4\pi} \right) C_{\text{TVPC}}^{(8)}$$

Interpretation

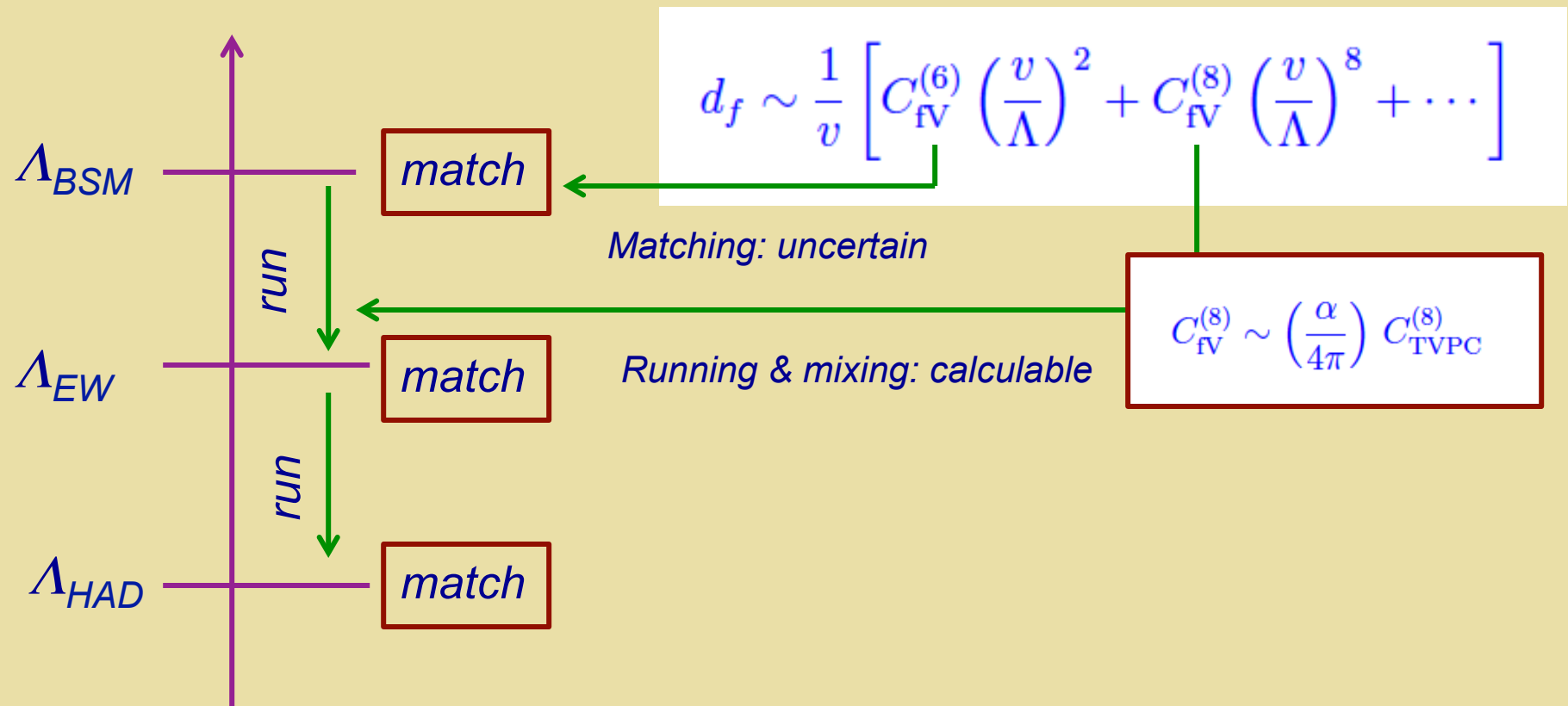


$$d_f \sim \frac{1}{v} \left[C_{fV}^{(6)} \left(\frac{v}{\Lambda} \right)^2 + C_{fV}^{(8)} \left(\frac{v}{\Lambda} \right)^8 + \dots \right]$$

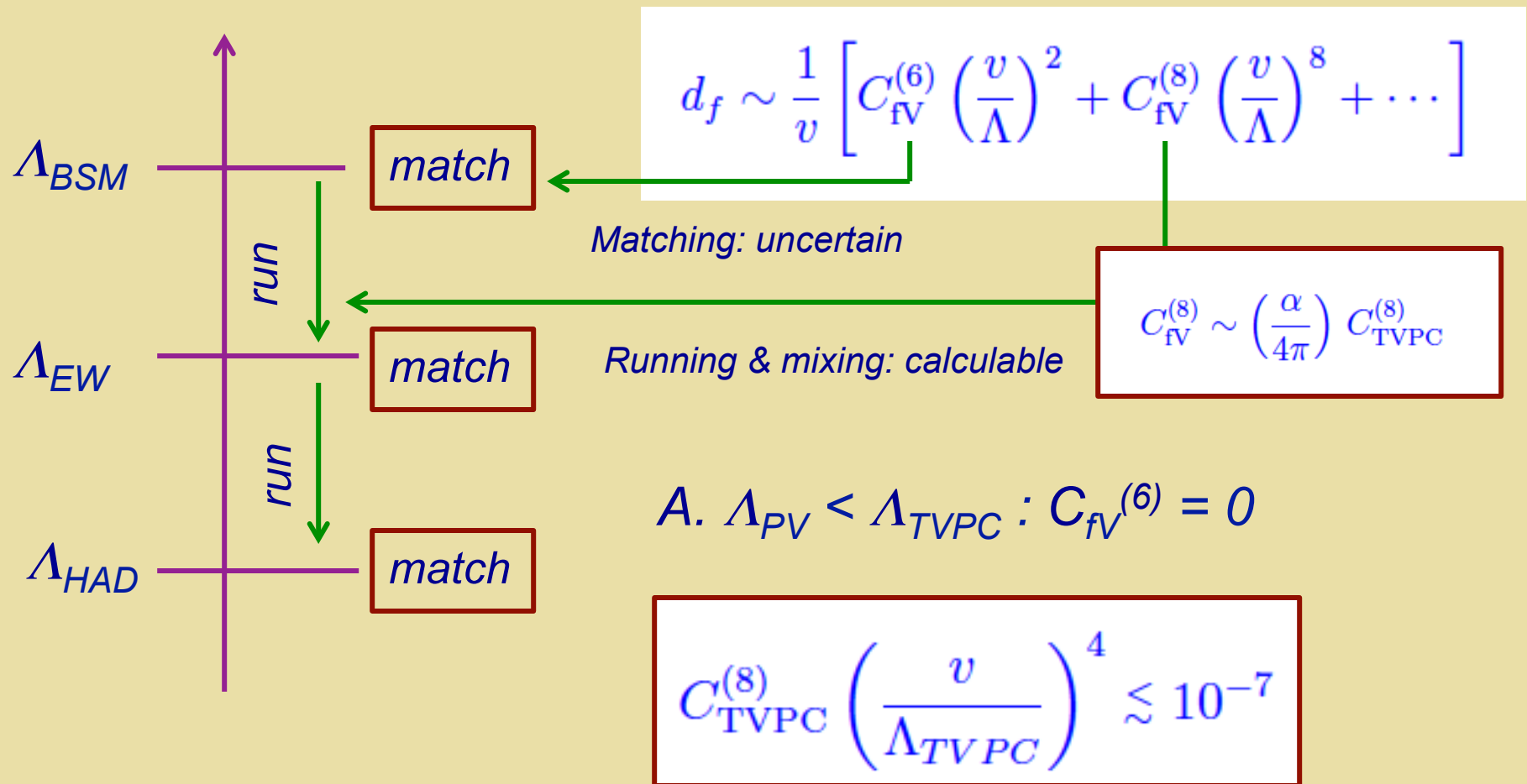
Interpretation



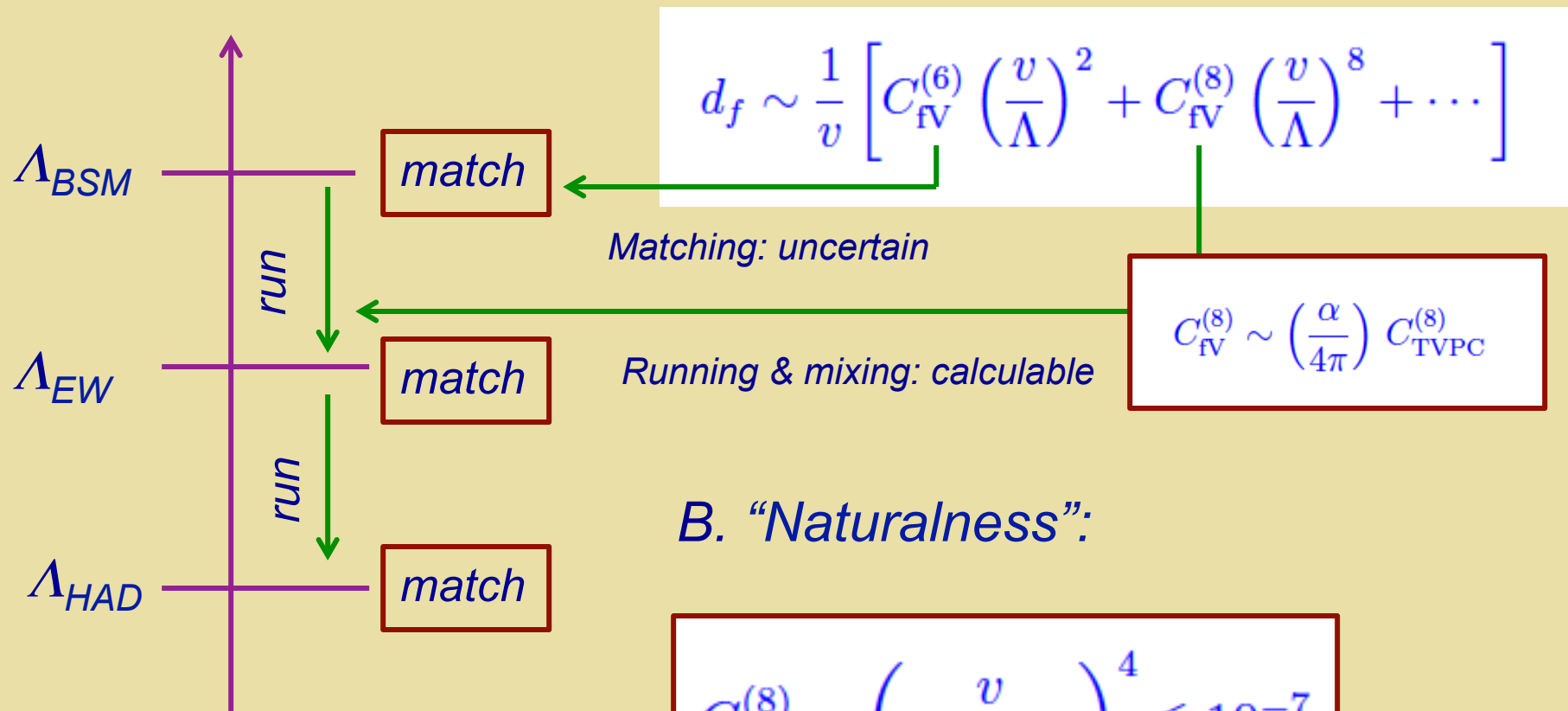
Interpretation



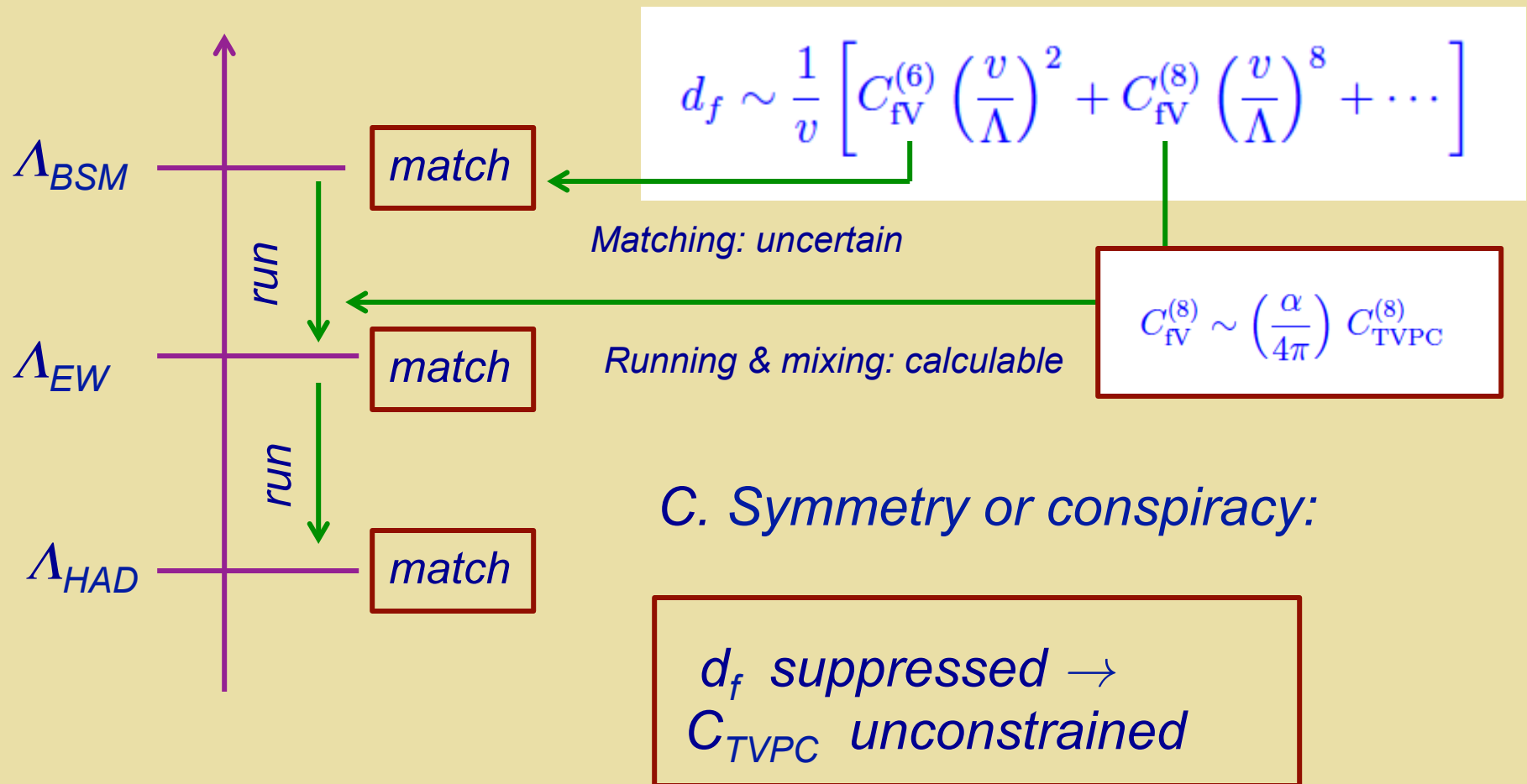
Limits: Short Distance Parity Cons



Limits: Naturalness



Limits: Symmetry or Conspiracy



Implications

A. $\Lambda_{PV} < \Lambda_{TVPC} : C_{fV}^{(6)} = 0$

B. “Naturalness”

$$\alpha_T = \frac{\langle f | \mathcal{O}_{TVPC}^{(8)} | i \rangle}{\langle f | \mathcal{O}_{QCD} | i \rangle} \sim C_{TVPC}^{(8)} \left(\frac{v}{\Lambda_{TVPC}} \right) \left(\frac{p}{\Lambda_{TVPC}} \right)^3 \lesssim 10^{-15}$$

for $\Lambda_{TVPC} \sim v$, $p \sim 1$ GeV

C. Symmetry or conspiracy

$$\alpha_T = \frac{\langle f | \mathcal{O}_{TVPC}^{(8)} | i \rangle}{\langle f | \mathcal{O}_{QCD} | i \rangle} \sim C_{TVPC}^{(8)} \left(\frac{v}{\Lambda_{TVPC}} \right) \left(\frac{p}{\Lambda_{TVPC}} \right)^3 \lesssim 10^{-7}$$

for $\Lambda_{TVPC} \sim v$, $p \sim 1$ GeV, and $C_{TVPC} \sim 1$

Implications: Further Thoughts

C. Symmetry or conspiracy

$$\mathcal{O}_{fWB}^{(8)} = \bar{F} \sigma^{\mu\nu} \frac{\tau^a}{2} H f_R \tilde{W}_{\mu\alpha}^a B_\nu^\alpha$$

$$\mathcal{O}_{fWW}^{(8)} = \bar{F} \sigma^{\mu\nu} H f_R \tilde{W}_{\mu\alpha}^a W_\mu^{a\alpha}$$

$$\mathcal{O}_{fBB}^{(8)} = \bar{F} \sigma^{\mu\nu} H f_R \tilde{B}_{\mu\alpha} B_\nu^\alpha$$

$$\mathcal{O}_{f\gamma\gamma}^{(8)} = \bar{F} \sigma^{\mu\nu} f_R \tilde{F}_{\mu\alpha} F_\nu^\alpha$$

$$\eta \rightarrow 3\gamma$$

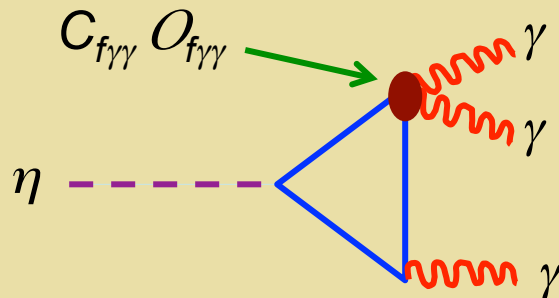
$$\mathcal{O}_{f\gamma Z}^{(8)} = \bar{F} \sigma^{\mu\nu} f_R \tilde{F}_{\mu\alpha} Z_\nu^\alpha$$

EDM

$$C_{fWB}, C_{fWW}, C_{fBB}$$



$$C_{f\gamma Z} = 0, C_{f\gamma\gamma} \neq 0$$



$$\frac{v}{\Lambda} \frac{1}{\Lambda^3} \eta F^3$$

Summary & Outlook

- *C-Violating \leftrightarrow TVPC interactions are a largely unexplored direction for fundamental symmetry tests*
- *Analyzing their effects for light quark systems requires an EFT approach, as they do not arise at tree-level via renormalizable gauge interactions*
- *In general, EDMs place stringent constraints on such interactions via EW radiative corrections from the standpoint of short distance parity restoration and/or naturalness*
- *Exceptions may exist in the presence of a conspiracy or new symmetry at the TVPC matching scale*
- *Magnitude of low-energy amplitude $\sim (p/\Lambda)^3 < 10^{-7}$ for $\Lambda > v$*
- *C-Violating \leftrightarrow TVPC interactions are an interesting direction worthy of further exploration*

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