

Status and Prospects of Higgs CP Properties with CMS and ATLAS

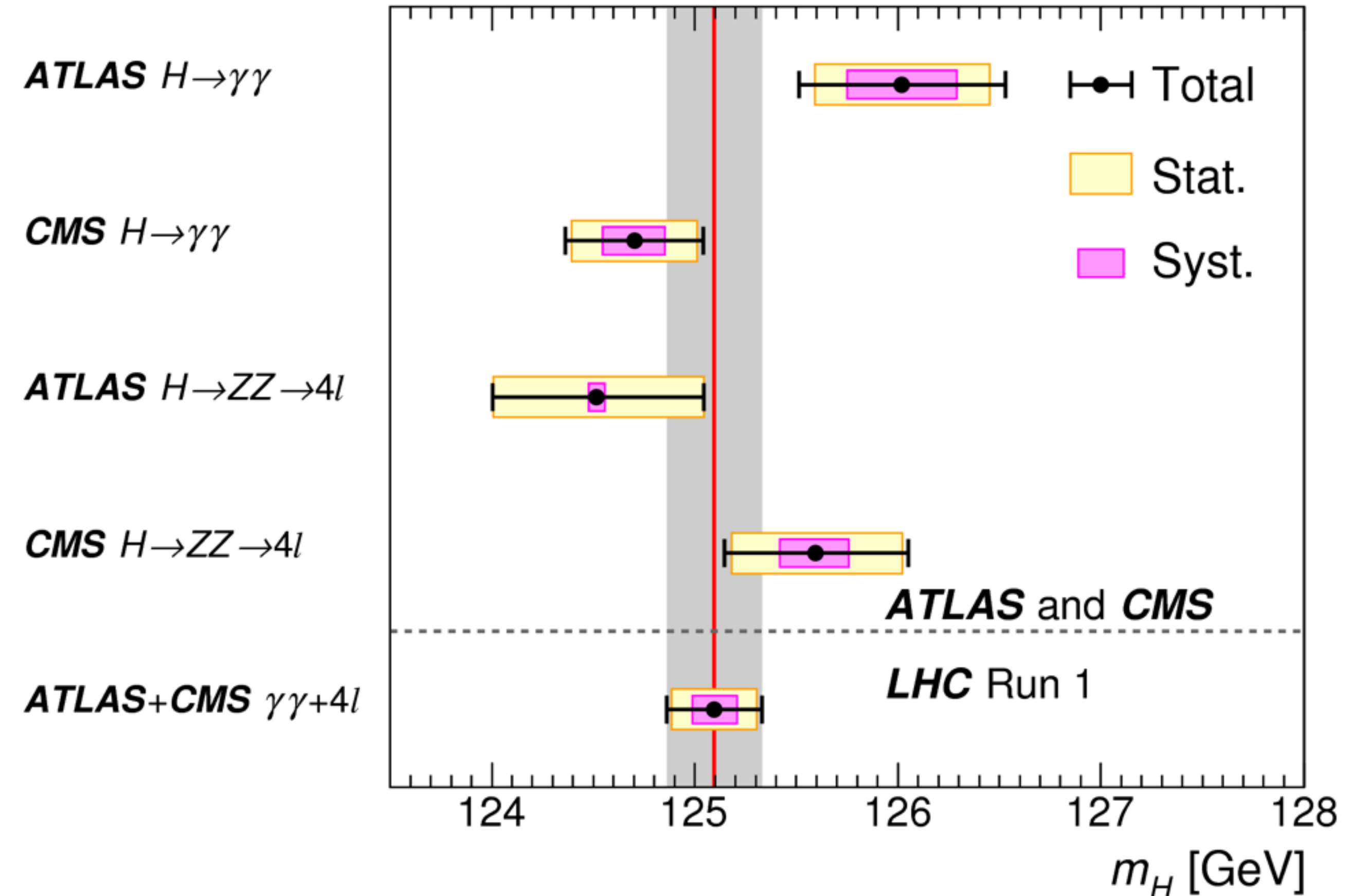
Andrew Whitbeck – Fermilab

*The CP Nature of the Higgs Boson Workshop
@ U. Mass. Amherst*

May 1, 2015

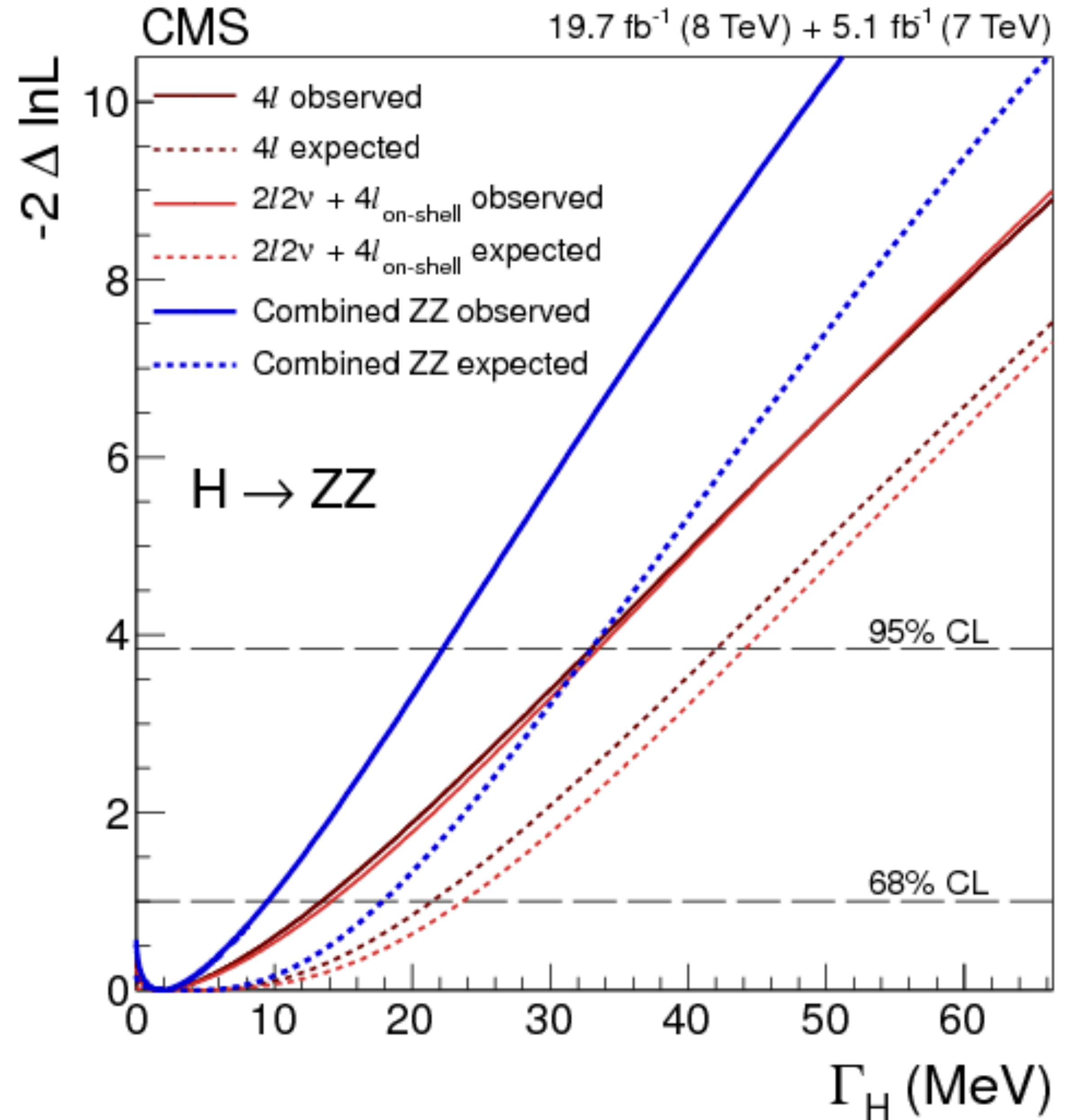
What we know about the Higgs...

- 125 GeV narrow width resonance
 - driven by the two high-resolution decay channels, $ZZ \rightarrow 4\ell$ & $\gamma\gamma$
 - mass has now been determined to be **~.2 GeV!**
- constraints from off-shell production suggest that the width is consistent with the SM prediction
- Branching ratios consistent with SM Higgs boson
 - Only moderate sensitivity, yet, in the fermionic decay channels
- Various production modes seen
 - overall rates are consistent with SM expectation
 - still not sensitive to SM strength ttH production
- In general, couplings to various SM fields (*constrained by yields*) are consistent with SM predictions



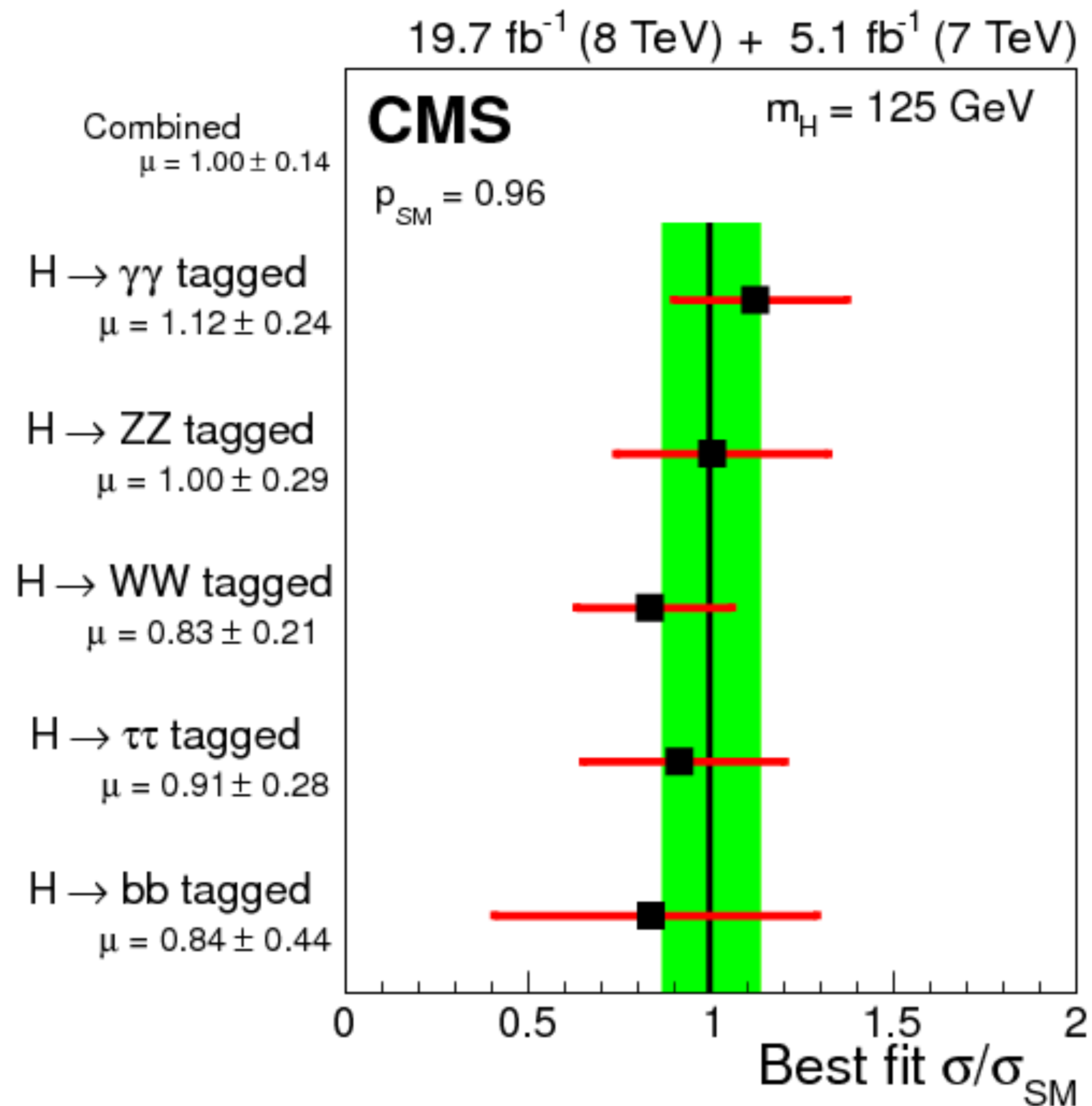
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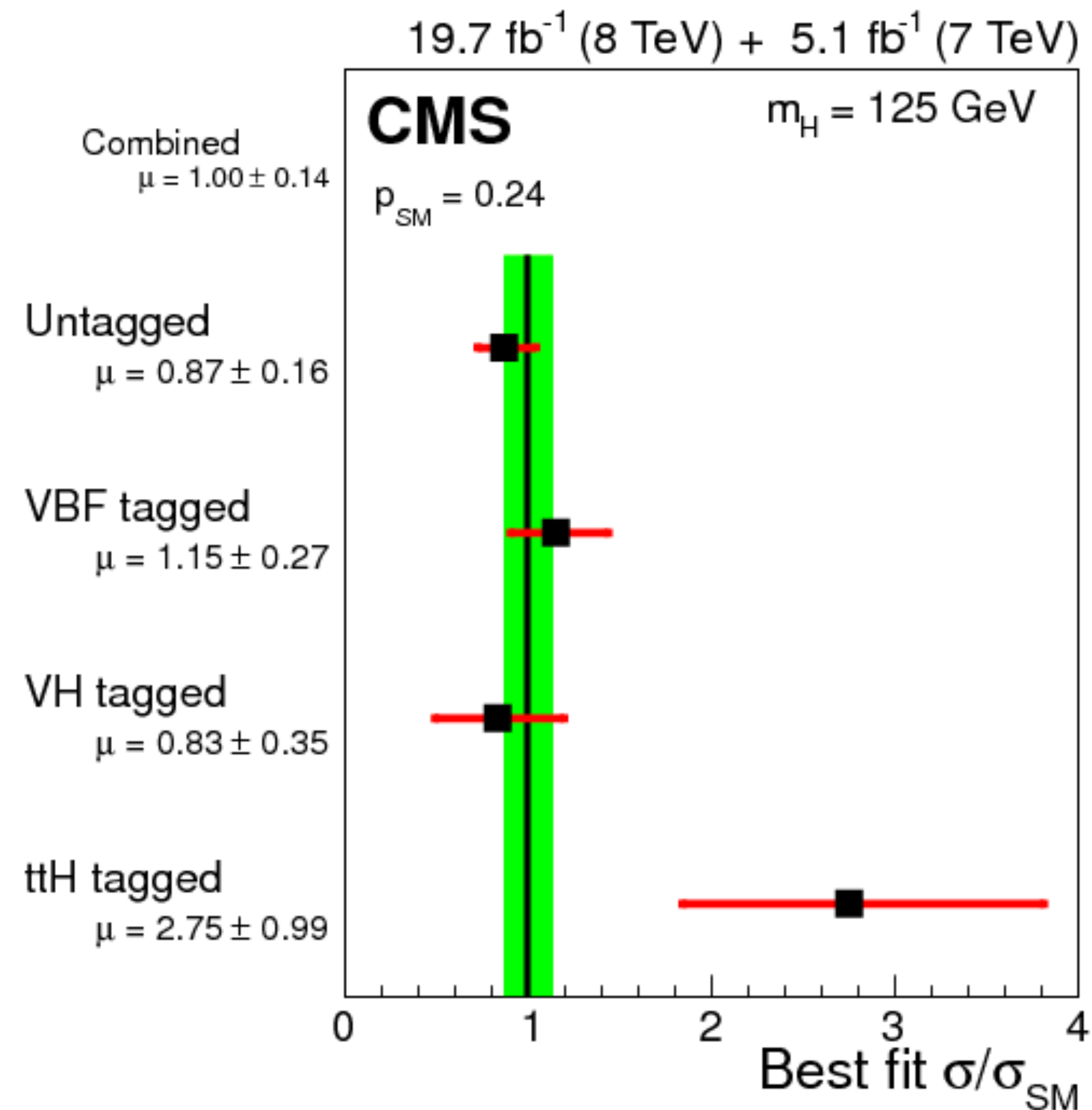
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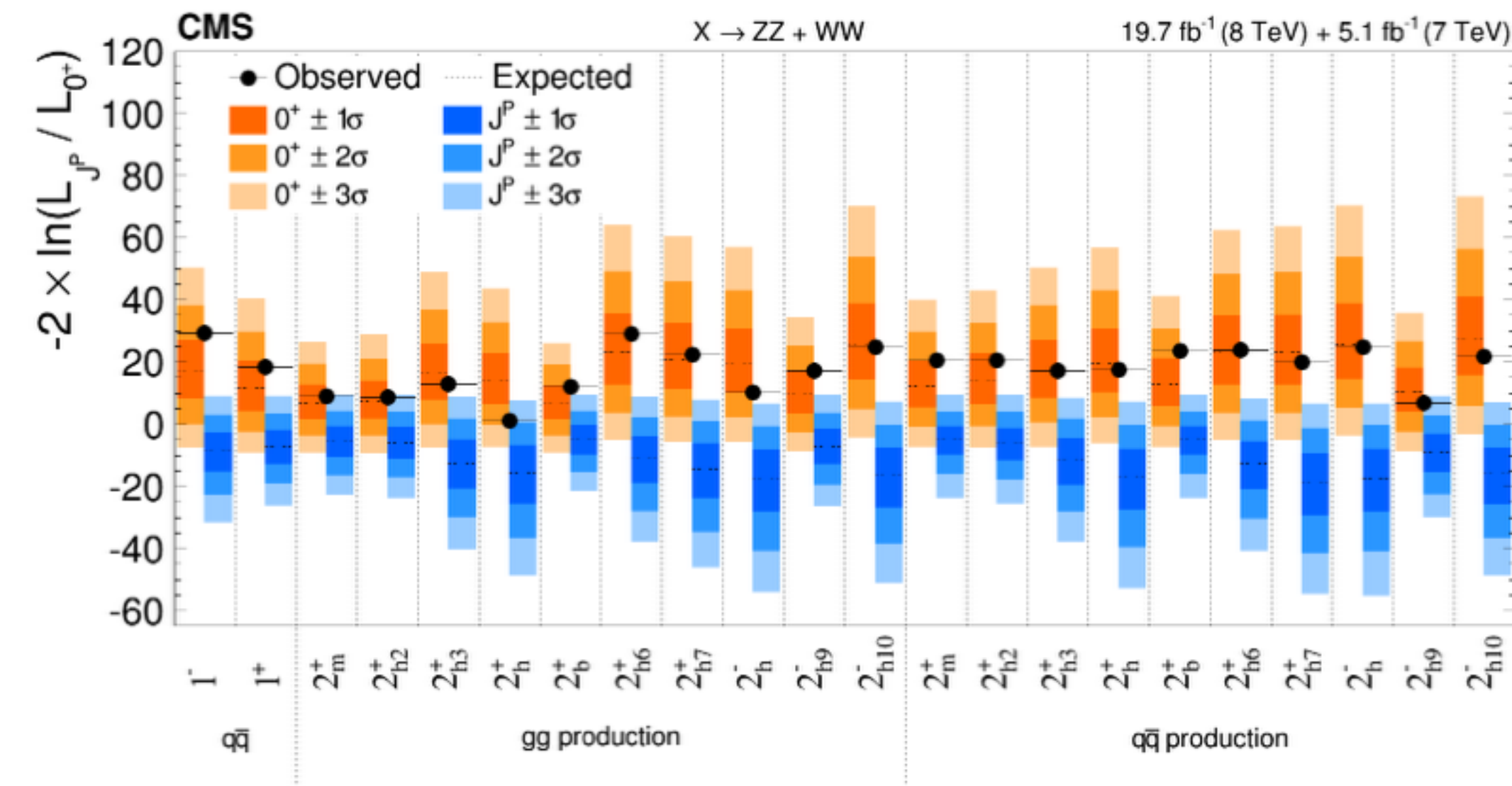
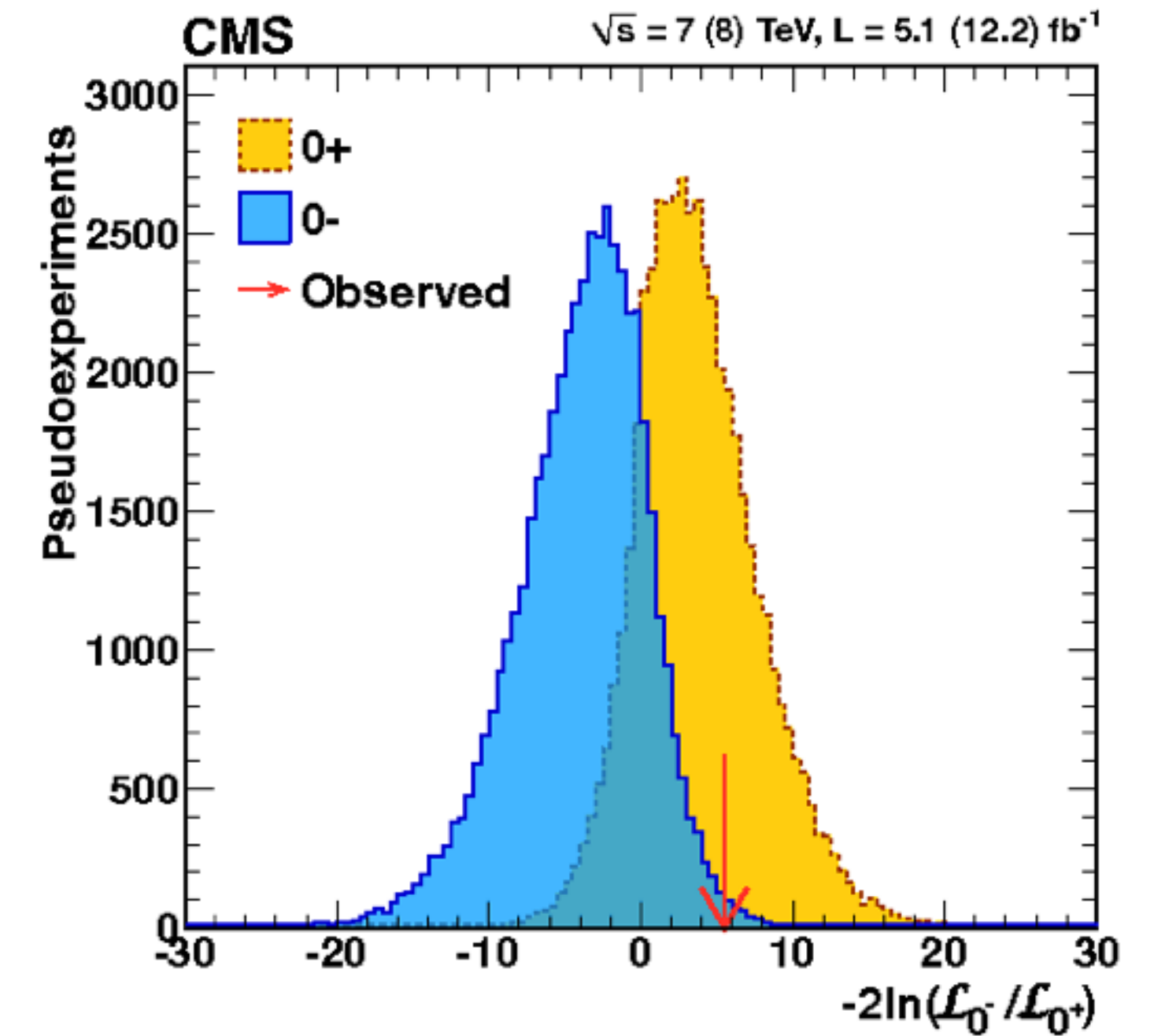
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What we know about the Higgs...

- Kinematic used to test various J^P signal hypotheses
- these measurements use *distributions* for characterizing signal models making them *complementary* to coupling fit based on rates
- early analyses focused on hypothesis separation, but now we are starting to constrain the *tensor structure of a generic spin-0 resonance decaying to a pair of vector bosons*
- *Anomalous HVV couplings* are currently the primary focus since these are the channels that have the most sensitivity to signal events



Outline

- CMS & ATLAS 4ℓ analyses
 - detailed explanation of observables, model parameters, and methods
- $W(\ell\nu)W(\ell\nu)$ analyses & ZZ-WW combination
- $\gamma\gamma$ & $Z\gamma$ couplings
- Prospects for the future

Model Parameters

- Starting from a more generic set of couplings \longrightarrow
- More convenient notation defines parameters which are functions of couplings squared & cross sections
- anomalous couplings are redefined onto a unit interval:

$$A(X_{J=0} \rightarrow V_1 V_2) \sim v^{-1} \left(\left[a_1 - e^{i\phi_{\Lambda_1}} \frac{q_{Z_1}^2 + q_{Z_2}^2}{(\Lambda_1)^2} \right] m_Z^2 \epsilon_{Z_1}^* \epsilon_{Z_2}^* \right. \\ \left. + a_2 f_{\mu\nu}^{*(Z_1)} f^{*(Z_2),\mu\nu} + a_3 f_{\mu\nu}^{*(Z_1)} \tilde{f}^{*(Z_2),\mu\nu} \right. \\ \left. + a_2^{Z\gamma} f_{\mu\nu}^{*(Z)} f^{*(\gamma),\mu\nu} + a_3^{Z\gamma} f_{\mu\nu}^{*(Z)} \tilde{f}^{*(\gamma),\mu\nu} \right. \\ \left. + a_2^{\gamma\gamma} f_{\mu\nu}^{*(\gamma_1)} f^{*(\gamma_2),\mu\nu} + a_3^{\gamma\gamma} f_{\mu\nu}^{*(\gamma_1)} \tilde{f}^{*(\gamma_2),\mu\nu} \right)$$


$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda_1} / (\Lambda_1)^4} \quad \phi_{a3} = \arg \left(\frac{a_3}{a_1} \right) \\ f_{a2} = \frac{|a_2|^2 \sigma_2}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda_1} / (\Lambda_1)^4} \quad \phi_{a2} = \arg \left(\frac{a_2}{a_1} \right) \\ f_{\Lambda_1} = \frac{\tilde{\sigma}_{\Lambda_1} / (\Lambda_1)^4}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda_1} / (\Lambda_1)^4} \quad \phi_{\Lambda_1},$$

a_1 : SM HZZ couplings
 a_2 : higher dim. 0^+ ZZ coupling
 a_3 : 0^+ ZZ coupling
 $a_2^{Z\gamma}$: 0^+ Z γ coupling
 $a_3^{Z\gamma}$: 0^+ Z γ coupling
 $a_2^{\gamma\gamma}$: 0^+ $\gamma\gamma$ coupling
 $a_3^{\gamma\gamma}$: 0^+ $\gamma\gamma$ coupling


where σ_i is the cross section of the process corresponding to $a_i = 1, a_{j \neq i} = 0$, while $\tilde{\sigma}_{\Lambda_1}$ is the effective cross section of the process corresponding to $\Lambda_1 > 0, a_{j \neq \Lambda_1} = 0$, given in units fb \cdot GeV⁴.

e.g. $f_{a3} = 1, f_{a2} = 0$ for pure pseudoscalar
 $f_{a3} = f_{a2} = 0$ for SM Higgs (up to higher order corrections)

Generic Amplitudes vs. Effective Field Theories



$$\begin{aligned}
 L(\text{HVV}) \sim & a_1 \frac{m_Z^2}{2} \text{HZ}^\mu \text{Z}_\mu + \frac{1}{(\Lambda_1)^2} m_Z^2 \text{HZ}_\mu \square \text{Z}^\mu - \frac{1}{2} a_2 \text{HZ}^{\mu\nu} \text{Z}_{\mu\nu} - \frac{1}{2} a_3 \text{HZ}^{\mu\nu} \tilde{\text{Z}}_{\mu\nu} \\
 & + a_1^{\text{WW}} \frac{m_W^2}{2} \text{HW}^\mu \text{W}_\mu + \frac{1}{(\Lambda_1^{\text{WW}})^2} m_W^2 \text{HW}_\mu \square \text{W}^\mu - \frac{1}{2} a_2^{\text{WW}} \text{HW}^{\mu\nu} \text{W}_{\mu\nu} - \frac{1}{2} a_3^{\text{WW}} \text{HW}^{\mu\nu} \tilde{\text{W}}_{\mu\nu} \\
 & + \frac{1}{(\Lambda_1^{\text{Z}\gamma})^2} m_Z^2 \text{HZ}_\mu \partial_\nu \text{F}^{\mu\nu} - a_2^{\text{Z}\gamma} \text{HF}^{\mu\nu} \text{Z}_{\mu\nu} - a_3^{\text{Z}\gamma} \text{HF}^{\mu\nu} \tilde{\text{Z}}_{\mu\nu} - \frac{1}{2} a_2^{\gamma\gamma} \text{HF}^{\mu\nu} \text{F}_{\mu\nu} - \frac{1}{2} a_3^{\gamma\gamma} \text{HF}^{\mu\nu} \tilde{\text{F}}_{\mu\nu},
 \end{aligned}$$



$$\begin{aligned}
 \mathcal{L}_0^V = & \left\{ c_\alpha \kappa_{\text{SM}} \left[\frac{1}{2} g_{\text{HZZ}} \text{Z}_\mu \text{Z}^\mu + g_{\text{HWW}} \text{W}_\mu^+ \text{W}^{-\mu} \right] \right. \\
 & - \frac{1}{4} \frac{1}{\Lambda} \left[c_\alpha \kappa_{\text{HZZ}} \text{Z}_{\mu\nu} \text{Z}^{\mu\nu} + s_\alpha \kappa_{\text{AZZ}} \text{Z}_{\mu\nu} \tilde{\text{Z}}^{\mu\nu} \right] \\
 & \left. - \frac{1}{2} \frac{1}{\Lambda} \left[c_\alpha \kappa_{\text{HWW}} \text{W}_{\mu\nu}^+ \text{W}^{-\mu\nu} + s_\alpha \kappa_{\text{AWW}} \text{W}_{\mu\nu}^+ \tilde{\text{W}}^{-\mu\nu} \right] \right\} X_0
 \end{aligned}$$

Conversion:

$$f_{g_i} = \frac{r_{i1}^2}{1 + r_{i1}^2}; \quad (i = 2, 4),$$

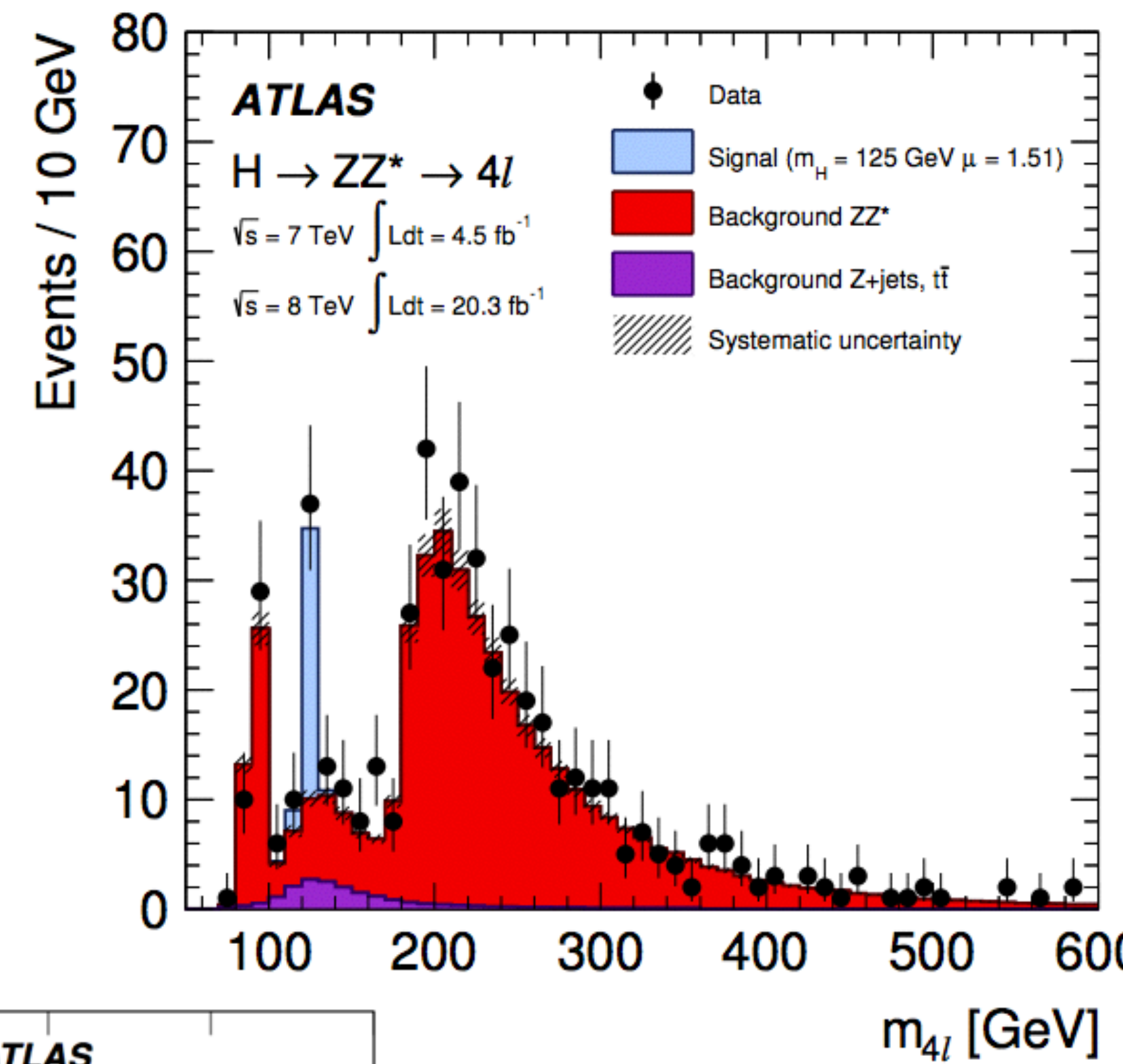
$$\sigma_{\text{HVV}} / \sigma_{\text{SM}} = 0.349$$

$$\sigma_{\text{AVV}} / \sigma_{\text{SM}} = 0.143$$

$$r_{21}^2 = \frac{\sigma_{\text{HVV}}}{\sigma_{\text{SM}}} \left(\frac{\tilde{k}_{\text{HVV}}}{k_{\text{SM}}} \right)^2, \quad \text{and} \quad r_{41}^2 = \frac{\sigma_{\text{AVV}}}{\sigma_{\text{SM}}} \left(\frac{\tilde{k}_{\text{AVV}}}{k_{\text{SM}}} \right)^2 \tan^2 \alpha.$$

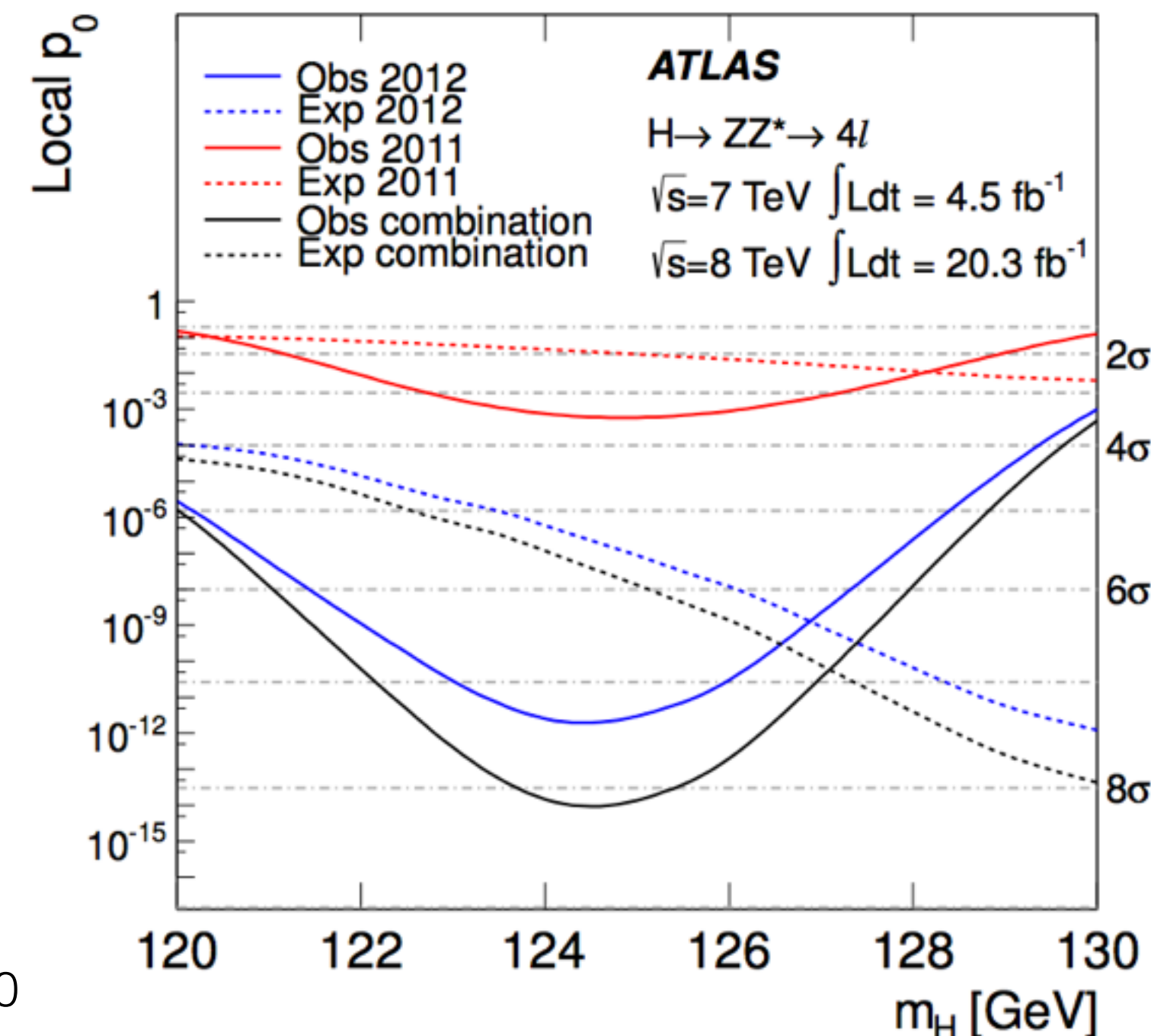
4 ℓ analyses

- two pairs of opposite-sign same flavor leptons (e, μ only)
- SM ZZ contributions taken from NLO MC simulation
- Z+X backgrounds estimated from a loose ID control region
- similar strategy employed by CMS
(see backup for details)



115 < $m_{4\ell}$ < 130 GeV

	Signal	ZZ*	$t\bar{t}, Z + \text{jets}$	Observed
$\sqrt{s} = 7 \text{ TeV}$				
4 μ	1.02 \pm 0.10	0.65 \pm 0.03	0.14 \pm 0.06	3
2 μ 2e	0.47 \pm 0.05	0.29 \pm 0.02	0.53 \pm 0.12	1
2e2 μ	0.64 \pm 0.06	0.45 \pm 0.02	0.13 \pm 0.05	2
4e	0.45 \pm 0.04	0.26 \pm 0.02	0.59 \pm 0.12	2
Total	2.58\pm0.25	1.65\pm0.09	1.39\pm0.26	8
$\sqrt{s} = 8 \text{ TeV}$				
4 μ	5.81 \pm 0.58	3.36 \pm 0.17	0.97 \pm 0.18	13
2 μ 2e	3.00 \pm 0.30	1.59 \pm 0.10	0.52 \pm 0.12	8
2e2 μ	3.72 \pm 0.37	2.33 \pm 0.11	0.84 \pm 0.14	9
4e	2.91 \pm 0.29	1.44 \pm 0.09	0.52 \pm 0.11	7
Total	15.4 \pm 1.5	8.72\pm0.47	2.85\pm0.39	37



arXiv:1408.5191v3



Observables

- An example of a complete set of observables describing the higgs decay kinematics

invariant masses: $m_{4\ell}, m_1, m_2$

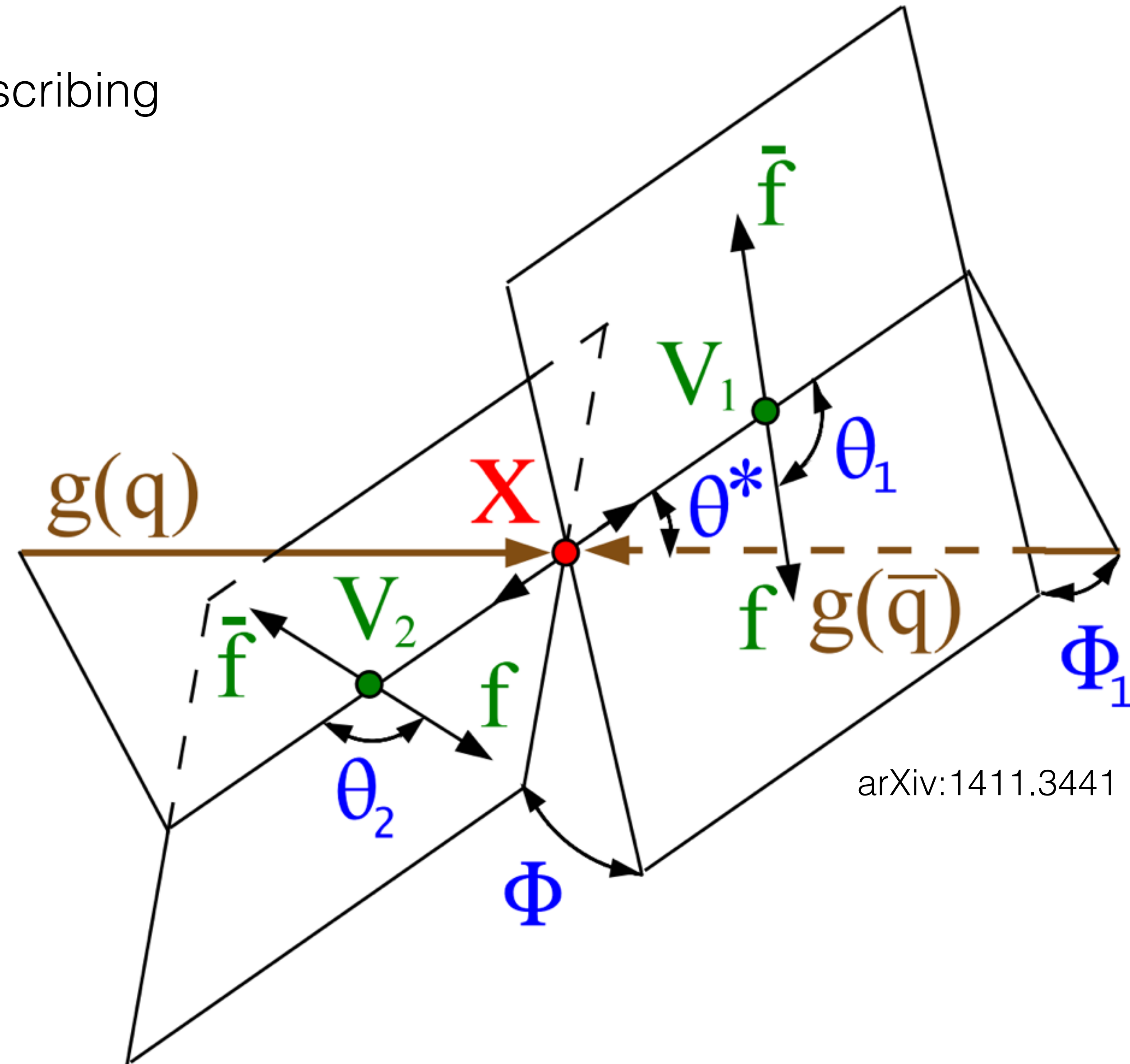
*powerful background discrimination,
insensitive to CP properties*

production angles: $\cos(\theta^*), \Phi_1$

*sensitive to the spin and polarization of the
 X resonance*

decay angles: $\cos(\theta_{1,2}), \Phi$

sensitive to the CP properties of X

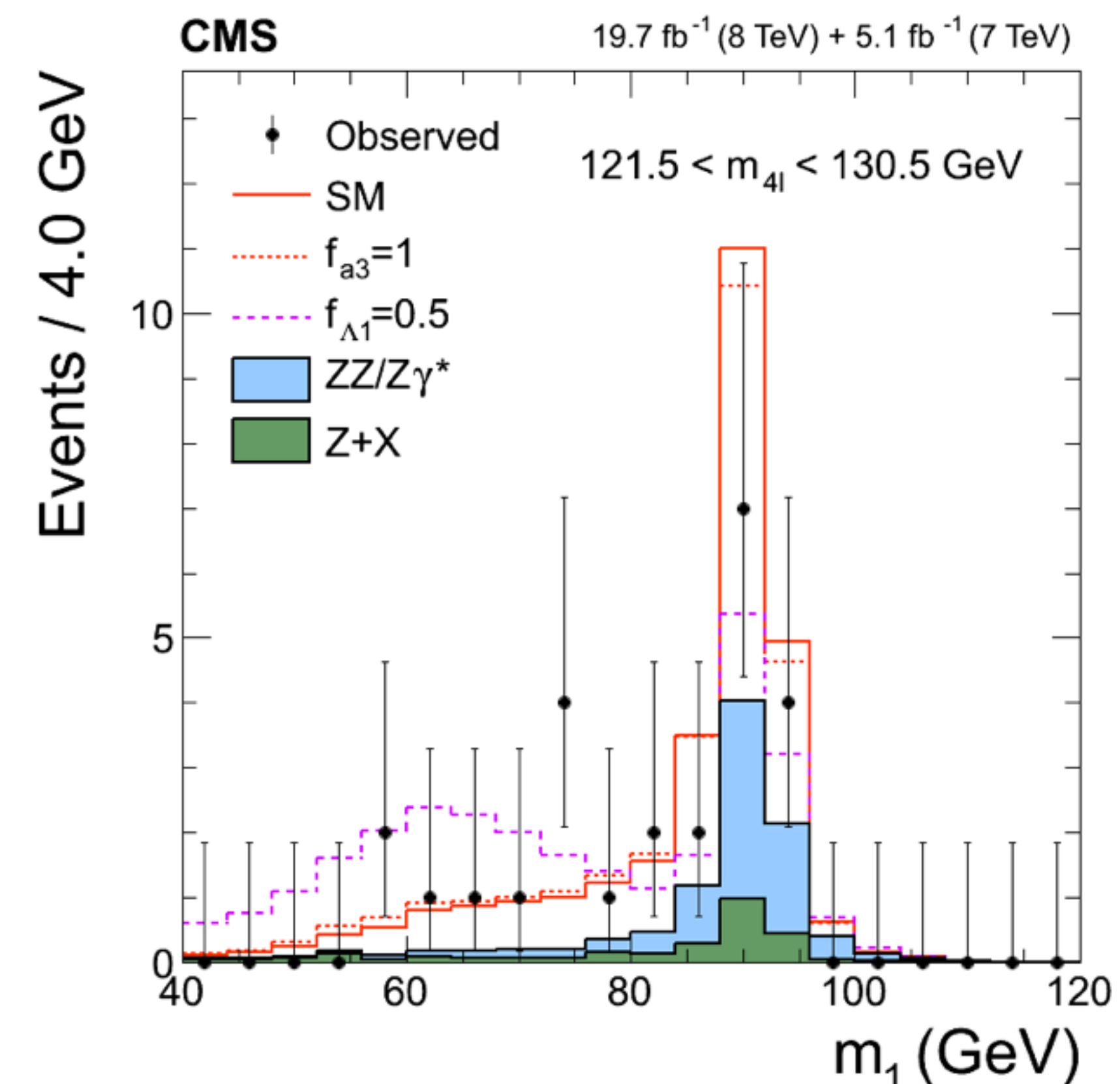
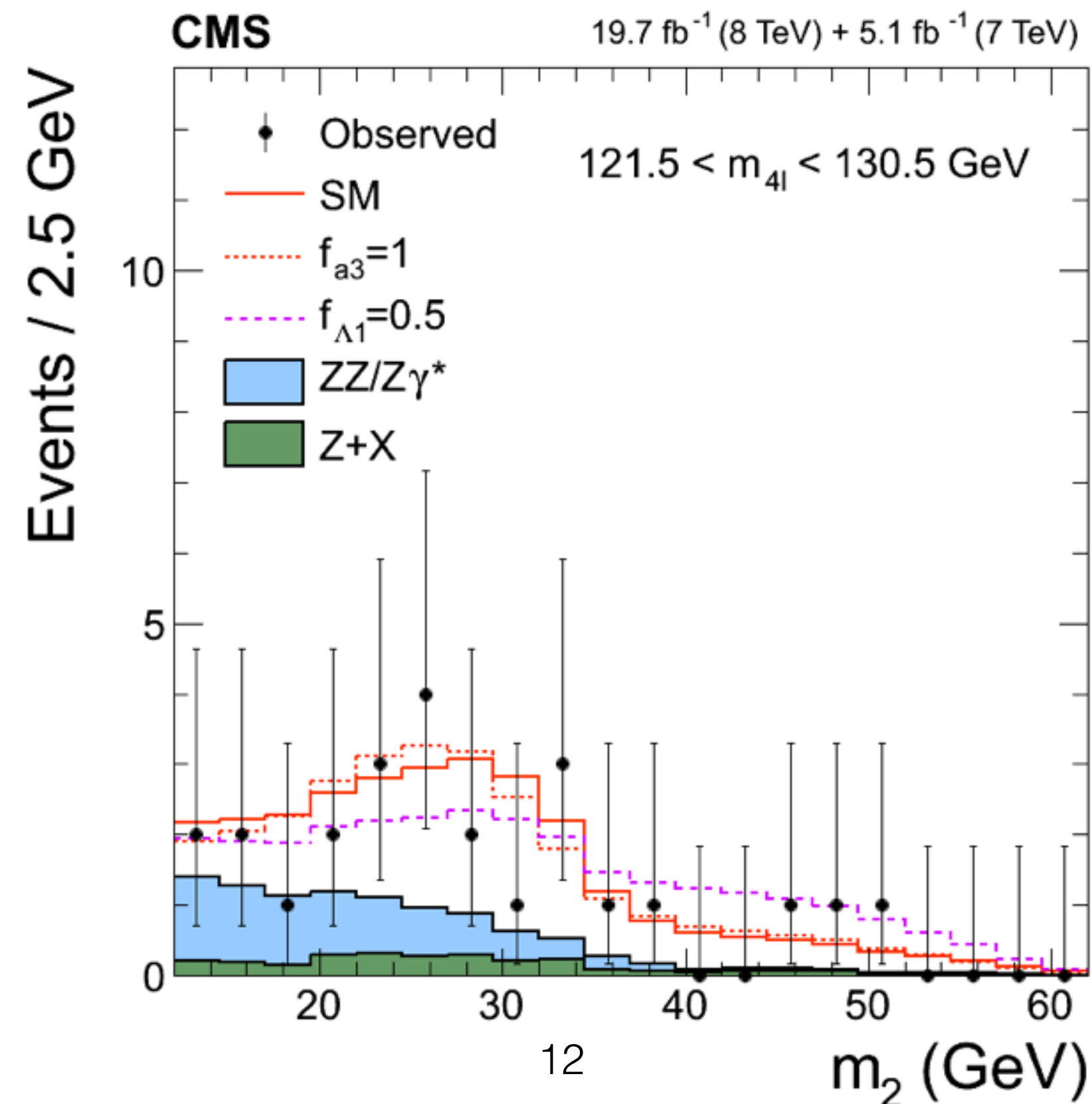
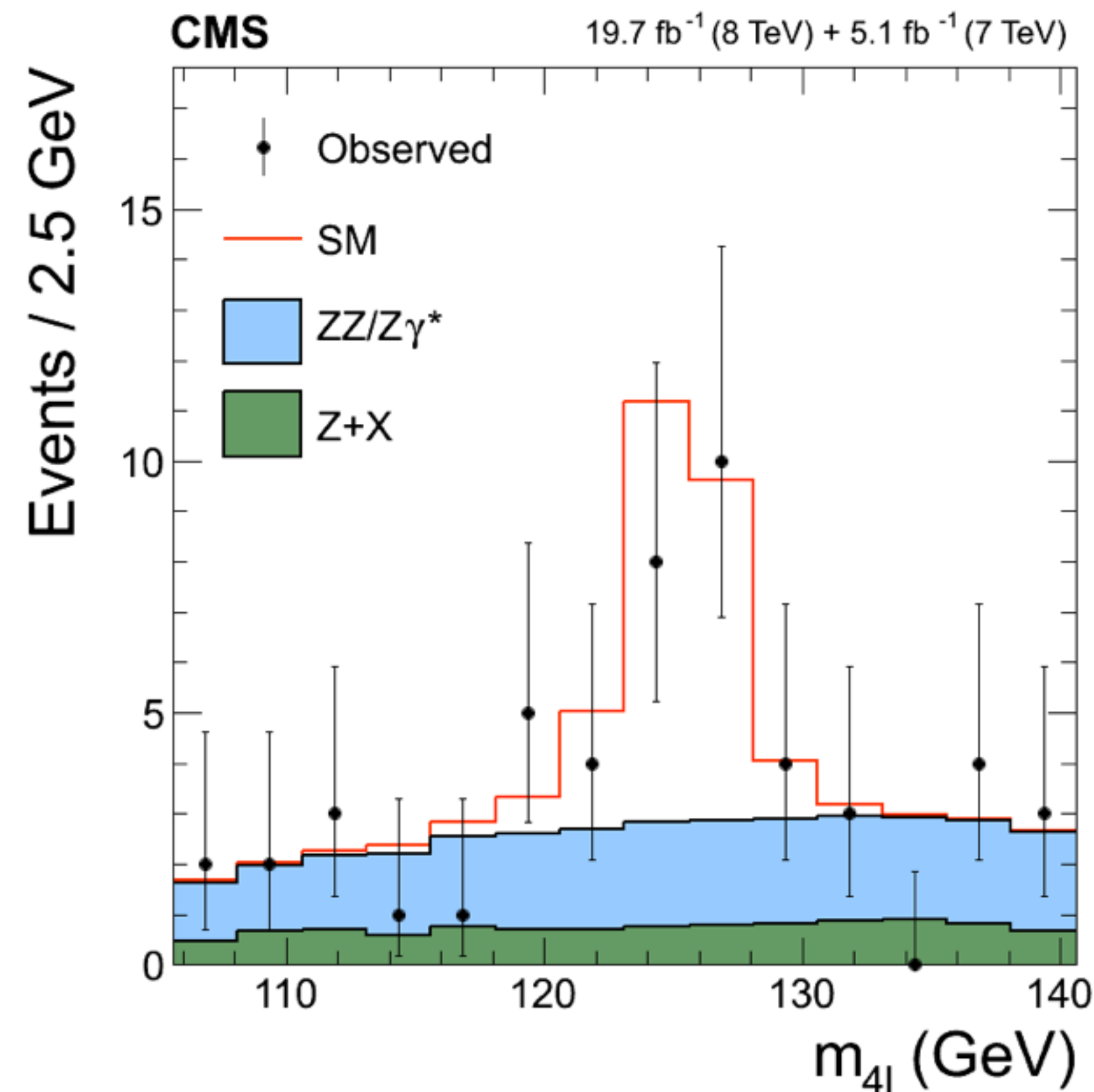


- other variables, p_T & η , are sensitive to NLO effects & insensitive to CP

Distributions masses

- In some sense, the power of the 4ℓ channel stems from the ability to reconstruct masses with very good resolution
- both CMS and ATLAS exploit the masses to discriminate signal from background
- not much sensitivity different CP quantum numbers
- sensitive to form factors

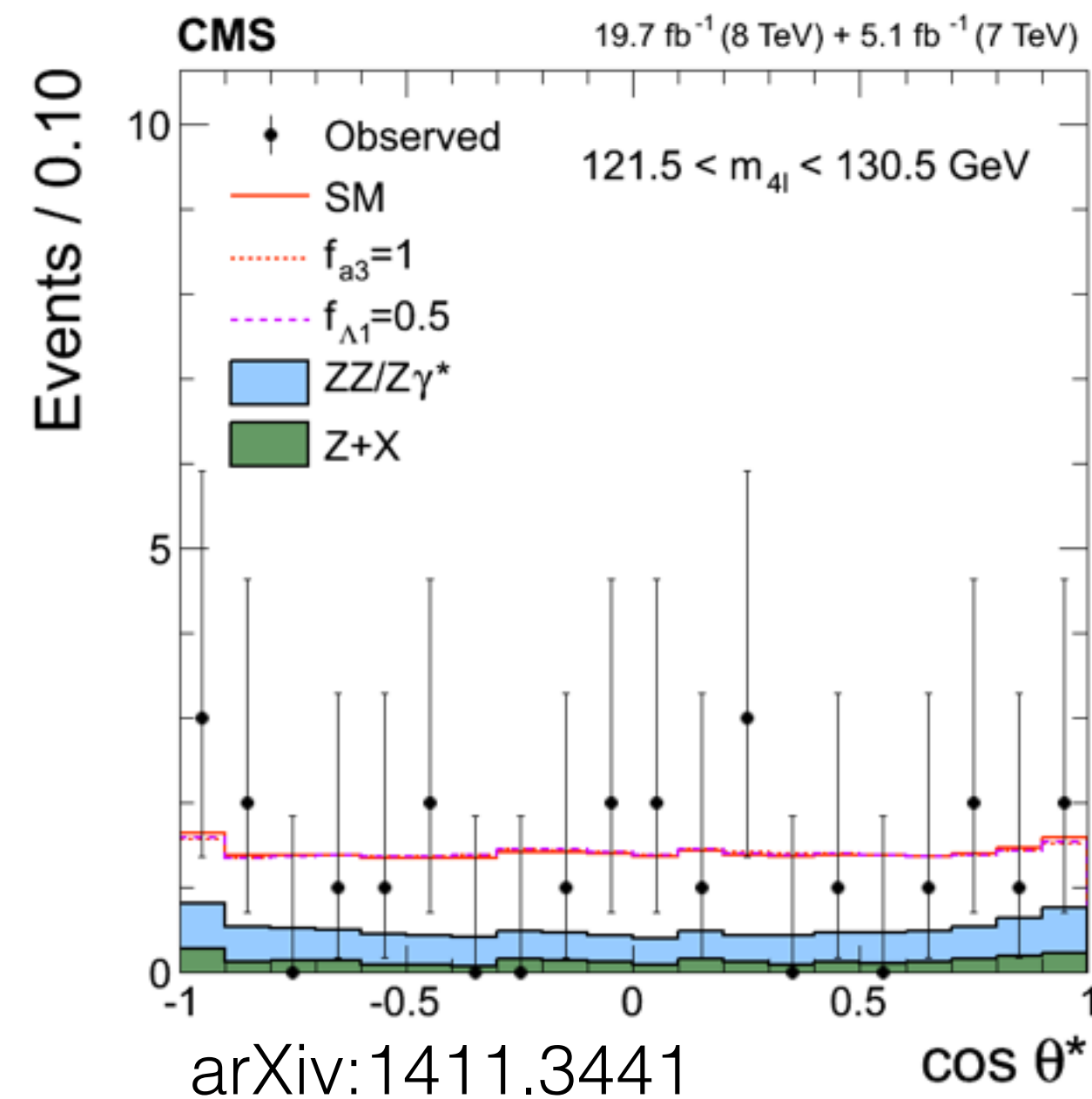
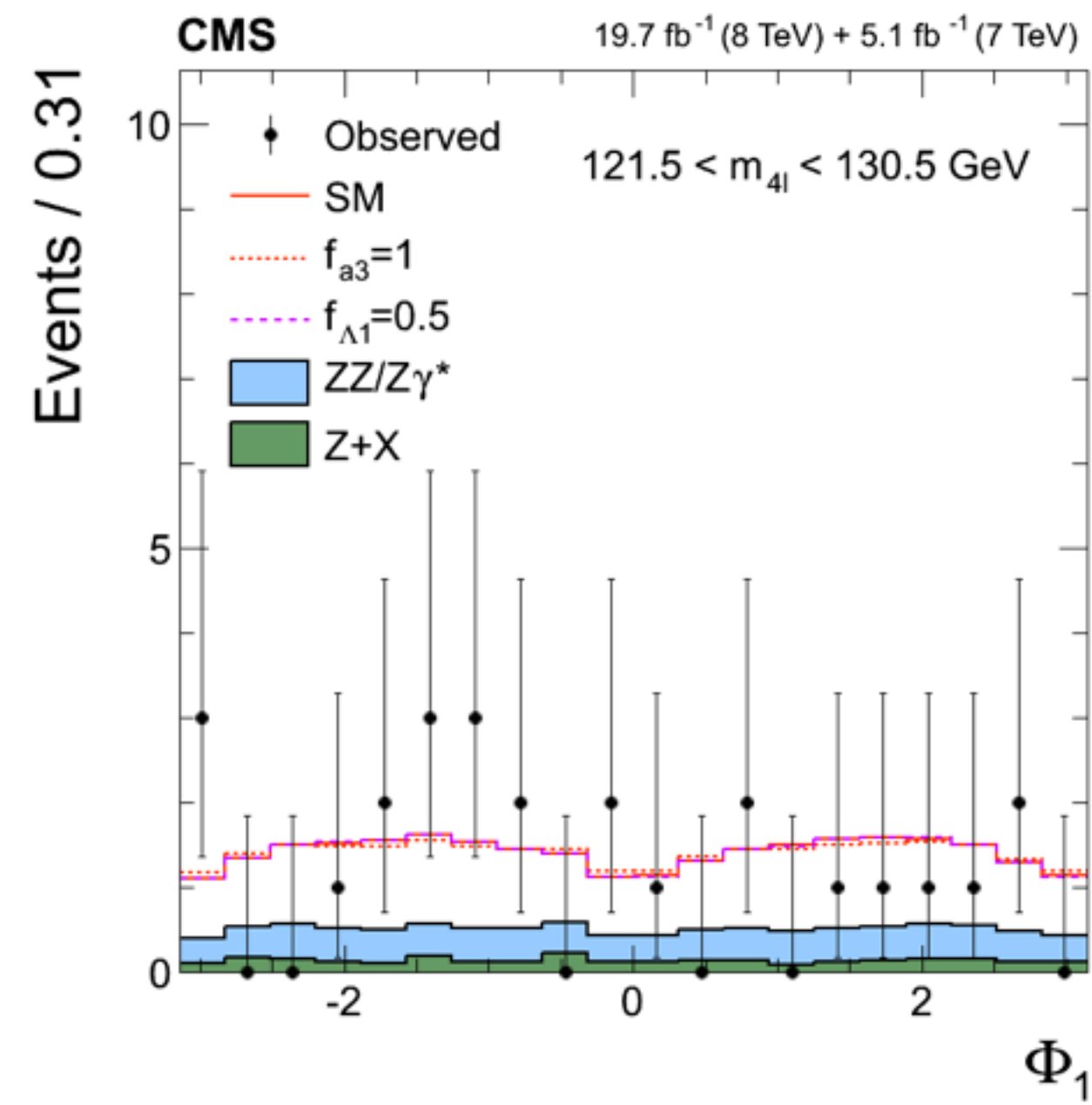
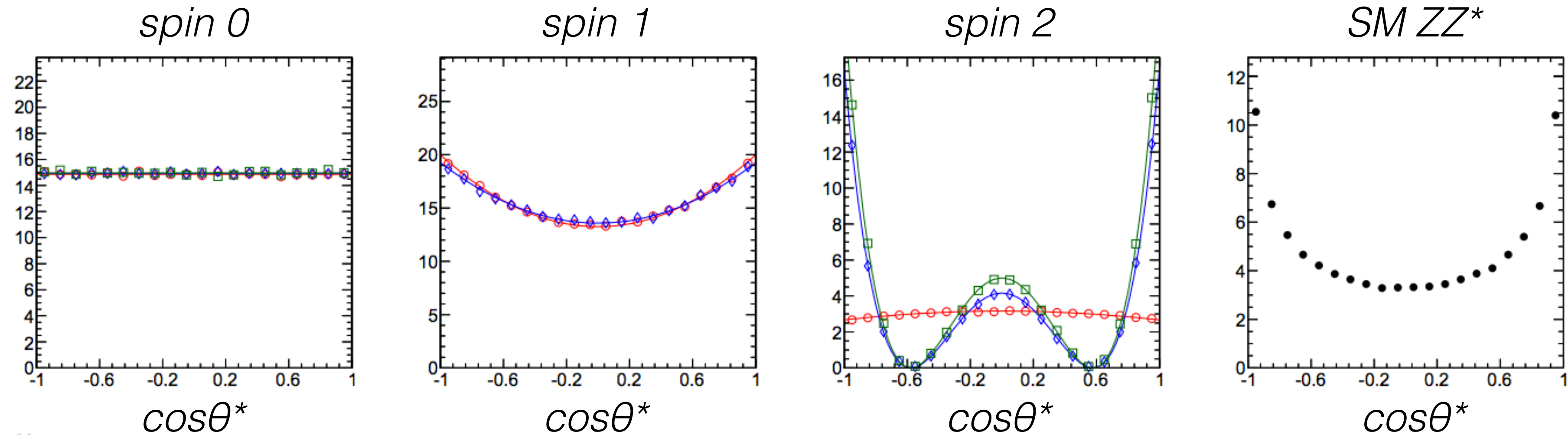
arXiv:1411.3441



Distributions production angles

- Angles describing Z's with respect to incoming partons
- Sensitive to the spin and polarization of X resonance, but aren't sensitive to various CP quantum numbers
- some help in distinguishing signal from backgrounds

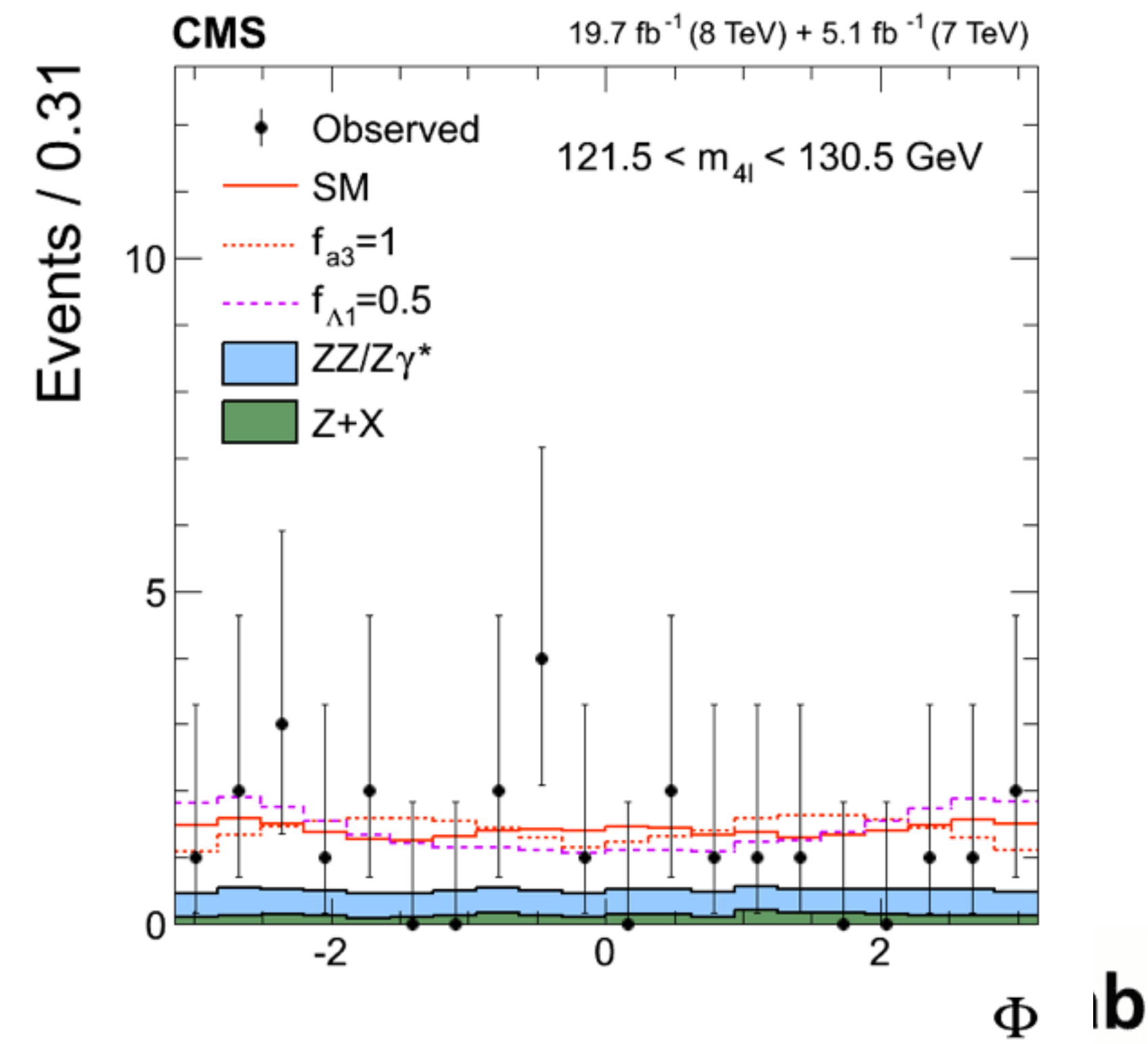
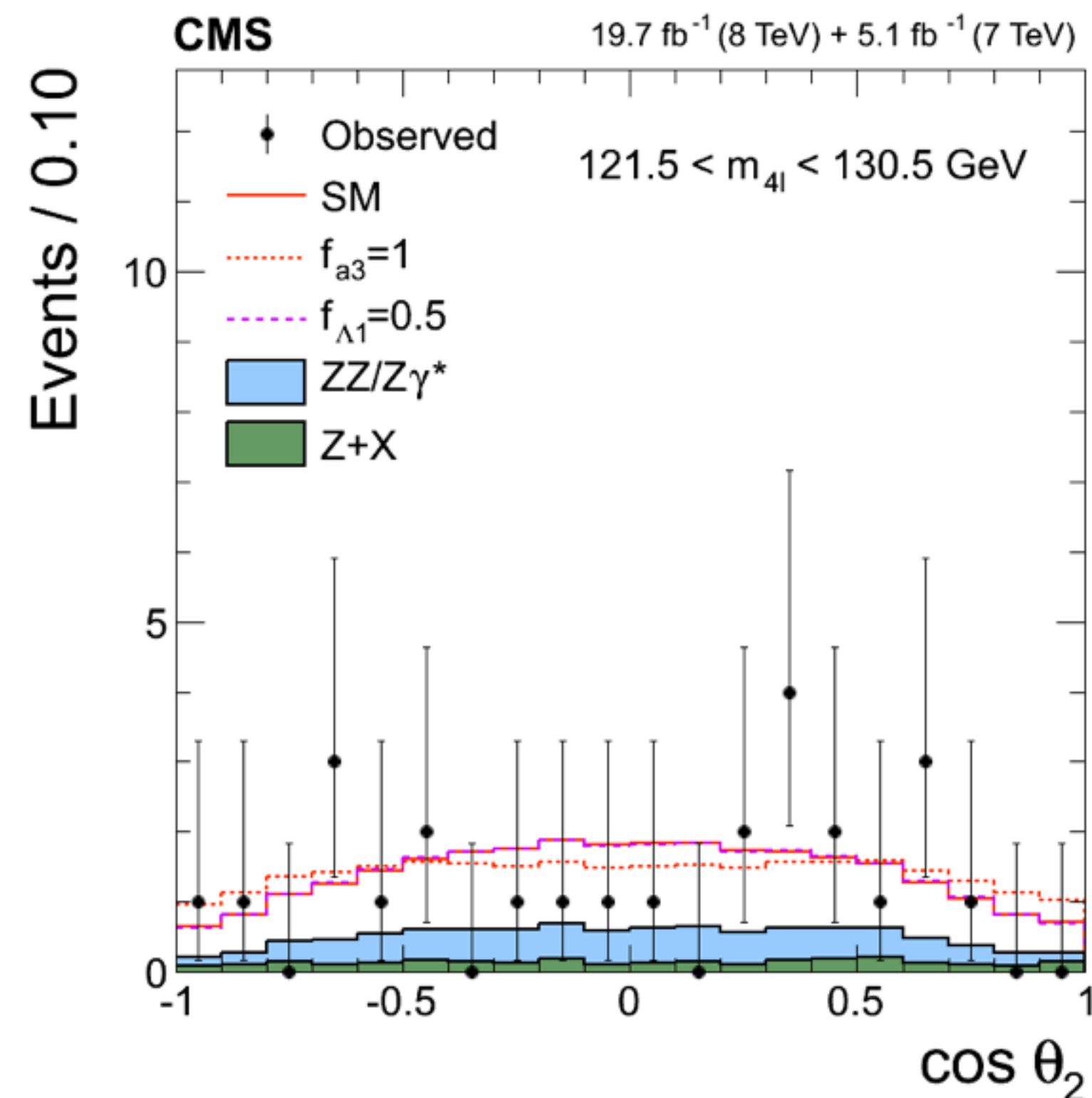
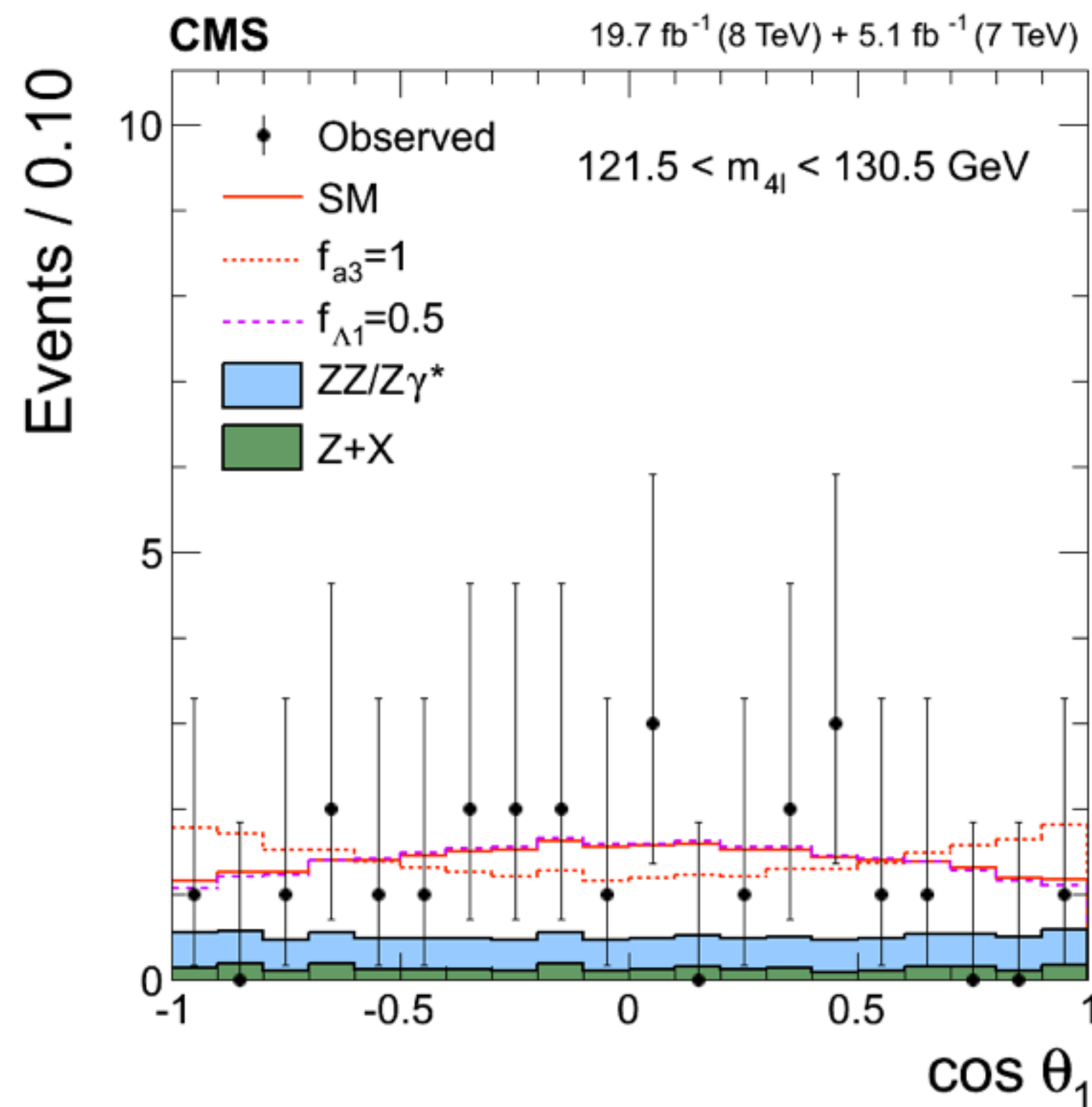
arXiv:1208.4018v3



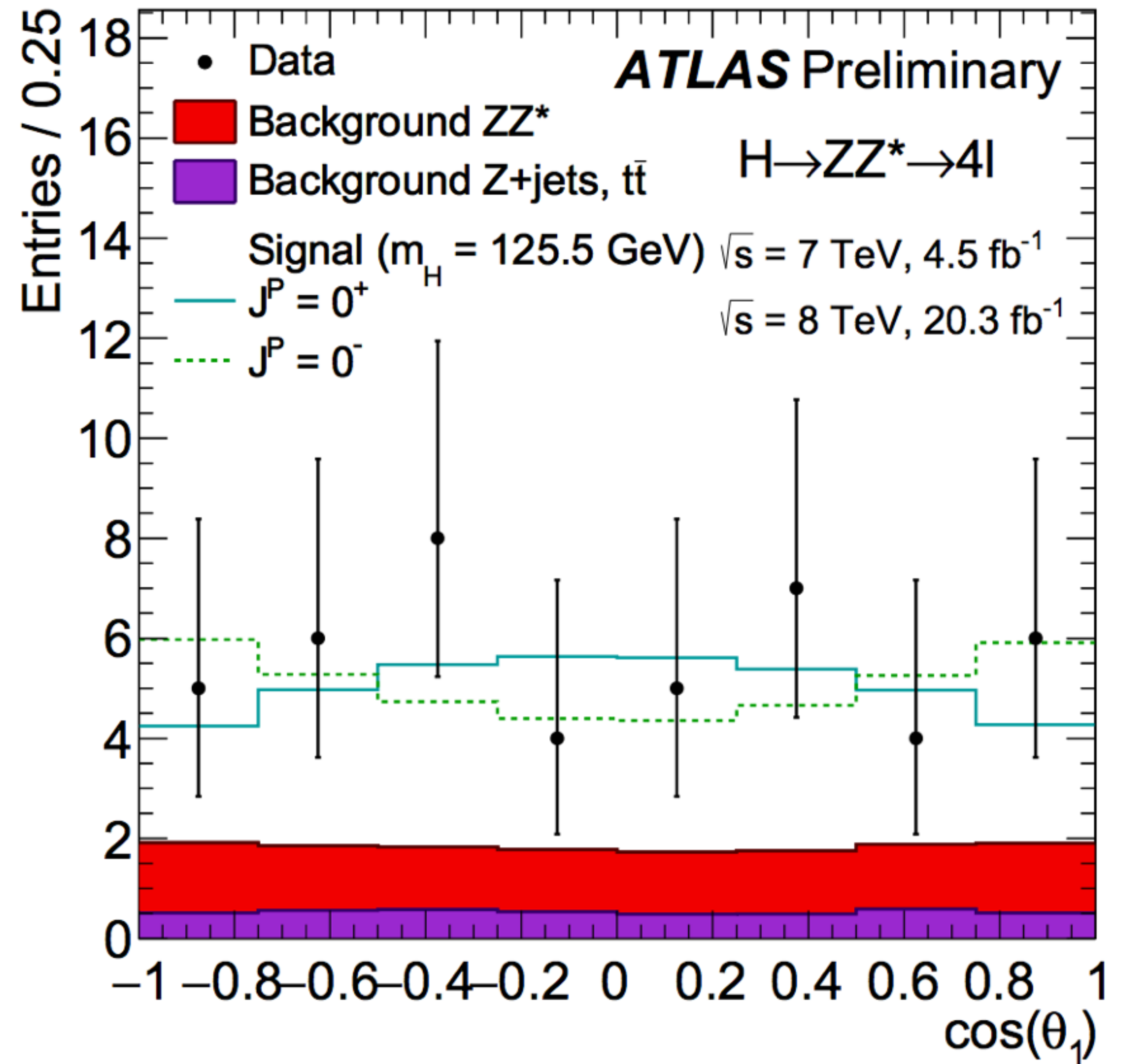
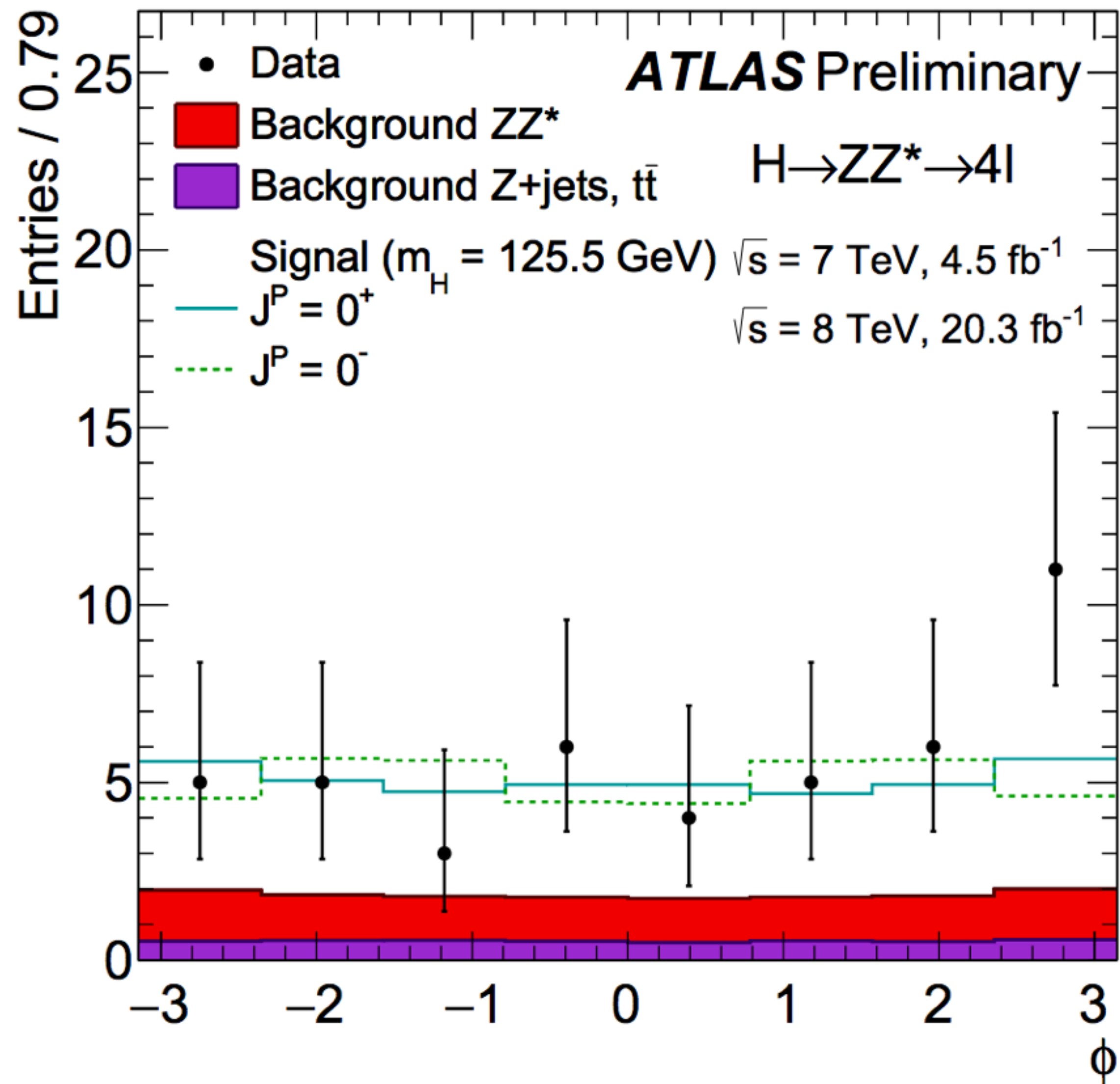
arXiv:1411.3441

Distributions *decay angles*

- Angles describing directions of leptons with respect to Z's
- These *are* sensitive to the CP properties!



ATLAS distributions



Multidimensional fits & modeling data

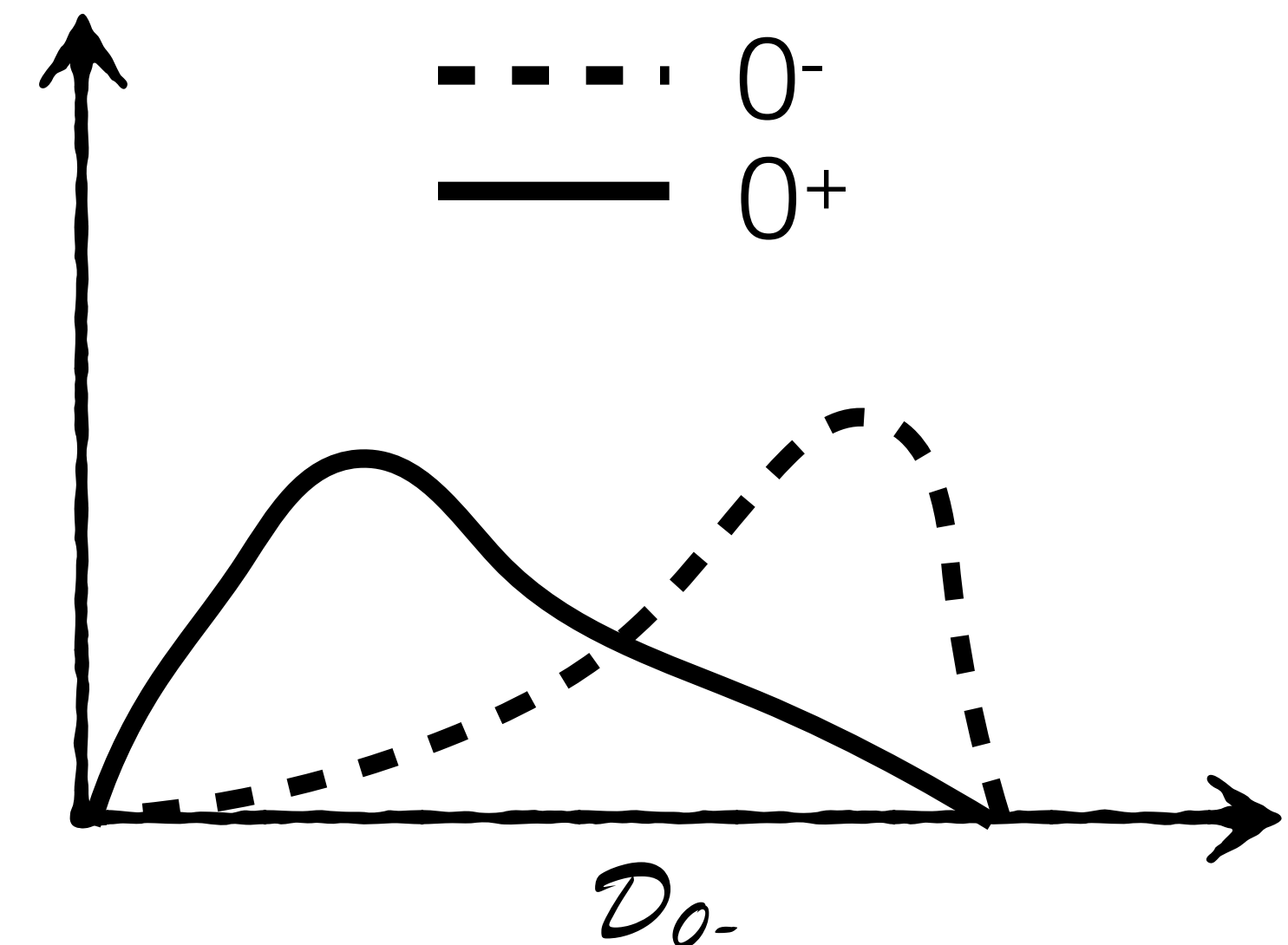
- modeling data with either
 - templates based on kinematic discriminants
 - discriminants based on LO ME or BDT
 - 3D templates describing various event categories
(one sensitive to signal/background,
one sensitive to CP eigenstates,
one sensitive to interference)
 - fully correlated multi-dim likelihood
(including approx. detector effects)
 - CMS: 8D
($m_{4\ell}, m_{1,2}, \cos\theta_{1,2}, \phi, \cos\theta^*, \phi_1$)
 - ATLAS: 9D
($p_T^{4\ell}, \eta_{4\ell}, m_{4\ell}, m_{1,2}, \cos\theta_{1,2}, \phi, \cos\theta^*$)
 - used to validate discriminant fits

e.g.:

$$|\mathcal{M}\mathcal{E}|^2 = |\mathcal{M}\mathcal{E}_{0+}|^2 + |\mathcal{M}\mathcal{E}_{0-}|^2 + 2\text{Re}(\mathcal{M}\mathcal{E}_{0+}^* \mathcal{M}\mathcal{E}_{0-})$$

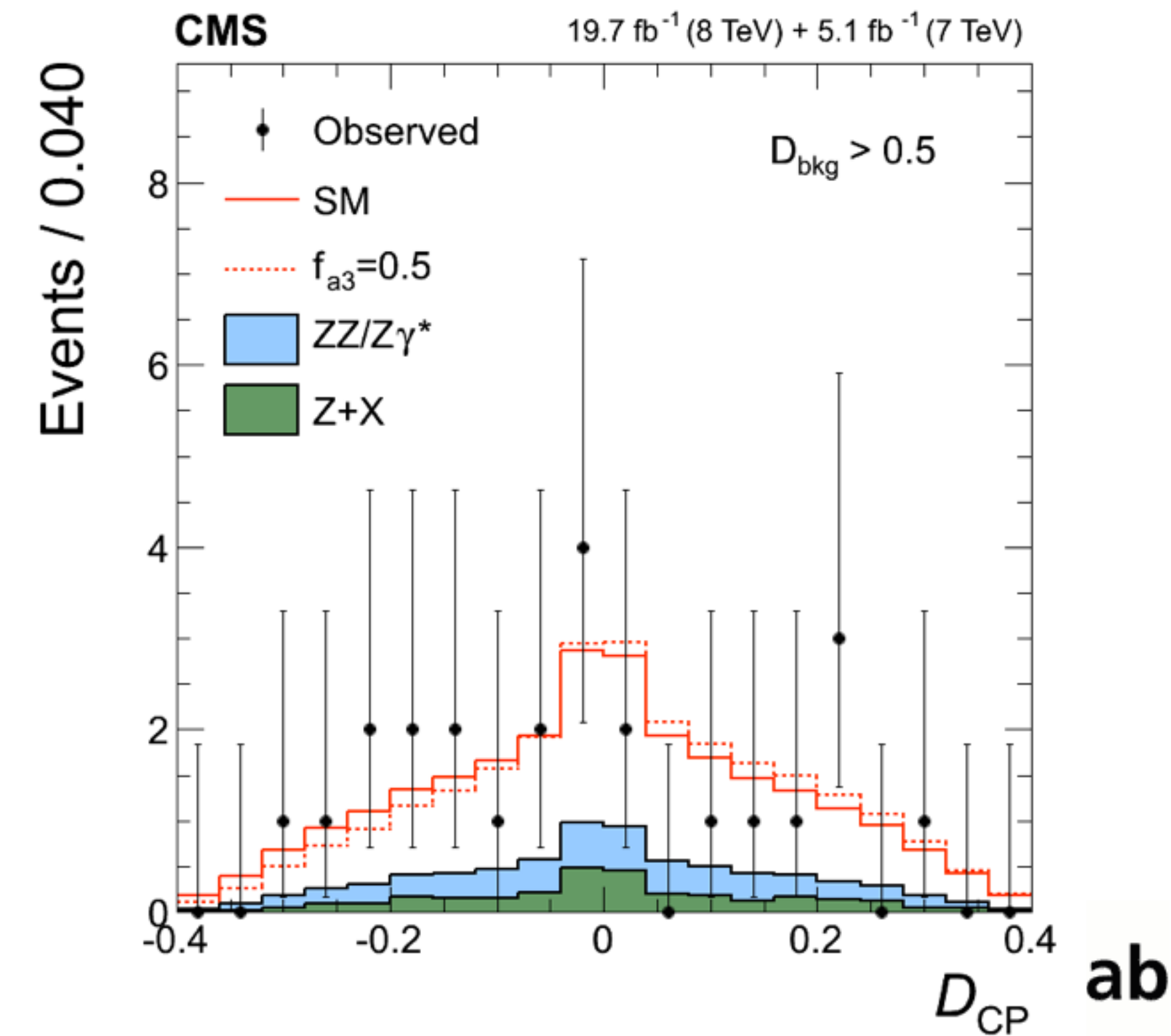
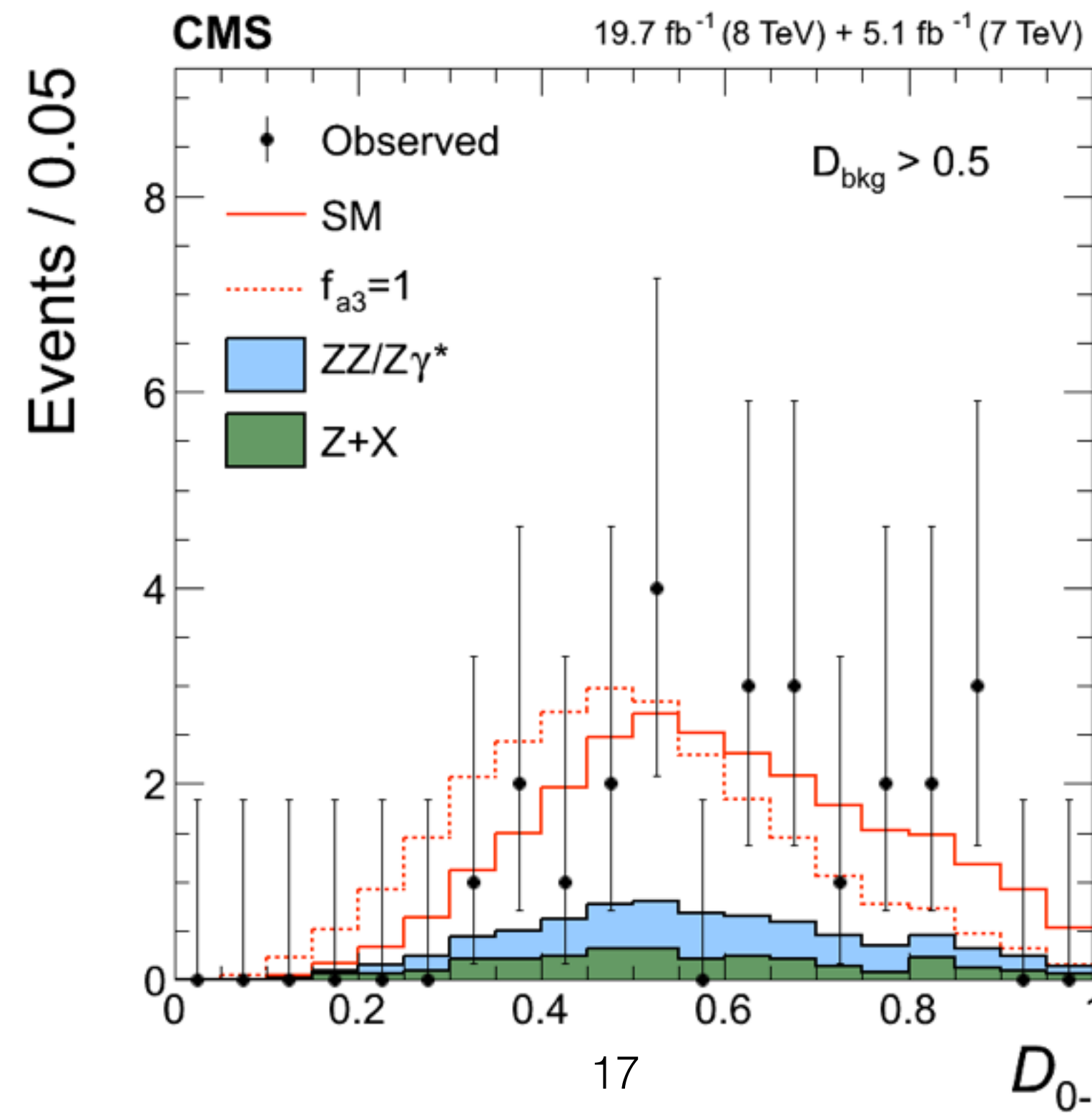
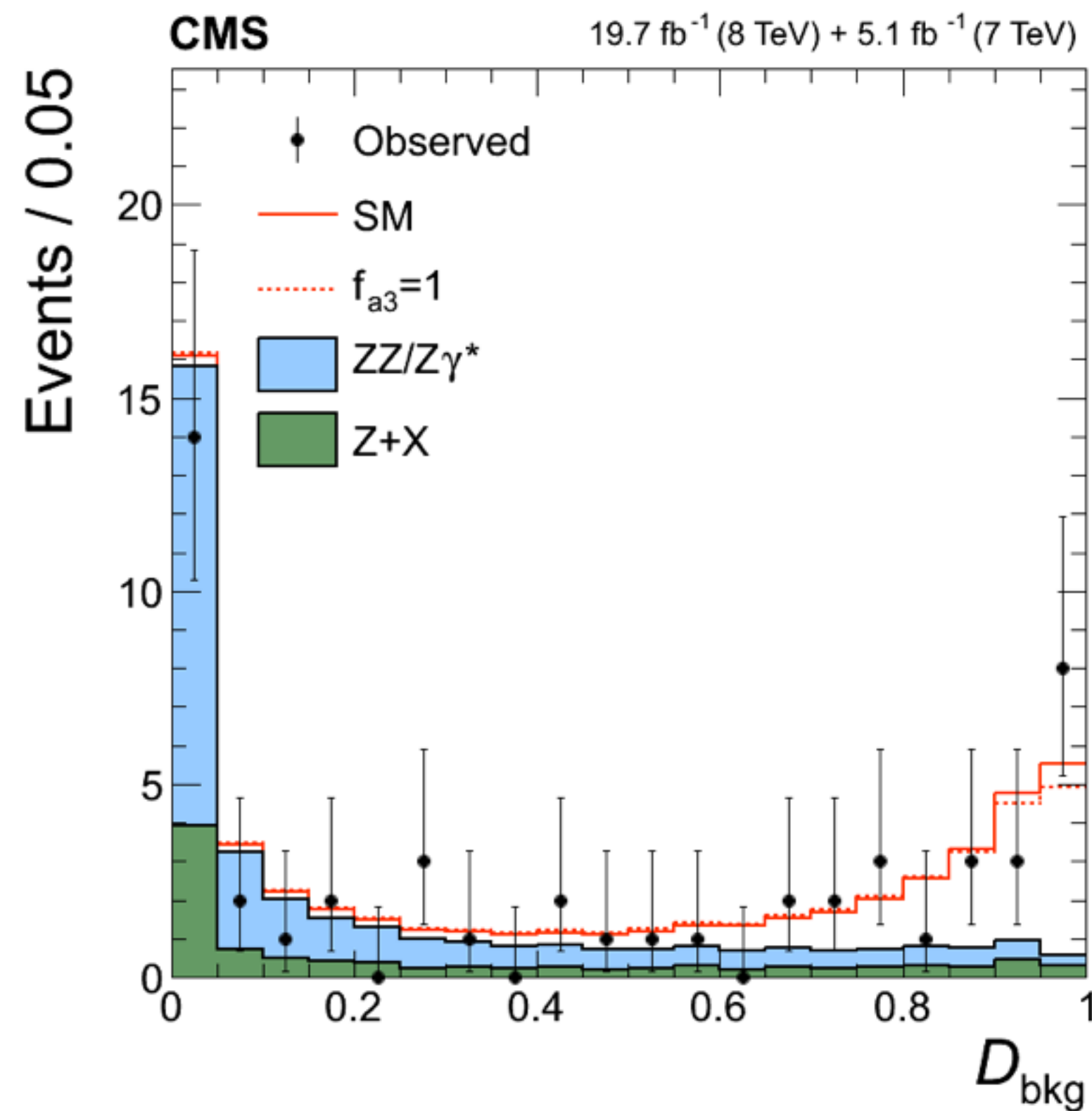
$$\mathcal{D}_{0-} = |\mathcal{M}\mathcal{E}_{0-}|^2 / [|\mathcal{M}\mathcal{E}_{0+}|^2 + |\mathcal{M}\mathcal{E}_{0-}|^2]$$

$$\mathcal{D}_{\text{int}} = 2\text{Re}(\mathcal{M}\mathcal{E}_{0+}^* \mathcal{M}\mathcal{E}_{0-}) / [|\mathcal{M}\mathcal{E}_{0+}|^2 + |\mathcal{M}\mathcal{E}_{0-}|^2]$$



Distributions discriminants

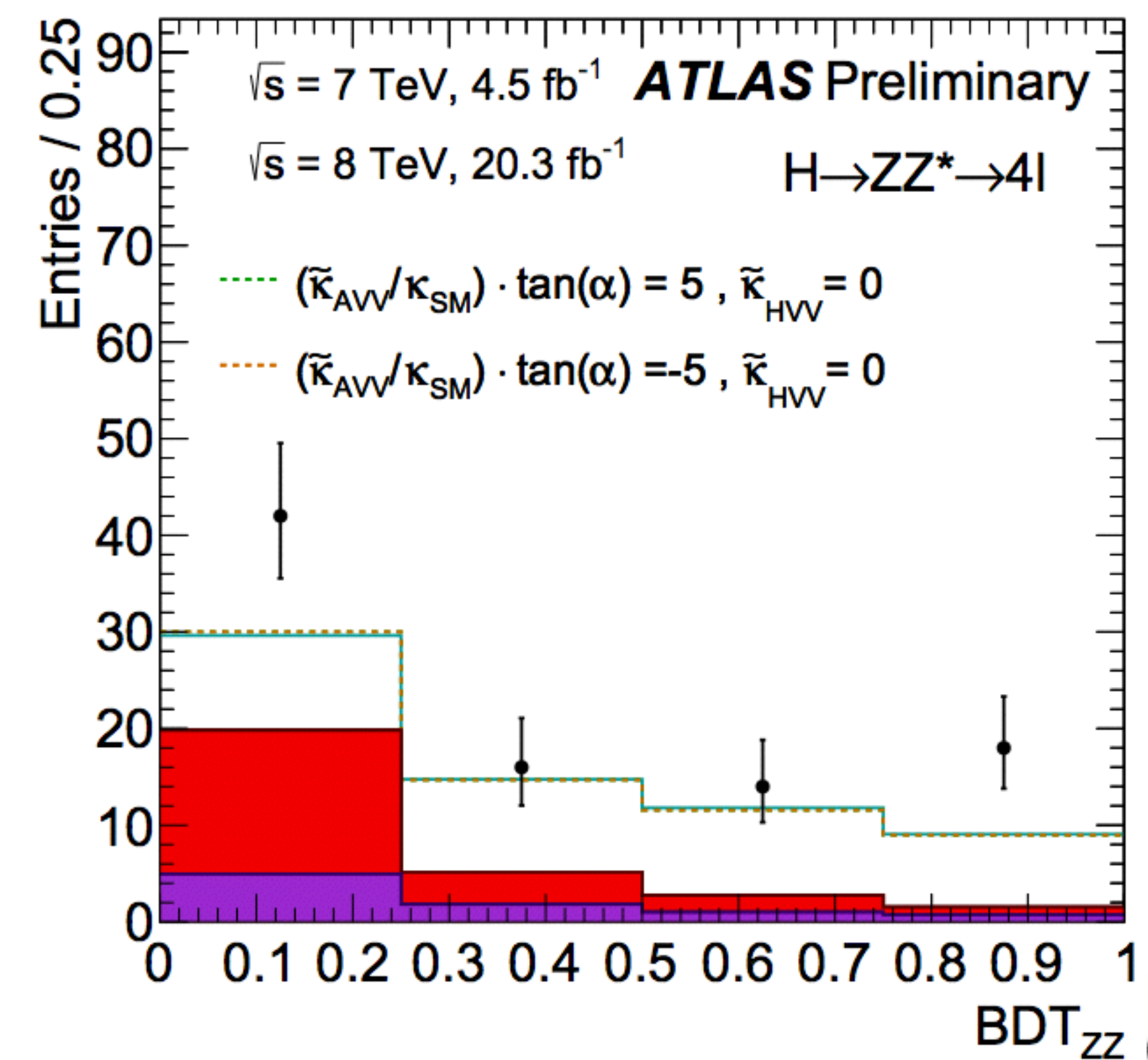
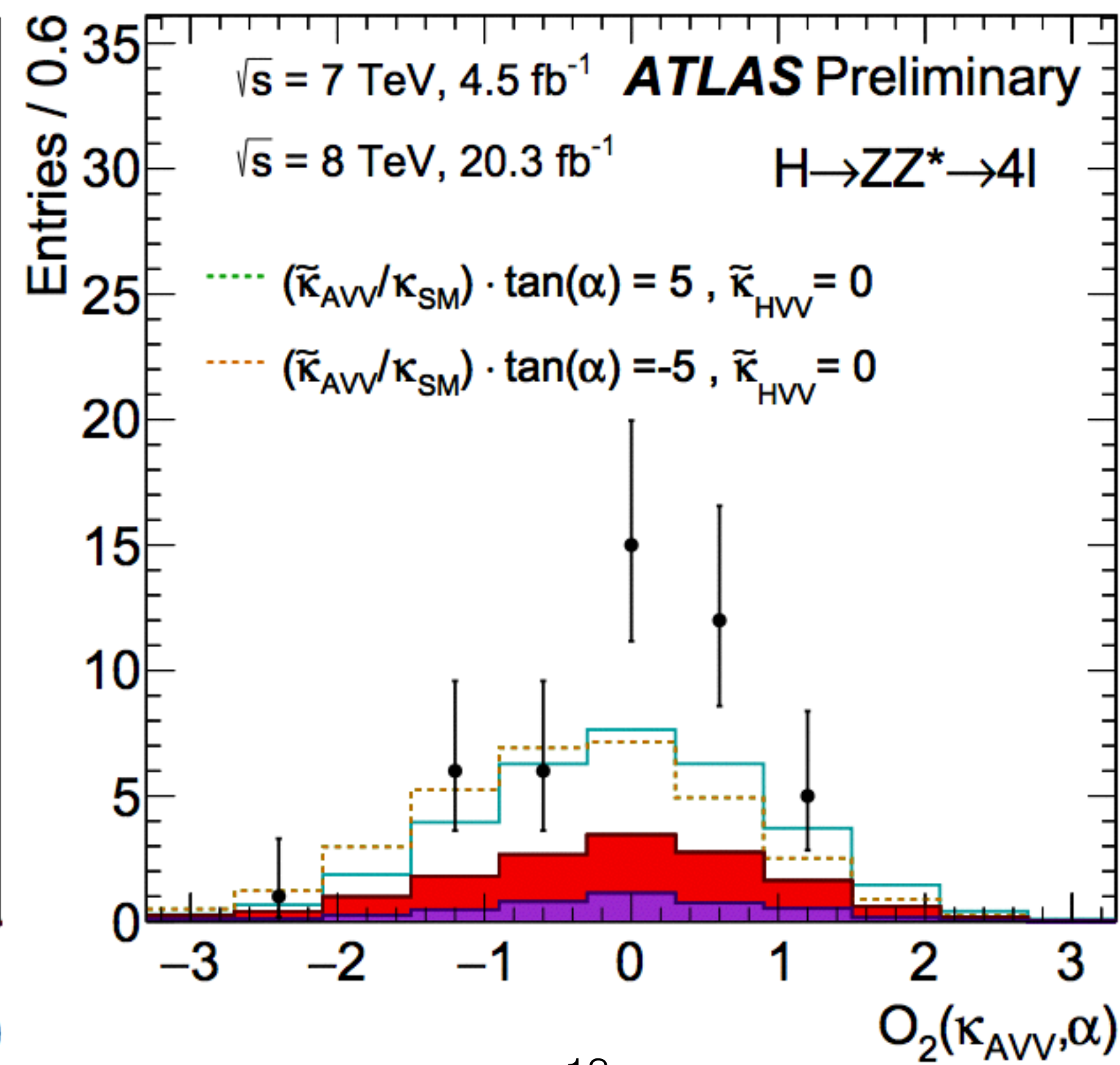
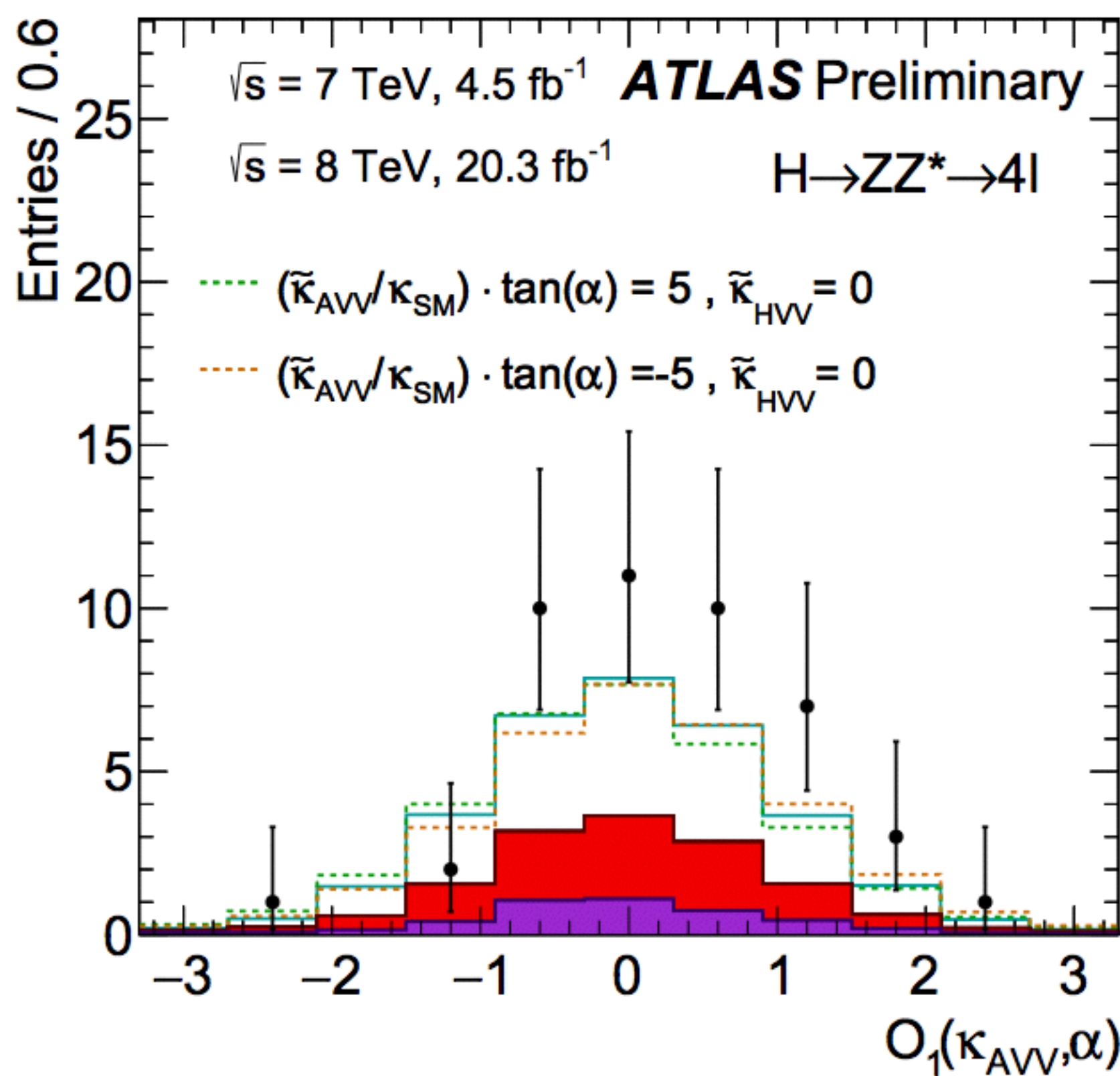
- D_{bkg} used to distinguish signal from background, dominant contribution from masses
- D_{0-} used to distinguish $0+$ from $0-$ neglecting kinematics from interference
- D_{CP} used to enhance interference effects, become important at small values of f_{a3}



ATLAS discriminants

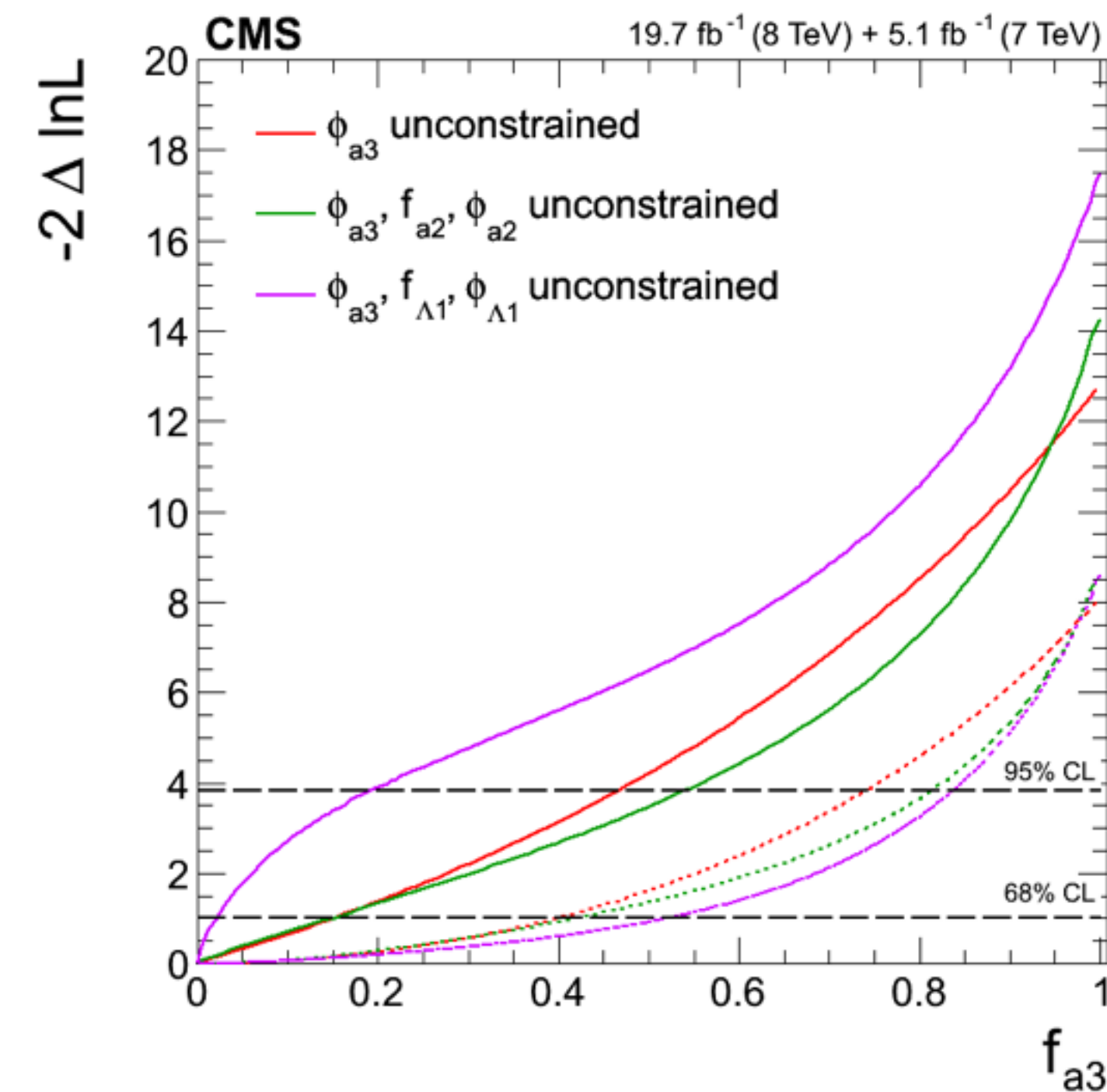
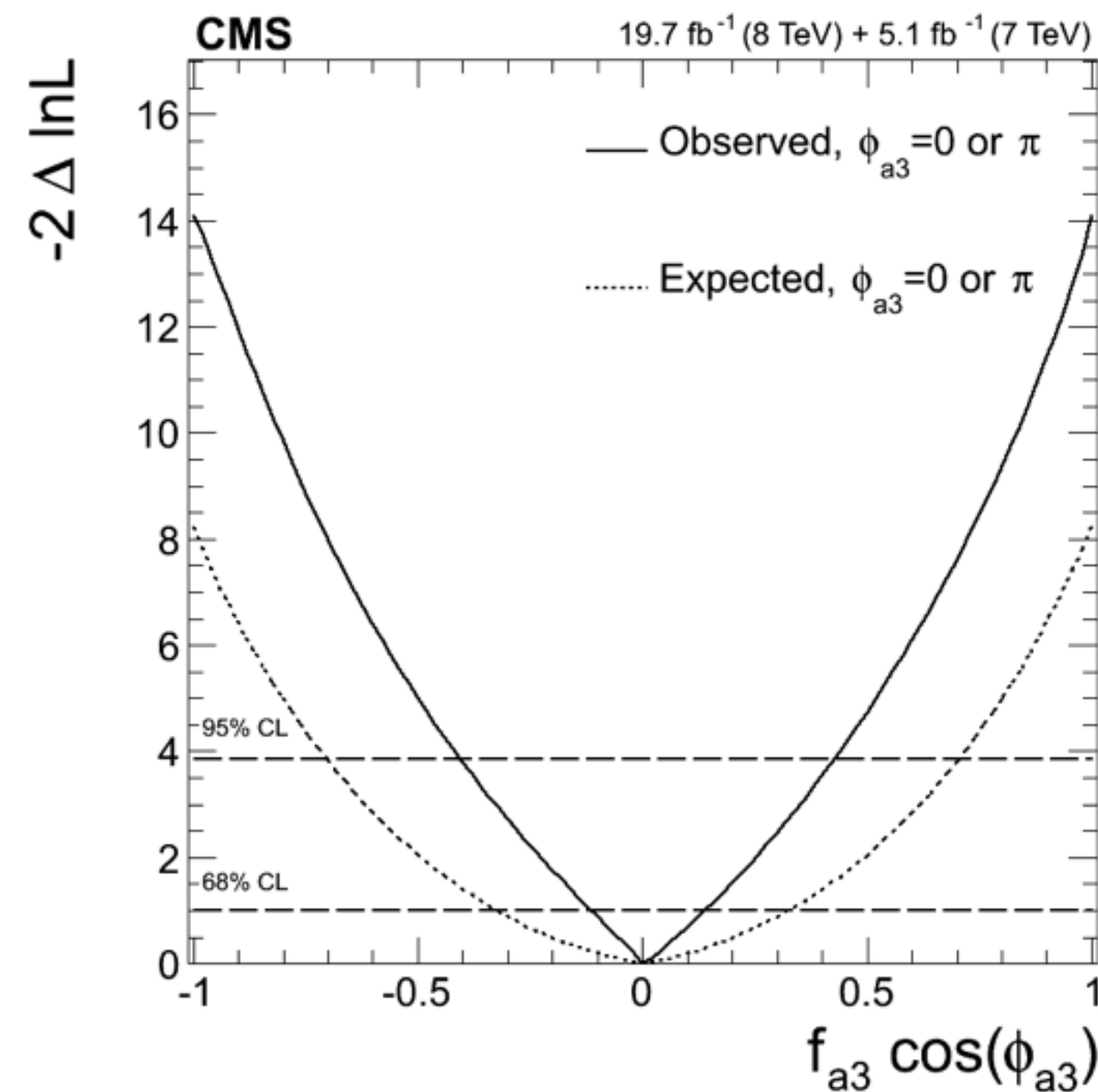
- Three discriminants are qualitatively similar to those used for the CMS analysis
- BDT_{ZZ} trained for signal versus backgrounds
(inputs: $\eta_{4\ell}$, $p_{T^{4\ell}}$, $m_{4\ell}$, $\cos(\theta^*)$, ϕ_1)

- Data
 - Background ZZ^*
 - Background Z +jets, $t\bar{t}$
 - SM: $\kappa_{\text{SM}} = 1$, $\tilde{\kappa}_{\text{HVV}} = 0$, $\tilde{\kappa}_{\text{AVV}} = 0$, $\alpha = 0$
- Signal ($m_H = 125.5$ GeV)



1D likelihood scans CMS

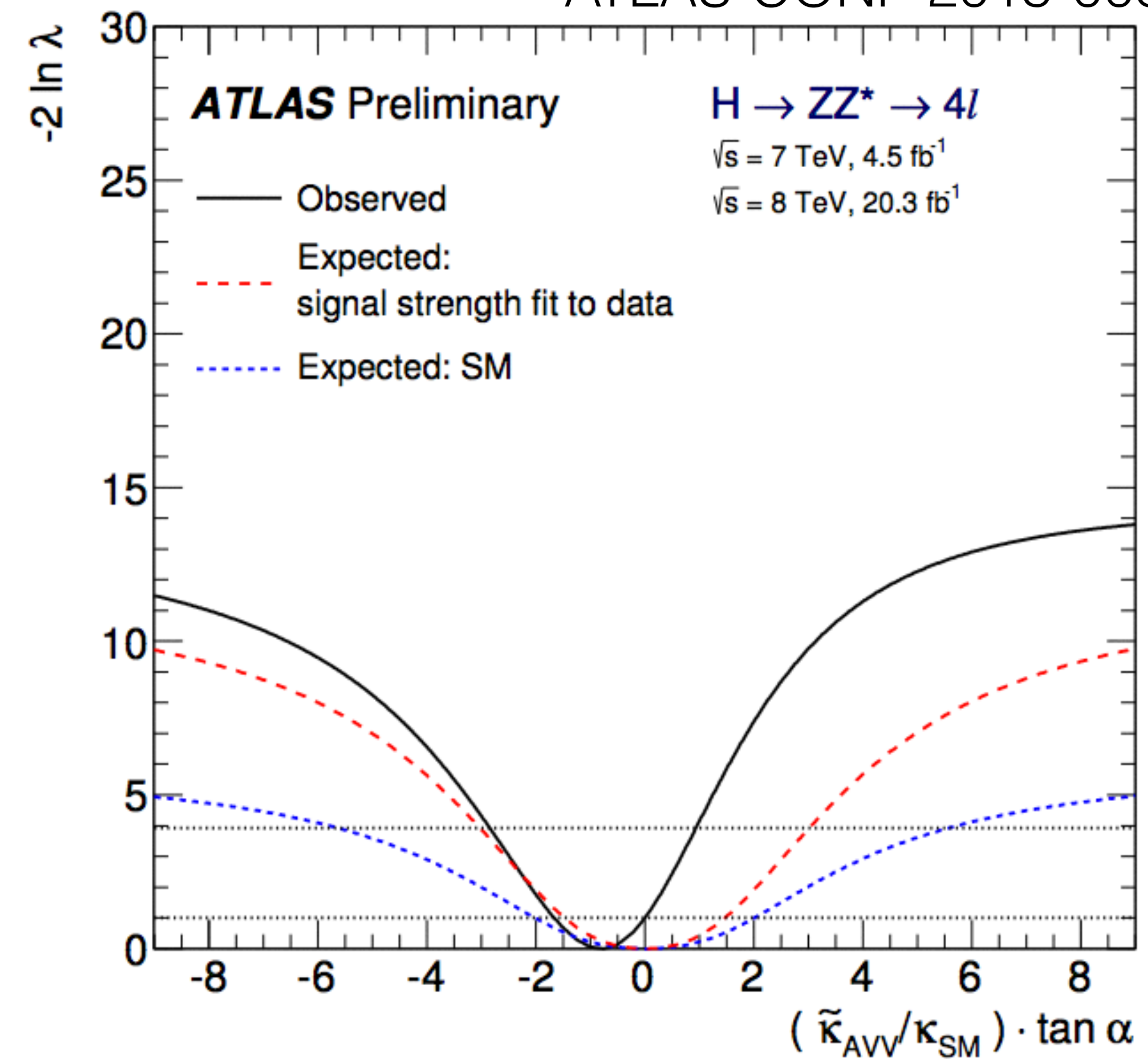
- log-likelihood scans for f_{a3} show consistency with the SM
- f_{a3} is a direct probe of CP-violating interactions - 95% C.L. upper limits set as low as $f_{a3} < 0.2$ (depending on the assumptions made)



Parameter	Observed	Expected
$f_{a3} \cos(\phi_{a3})$	$0.00^{+0.14}_{-0.11} [-0.40, 0.43]$	$0.00^{+0.33}_{-0.33} [-0.70, 0.70]$
a_3/a_1	$[-2.05, 2.19]$	$[-3.85, 3.85]$

1D likelihood scans ATLAS

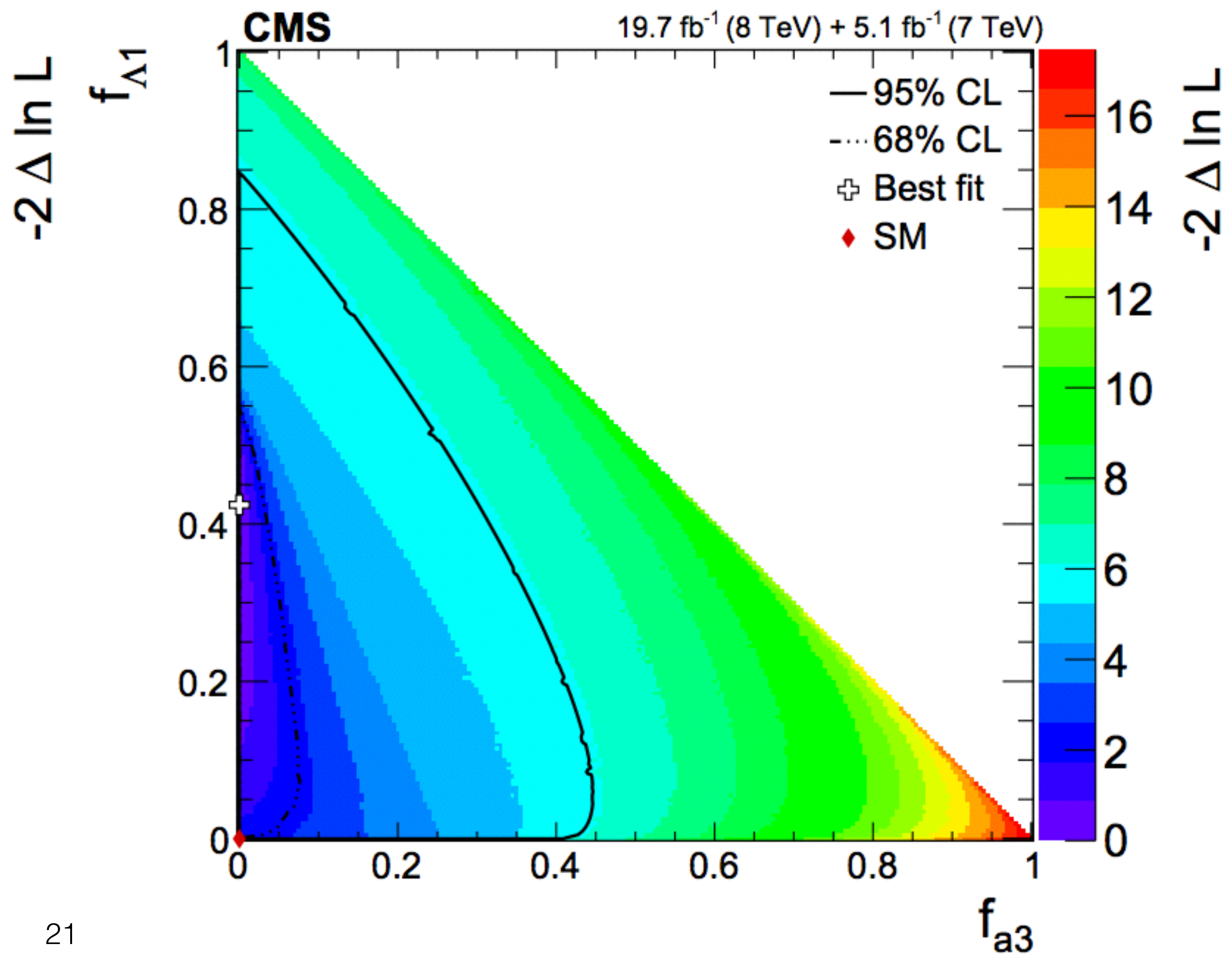
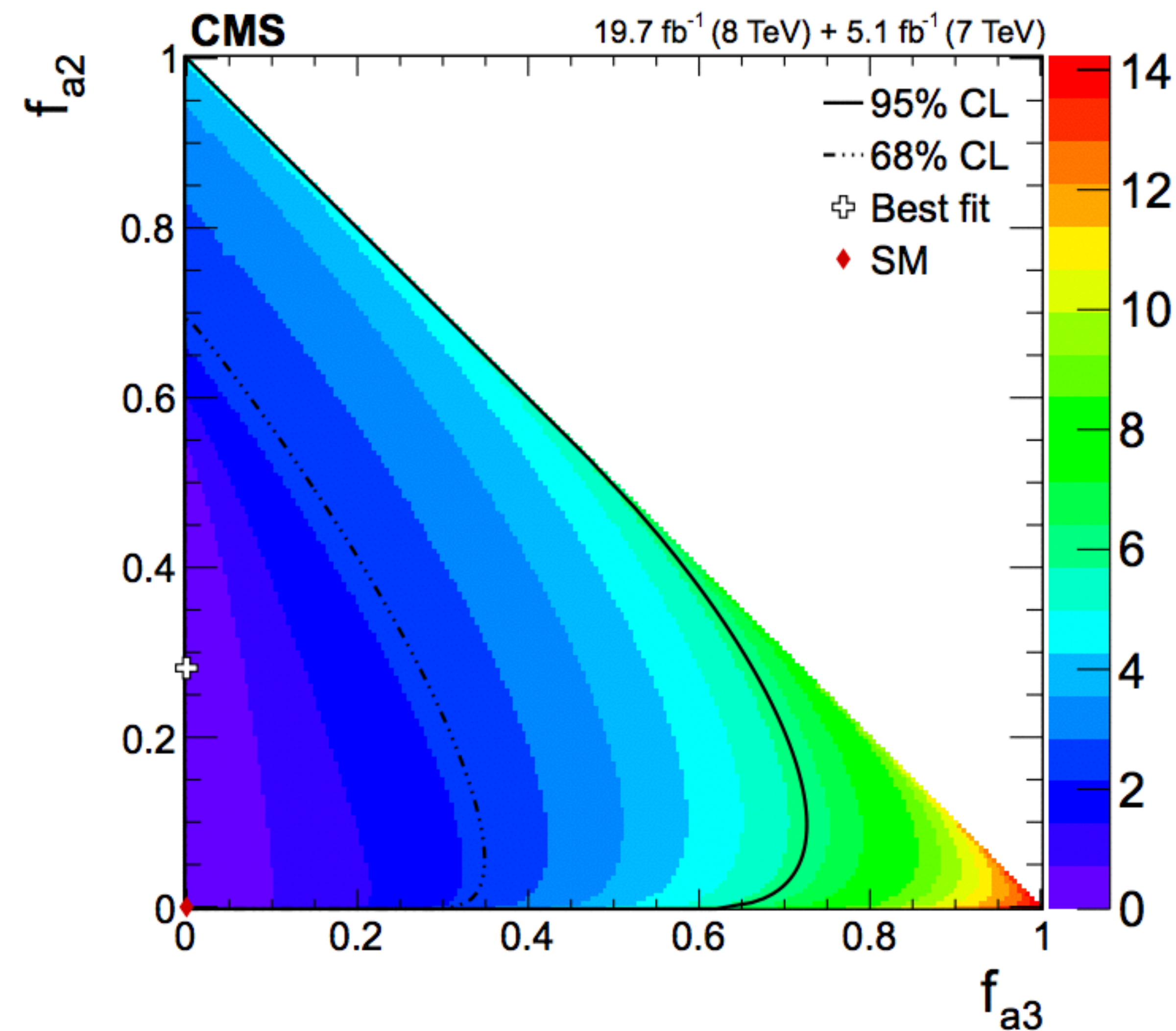
- Results consistent with SM expectations ($\kappa_{AVV}/\kappa_{SM}=0$)
 - small deviation from zero observed, only 1σ effect
- Results similar to CMS



Coupling ratio $H \rightarrow ZZ^* \rightarrow 4\ell$	Best fit value		95% CL Exclusion Regions	
	Expected	Observed	Expected	Observed
$\tilde{\kappa}_{HV V} / \kappa_{SM}$	0.0	-0.2	$(-\infty, -0.75] \cup [6.95, \infty)$	$(-\infty, -0.75] \cup [2.45, \infty)$
$(\tilde{\kappa}_{AV V} / \kappa_{SM}) \cdot \tan \alpha$	0.0	-0.8	$(-\infty, -2.95] \cup [2.95, \infty)$	$(-\infty, -2.85] \cup [0.95, \infty)$

2D likelihood scans **CMS**

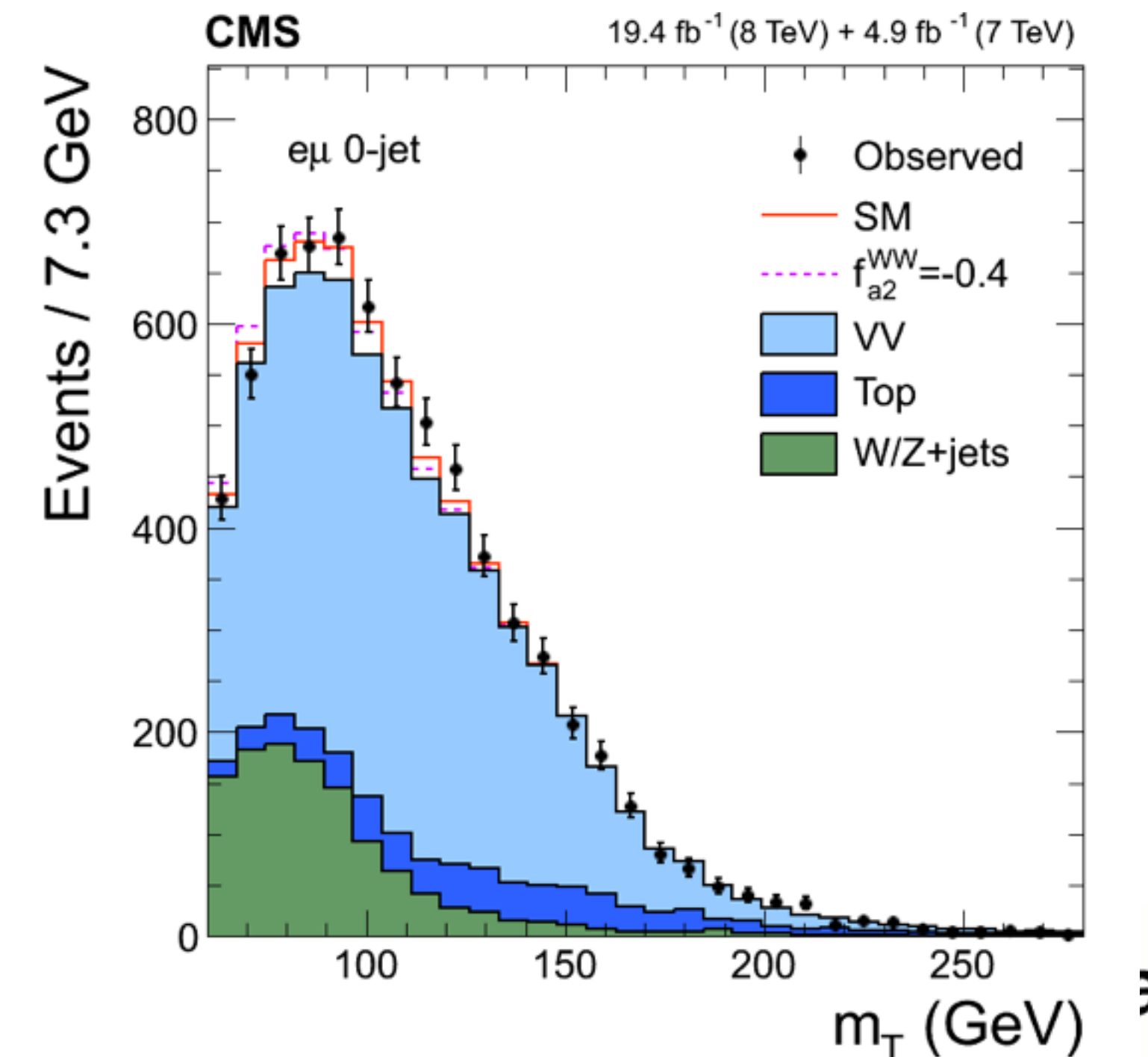
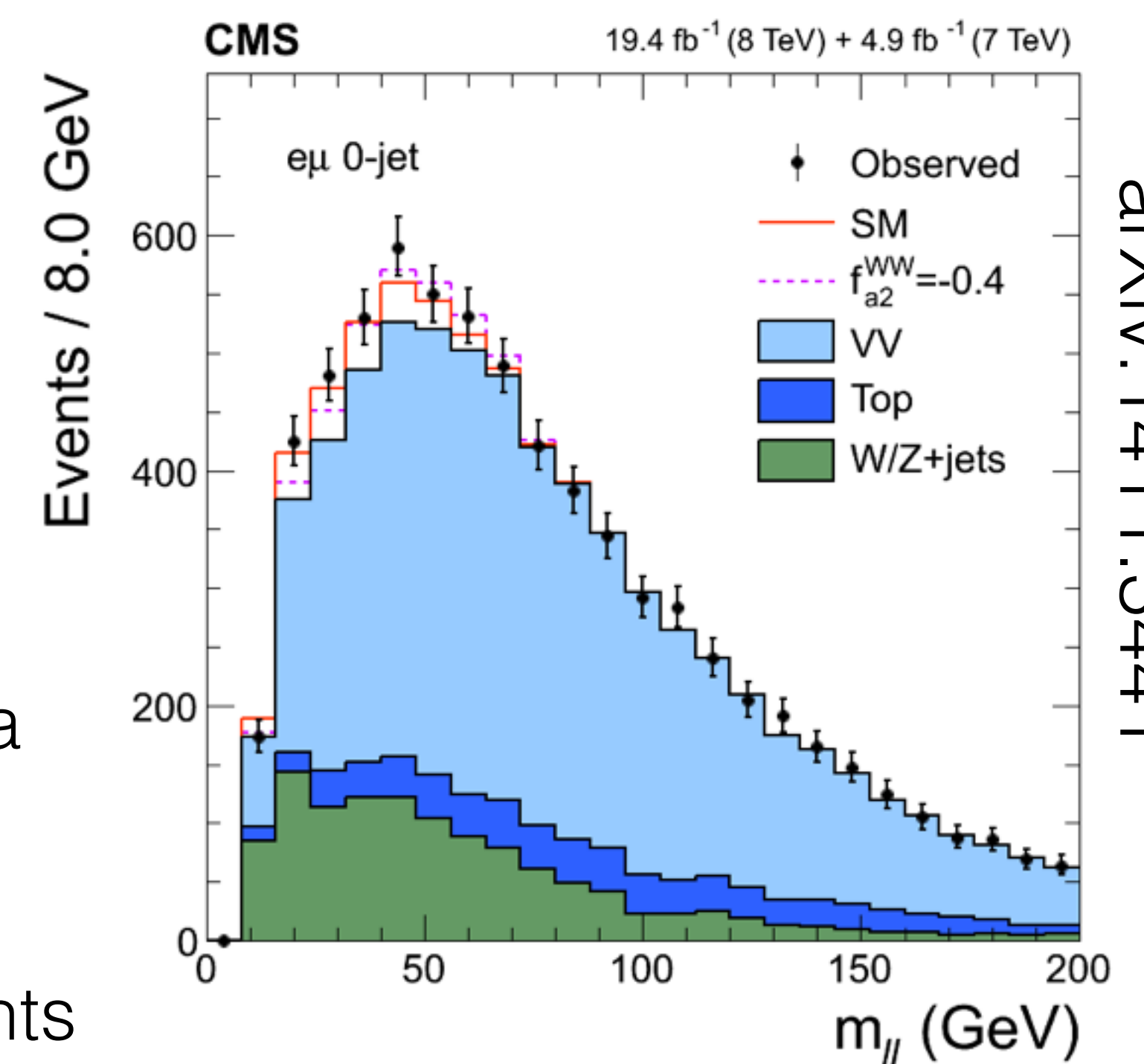
- For completeness: 2D scan of f_{a2} , $f_{\Lambda 1}$ versus f_{a3} are shown



WW analysis

- opposite sign e- μ events selected
- Events classified based on the exclusive numbers of jets (0 or 1 jet)
- Main backgrounds, WW, ttbar/tW, and Z/W+jets are extrapolated from data control regions
- CMS: 2D distribution of m_{ll} and m_T distributions are used to describe events
ATLAS: 2D distributions of BDT discriminants are used

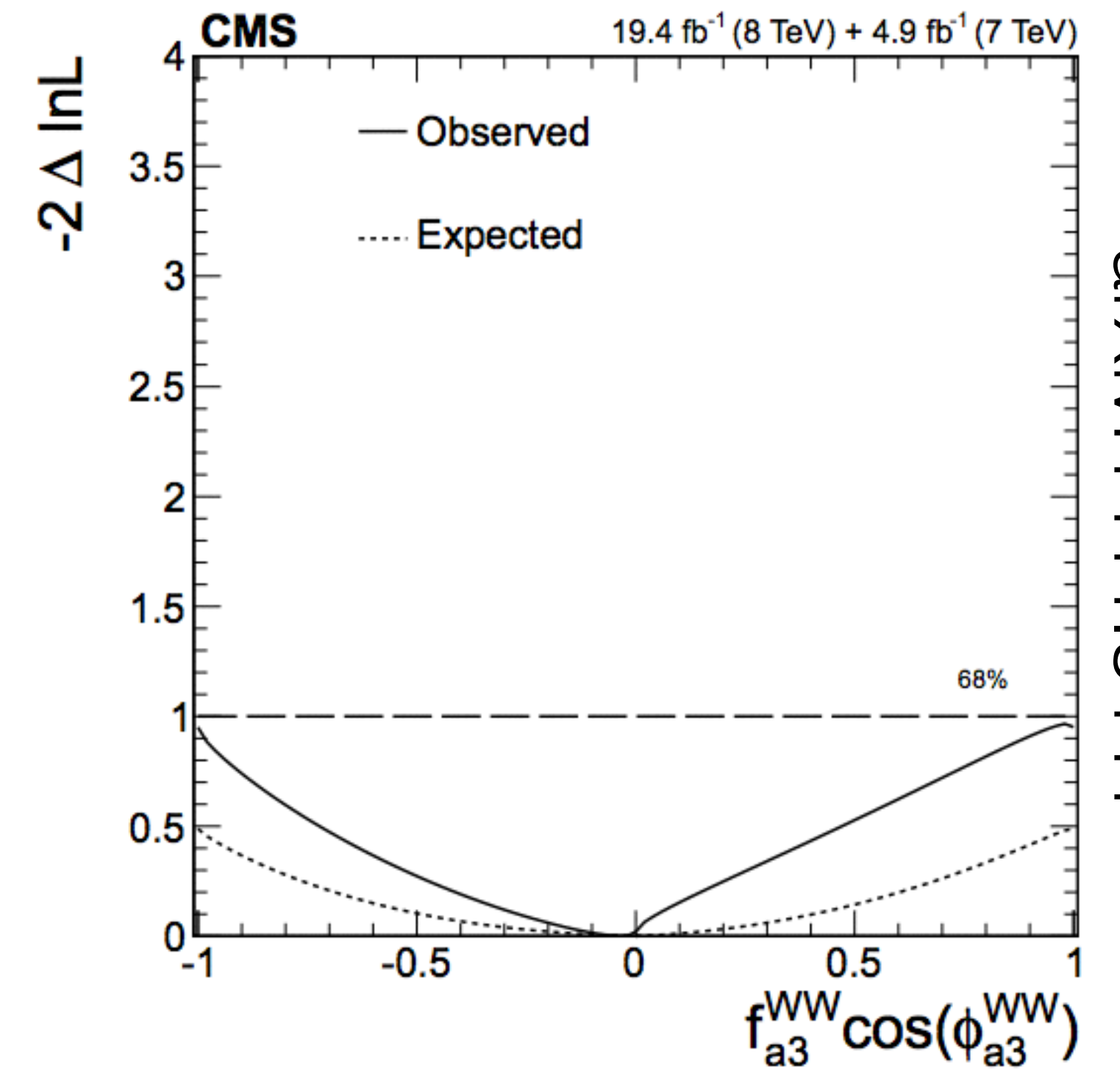
Channel	0-jet		1-jet	
Energy	7 TeV	8 TeV	7 TeV	8 TeV
WW	861 ± 12	4185 ± 63	249.9 ± 4.0	1268 ± 21
WZ + ZZ + Z/ γ^*	22.7 ± 1.2	178.3 ± 9.5	26.4 ± 1.4	193 ± 11
$t\bar{t} + tW$	91 ± 20	500 ± 96	226 ± 14	1443 ± 46
W+jets	150 ± 39	620 ± 160	60 ± 16	283 ± 72
W $\gamma^{(*)}$	68 ± 20	282 ± 76	10.1 ± 2.8	55 ± 14
Background	1193 ± 50	5760 ± 210	573 ± 22	3242 ± 90
Signal gg \rightarrow H	50 ± 10	227 ± 46	17.1 ± 5.5	88 ± 28
Signal VBF+VH	0.44 ± 0.03	10.27 ± 0.41	2.09 ± 0.12	19.83 ± 0.81
Observed	1207	5747	589	3281



1D likelihood scans **WW**

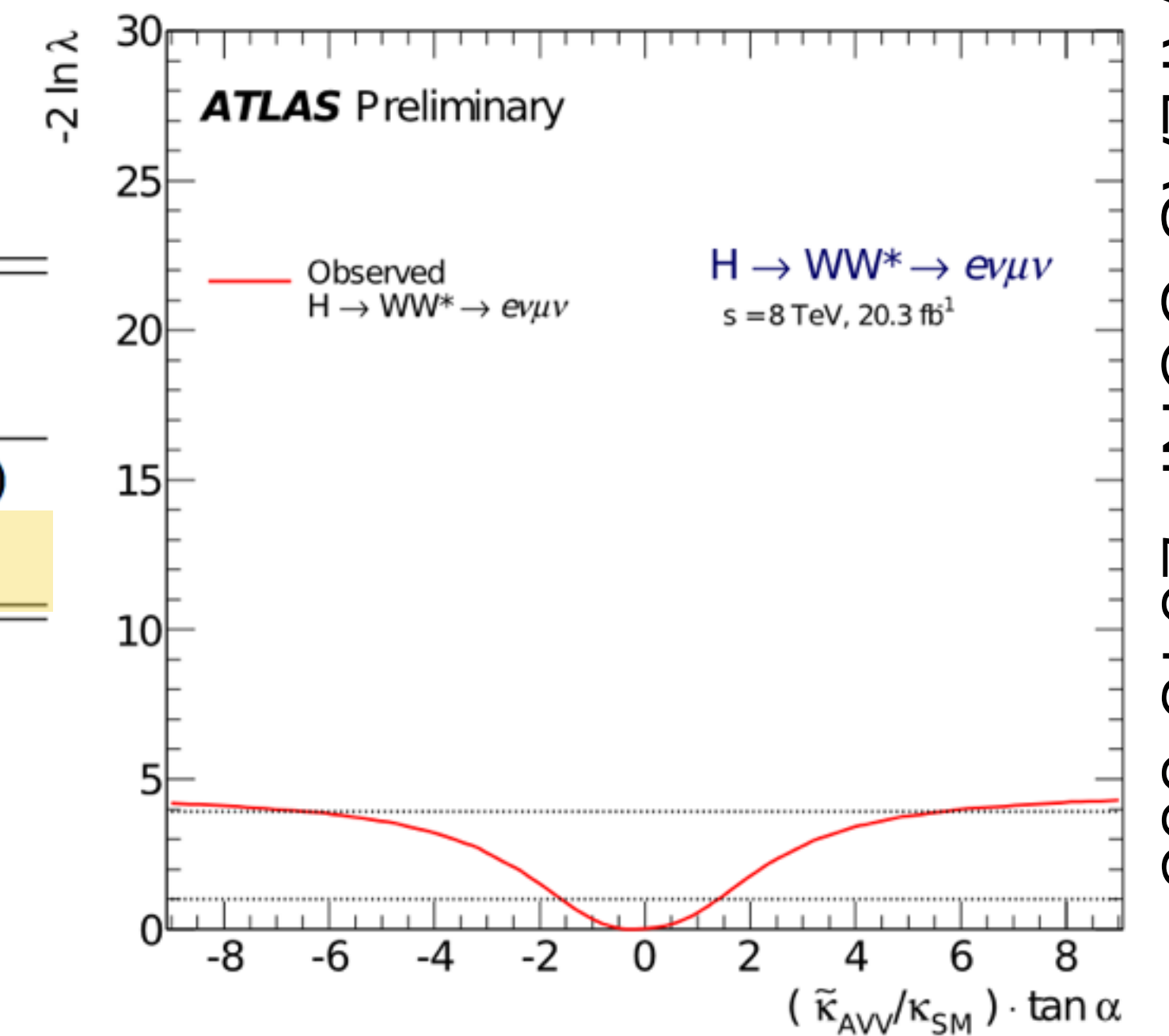
CMS

Parameter	Observed	Expected
$f_{\Lambda 1}^{WW} \cos(\phi_{\Lambda 1}^{WW})$	$0.21^{+0.18}_{-1.21} [-1.00, 1.00]$	$0.00^{+0.34}_{-1.00} [-1.00, 0.41]$ $\cup [0.49, 1.00]$
$f_{a2}^{WW} \cos(\phi_{a2}^{WW})$	$-0.02^{+1.02}_{-0.16} [-1.00, -0.54]$ $\cup [-0.29, 1.00]$	$0.00^{+1.00}_{-0.12} [-1.00, -0.58]$ $\cup [-0.22, 1.00]$
$f_{a3_}^{WW} \cos(\phi_{a3_}^{WW})$	$-0.03^{+1.03}_{-0.97} [-1.00, 1.00]$	$0.00^{+1.00}_{-1.00} [-1.00, 1.00]$



ATLAS

Coupling ratio	Best fit value		95% CL Exclusion Regions	
	Expected	Observed	Expected	Observed
$H \rightarrow WW^* \rightarrow e\nu\mu\nu$				
$\tilde{\kappa}_{HVV}/\kappa_{SM}$	0.0	-1.3	$[-1.2, -0.7]$	$(-\infty, -2.2] \cup [-1, -0.85] \cup [0.4, \infty)$
$(\tilde{\kappa}_{AVV}/\kappa_{SM}) \cdot \tan \alpha$	0.0	-0.2	-	$(-\infty, -6] \cup [5, \infty)$



arXiv:1411.3441

ATLAS-CONF-2015-008

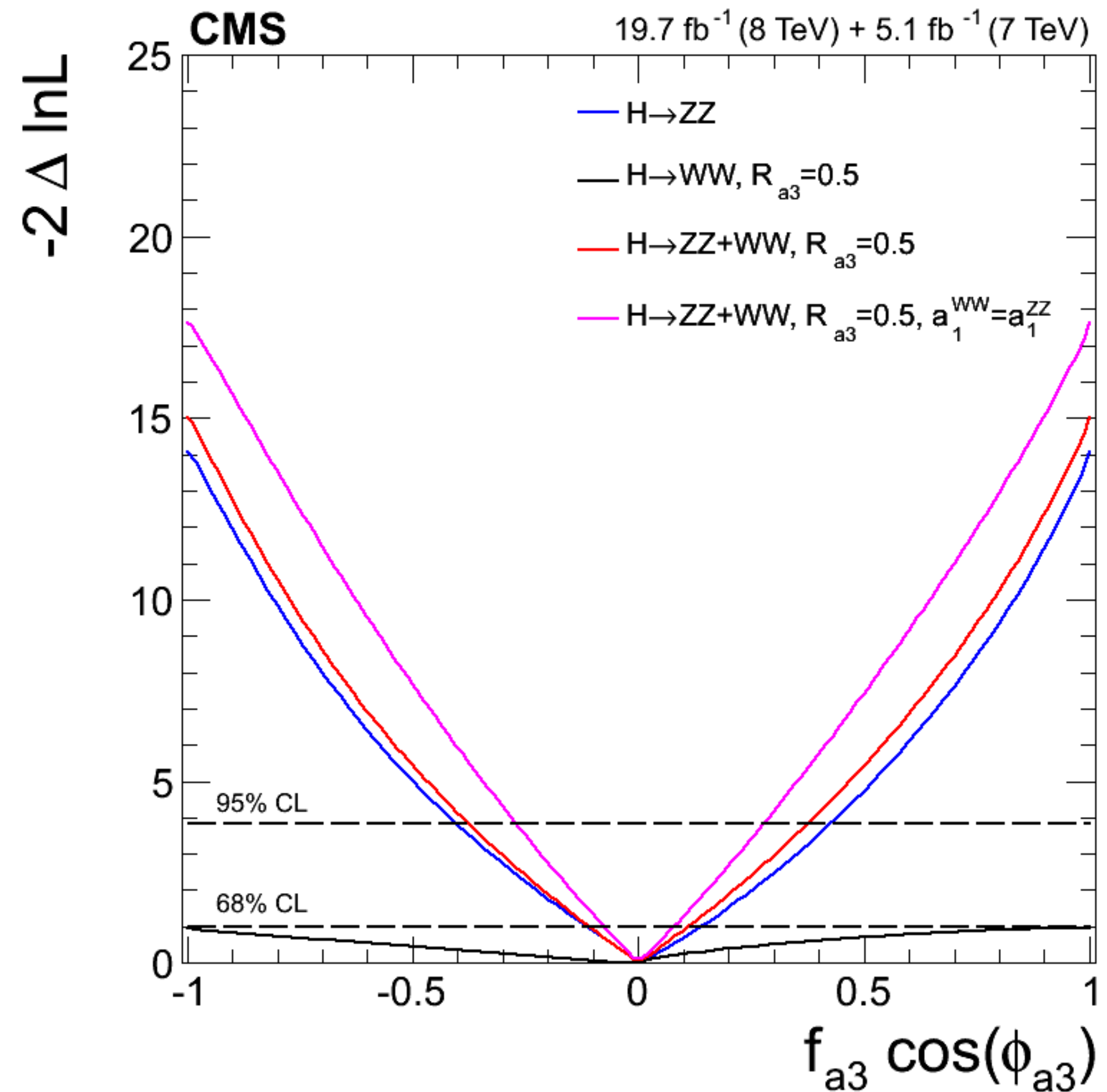
b

Combined fits (ZZ&WW) CMS

- The combination of ZZ & WW channels involves twice as many degrees of freedom
- We can constrain the size of the SM couplings (a_1) to be the same — custodial symmetry
 - fits are done both **with** and **without** this assumption
- The relative strange of a_3/a_1 can also be fixed

$$r_{ai} = \frac{a_i^{WW} / a_1^{WW}}{a_i / a_1}, \text{ or } R_{ai} = \frac{r_{ai} |r_{ai}|}{1 + r_{ai}^2}.$$

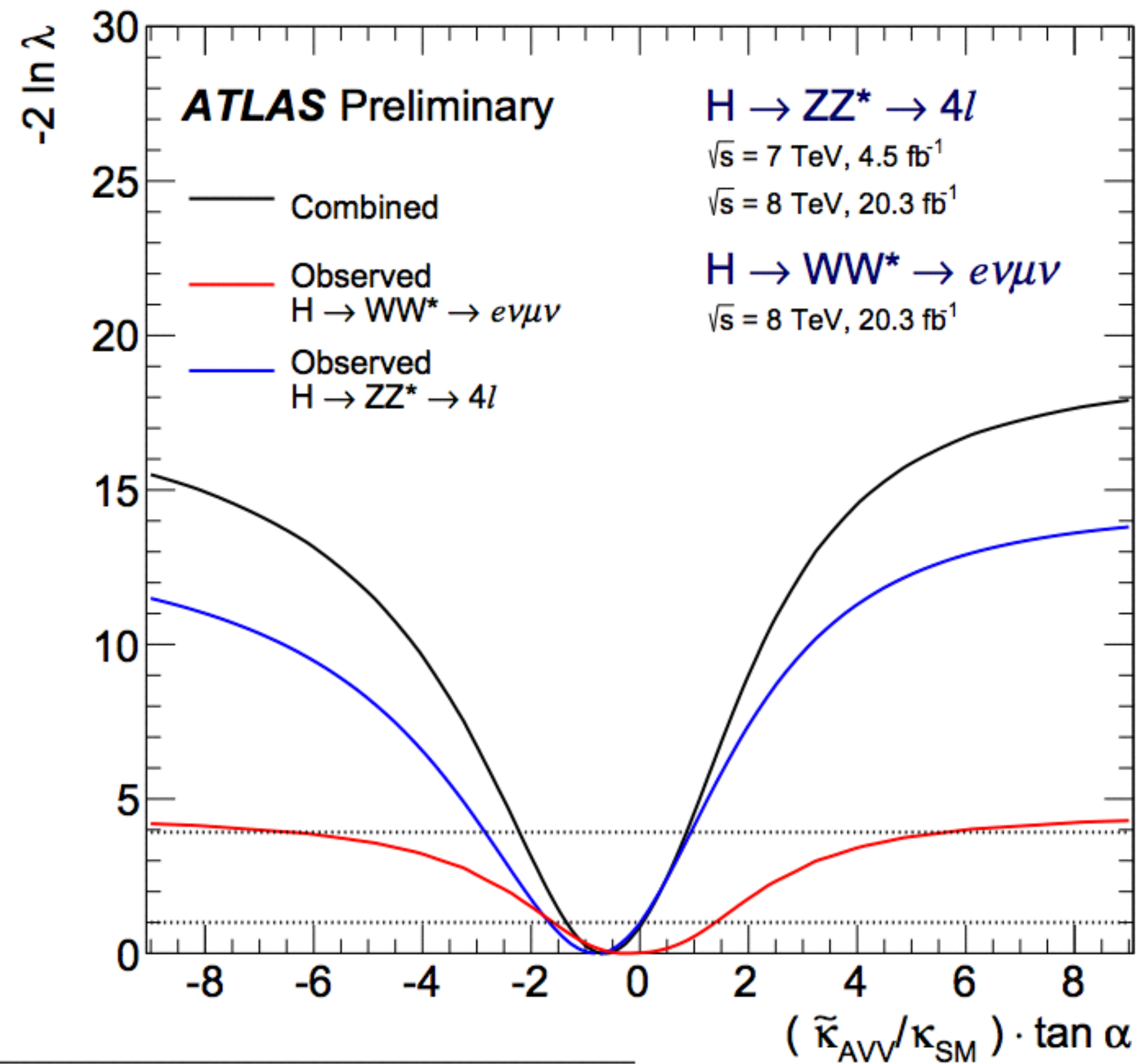
- custodial symmetry & $R_{a3} = 0.5$ implies that $a_1^{WW} = a_1$ & $a_3^{WW} = a_3$



Parameter	Observed	Expected
$f_{\Lambda 1} \cos(\phi_{\Lambda 1})$	$0.21^{+0.11}_{-0.17} [-0.42, 0.38]$	$0.00^{+0.15}_{-0.80} [-1, 0.27] \cup [0.95, 1]$
$f_{a2} \cos(\phi_{a2})$	$-0.01^{+0.02}_{-0.05} [-0.11, 0.17]$	$0.00^{+0.08}_{-0.03} [-0.07, 0.51]$
$f_{a3} \cos(\phi_{a3})$	$0.00^{+0.08}_{-0.08} [-0.27, 0.28]$	$0.00^{+0.23}_{-0.23} [-0.53, 0.53]$

Combined fits (ZZ & WW) ATLAS

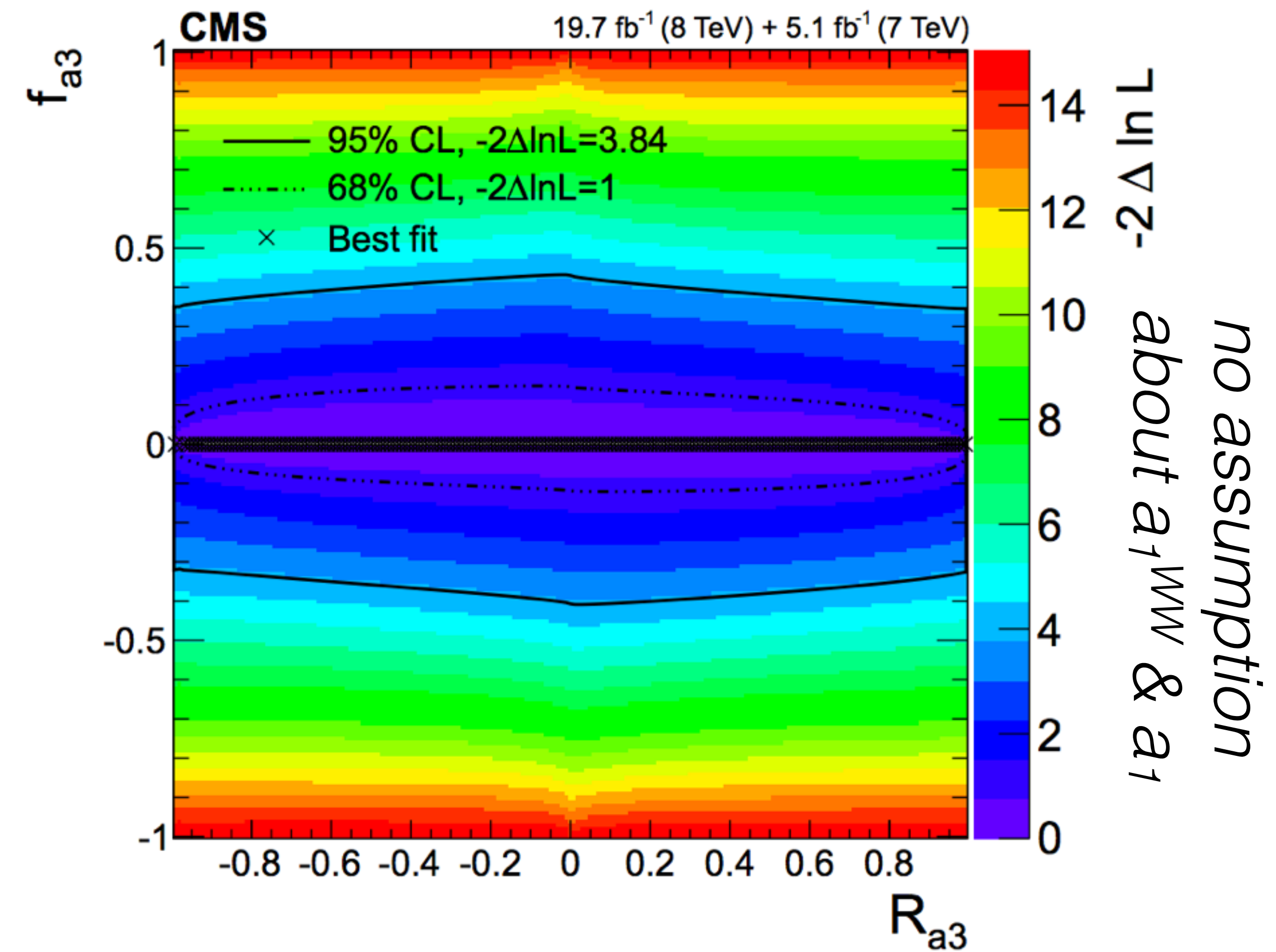
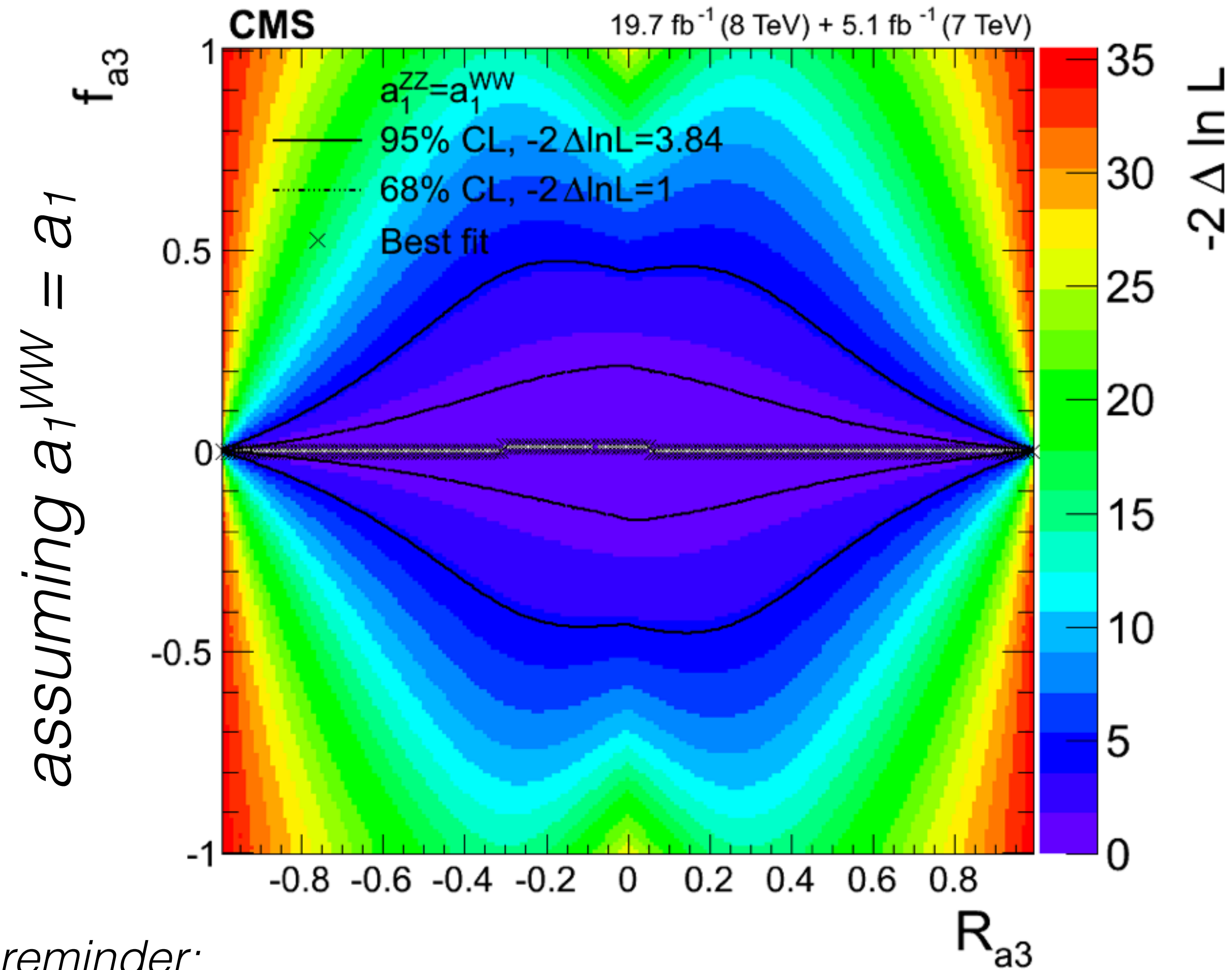
- Assuming $a_1/a_3 = a_1^{WW}/a_3^{WW}$ ($R_{a3} = 0.5$)
- Deviation slightly reduced when combining ZZ & WW



Coupling ratio	Best fit value		95% CL Exclusion Regions	
	Expected	Observed	Expected	Observed
$\tilde{\kappa}_{HVV}/\kappa_{SM}$	0.0	-0.48	$(-\infty, -0.55] \cup [4.80, \infty)$	$(-\infty, -0.73] \cup [0.63, \infty)$
$(\tilde{\kappa}_{AVV}/\kappa_{SM}) \cdot \tan \alpha$	0.0	-0.68	$(-\infty, -2.33] \cup [2.30, \infty)$	$(-\infty, -2.18] \cup [0.83, \infty)$

Combined fits (ZZ&WW) CMS

- This combination takes advantage of the relationship between the expected yields in the two channels, significantly enhancing the overall sensitivity.



reminder:

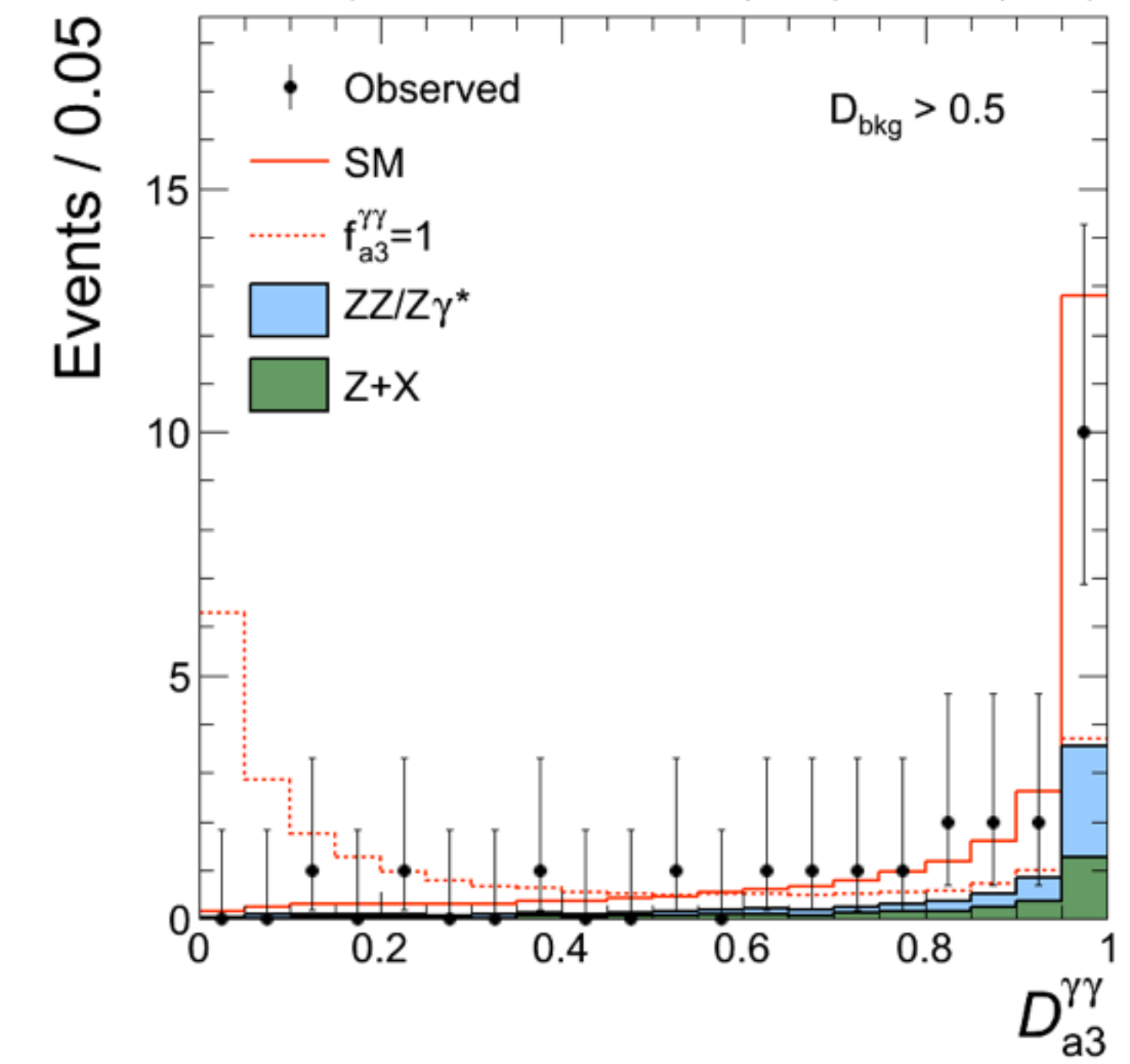
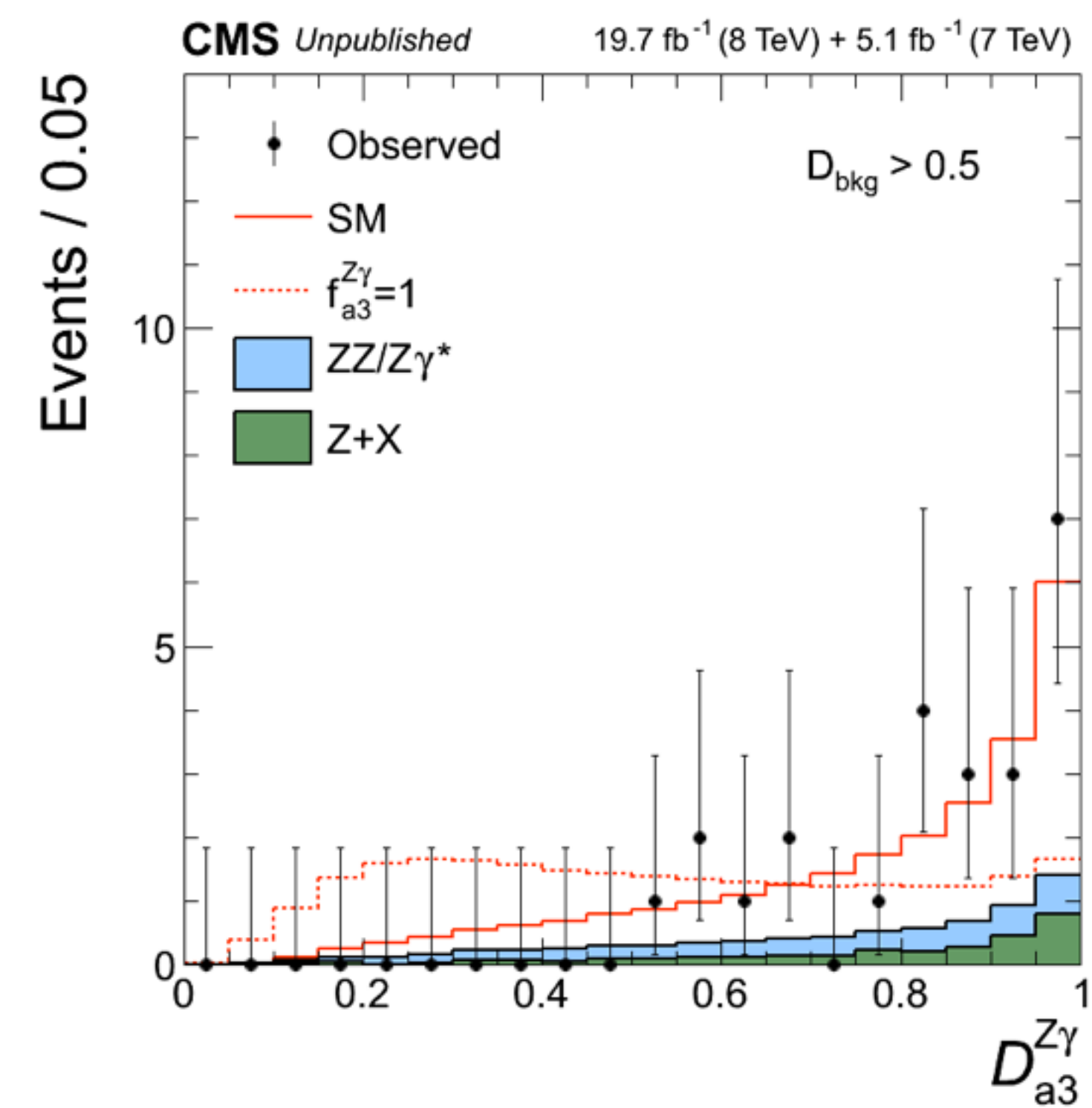
$$r_{ai} = \frac{a_i^{WW} / a_1^{WW}}{a_i / a_1}, \text{ or } R_{ai} = \frac{r_{ai} |r_{ai}|}{1 + r_{ai}^2}.$$

$\gamma\gamma$ & $Z\gamma$ couplings

- Can even start to probe $Z\gamma$ and $\gamma\gamma$ couplings with 4ℓ events
- Different discriminants used in order to be sensitive to these couplings

Measurement	Observables \vec{x}		
$f_{a3}^{\gamma\gamma}$	\mathcal{D}_{bkg}	$\mathcal{D}_{a3}^{\gamma\gamma}$	$\mathcal{D}_{\text{CP}}^{\gamma\gamma}$
$f_{a3}^{Z\gamma}$	\mathcal{D}_{bkg}	$\mathcal{D}_{a3}^{Z\gamma}$	$\mathcal{D}_{\text{CP}}^{Z\gamma}$

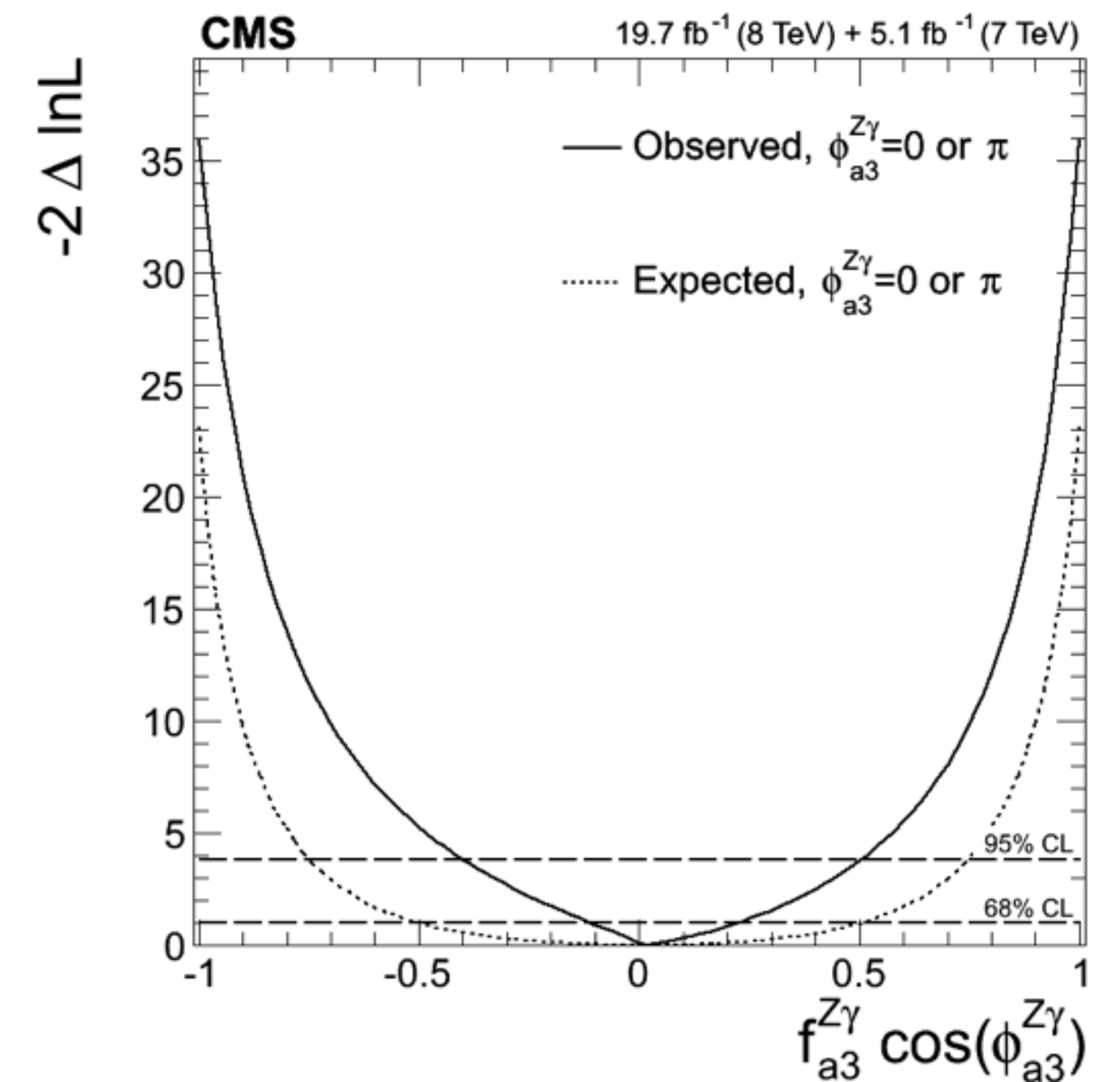
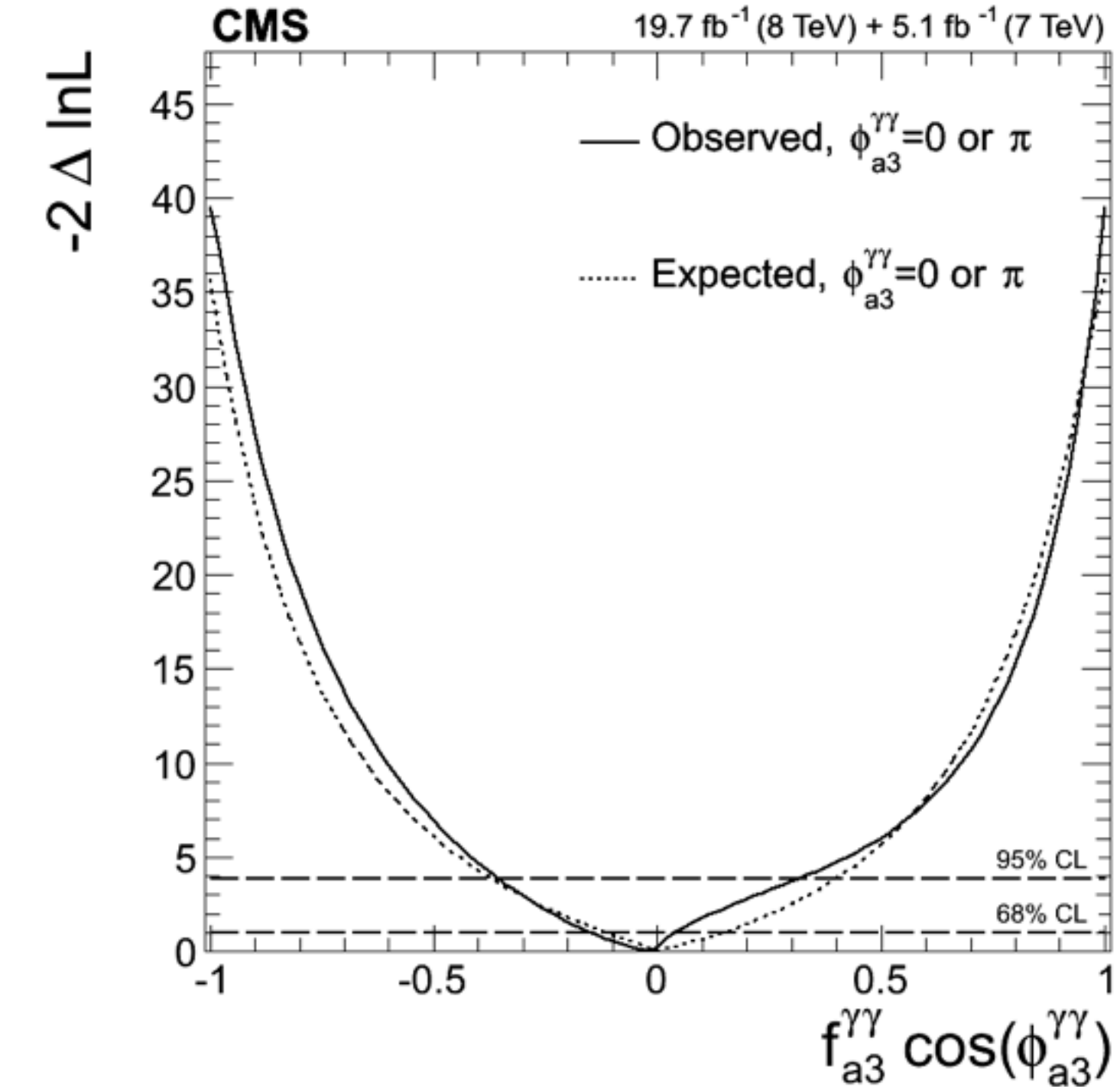
- note, sensitivity to these events is much weaker with 4ℓ events than search with on-shell photons
- direct searches with on-shell photons are not sensitive to different CP states (only to the total cross section)



$\gamma\gamma$ & $Z\gamma$ couplings

- Overall sensitivity still far from SM expectation
 - orders of magnitude small than searches with on-shell photons
 - 95% C.L. upper limit on $\mu^{Z\gamma}$ ($\mu^{\gamma\gamma}$) \sim 170 (730) with 4ℓ events
 - 95% (68%) C.L upper limit on $\mu^{Z\gamma} = 9.5$ with $2\ell + \gamma$ events

Parameter	Observed	Expected	$f_{ai}^{VV} = 1$
$f_{\Delta 1}^{Z\gamma} \cos(\phi_{\Delta 1}^{Z\gamma})$	$-0.27^{+0.34}_{-0.49} [-1.00, 1.00]$	$0.00^{+0.83}_{-0.53} [-1.00, 1.00]$	26% (16%)
$f_{a2}^{Z\gamma} \cos(\phi_{a2}^{Z\gamma})$	$0.00^{+0.14}_{-0.20} [-0.49, 0.46]$	$0.00^{+0.51}_{-0.51} [-0.78, 0.79]$	<0.01% (0.01%)
$f_{a3}^{Z\gamma} \cos(\phi_{a3}^{Z\gamma})$	$0.02^{+0.21}_{-0.13} [-0.40, 0.51]$	$0.00^{+0.51}_{-0.51} [-0.75, 0.75]$	<0.01% (<0.01%)
$f_{a2}^{\gamma\gamma} \cos(\phi_{a2}^{\gamma\gamma})$	$0.12^{+0.20}_{-0.11} [-0.04, +0.51]$	$0.00^{+0.11}_{-0.09} [-0.32, 0.34]$	<0.01% (<0.01%)
$f_{a3}^{\gamma\gamma} \cos(\phi_{a3}^{\gamma\gamma})$	$-0.02^{+0.06}_{-0.13} [-0.35, 0.32]$	$0.00^{+0.15}_{-0.11} [-0.37, 0.40]$	<0.01% (<0.01%)



Prospects for the LHC CMS

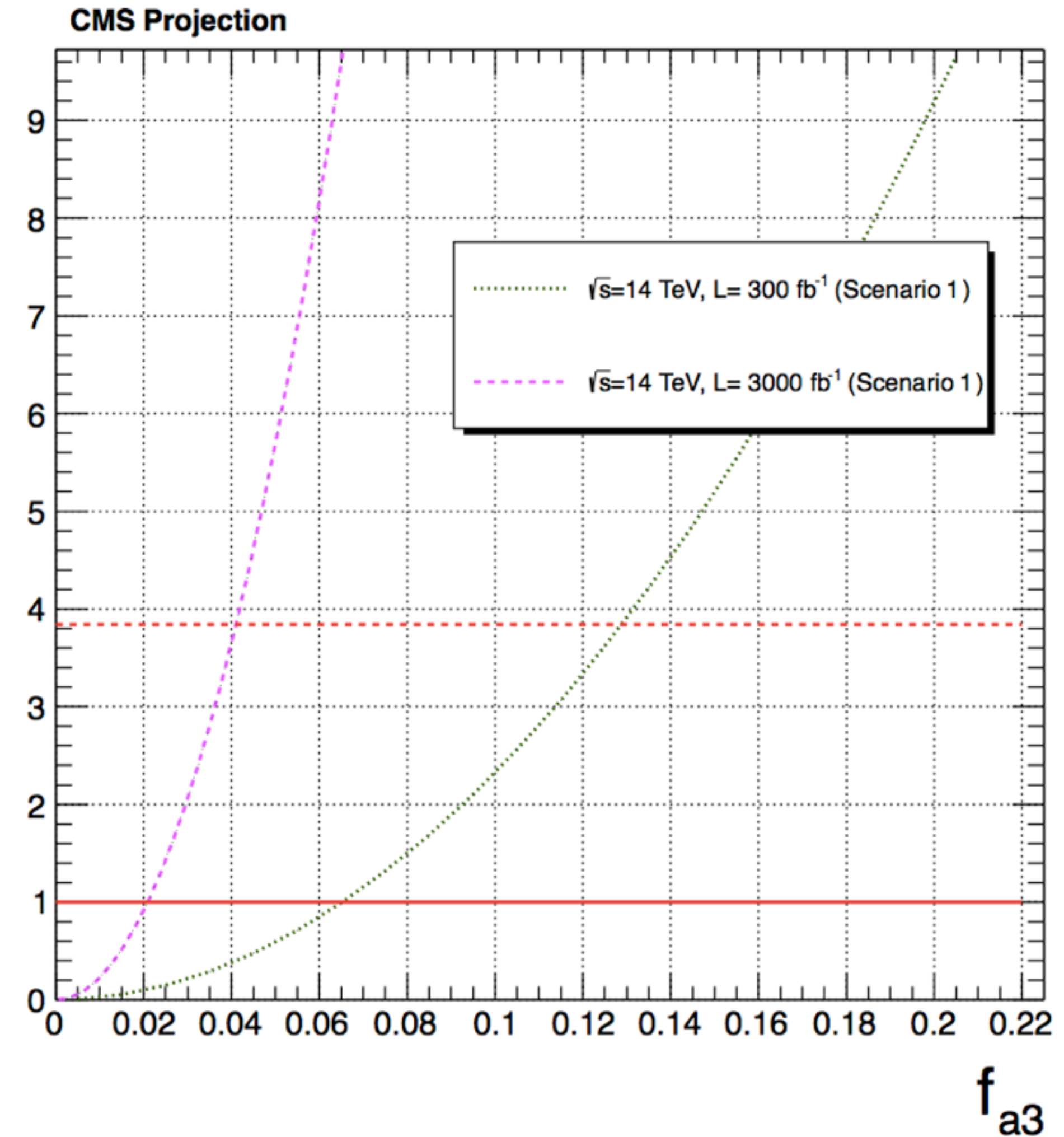
- Projections out to 300/fb & 3000/fb
- fits are performed with templates in which events are described by 2D distributions of kinematic discriminants:

$$\mathcal{D}_{\text{bkg}} = \left[1 + \frac{\mathcal{P}_{\text{bkg}}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell}) \times \mathcal{P}_{\text{bkg}}^{\text{mass}}(m_{4\ell})}{\mathcal{P}_{0^+}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell}) \times \mathcal{P}_{\text{sig}}^{\text{mass}}(m_{4\ell} | m_{0^+})} \right]^{-1}$$

$$\mathcal{D}_{J^P} = \left[1 + \frac{\mathcal{P}_{J^P}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})}{\mathcal{P}_{0^+}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})} \right]^{-1}$$

- signals model with LO MC
- scaling up all background predictions
- all systematic uncertainties are assumed to be the same:
still dominated by statistical uncertainties
- can reach 95% C.L. upper limits of $f_{a3} < 0.13$ (0.04)

2 x Δ(NLL)



Prospects for the LHC ATLAS

- Exploring sensitivity to couplings @ 300 & 3000/fb
- background: assuming qq->ZZ background scaled up to account for irreducible background
- signal: LO MC - reweighing to morph signal hypothesis
- Fitting either
 - templates based on ME-based kinematic discriminants
 - fully correlated 8D likelihood
(including approx. detector effects)
- Systematics uncertainties approximated to be: 3% luminosity, 5% signal and background yields (lepton reco ϵ), and 9.4% (7.4%) for 300/fb (3000/fb) for background yield

Observable	Sensitivity
$\ln \frac{ \text{ME}(g_1=1, g_2=0, g_4=0) ^2}{ \text{ME}(g_1=0, g_2=0, g_4=1) ^2}$	$ g_4 /g_1$
$\ln \frac{ \text{ME}(g_1=1, g_2=0, g_4=-2+2i) ^2}{ \text{ME}(g_1=1, g_2=0, g_4=2+2i) ^2}$	$\Re(g_4)/g_1$
$\ln \frac{ \text{ME}(g_1=1, g_2=0, g_4=2-2i) ^2}{ \text{ME}(g_1=1, g_2=0, g_4=2+2i) ^2}$	$\Im(g_4)/g_1$
$\ln \frac{ \text{ME}(g_1=1, g_2=0, g_4=0) ^2}{ \text{ME}(g_1=1, g_2=1, g_4=0) ^2}$	$ g_2 /g_1$
$\ln \frac{ \text{ME}(g_1=1, g_2=-1+i, g_4=0) ^2}{ \text{ME}(g_1=1, g_2=1+i, g_4=0) ^2}$	$\Re(g_2)/g_1$
$\ln \frac{ \text{ME}(g_1=1, g_2=1-i, g_4=0) ^2}{ \text{ME}(g_1=1, g_2=1+i, g_4=0) ^2}$	$\Im(g_2)/g_1$

$g_1(g_4)$ is effectively the same as $(a_1) a_3$

Final State	Signal	Background
4e	871	474
4 μ	1186	641
2e2 μ	1035	574
2 μ 2e	867	431

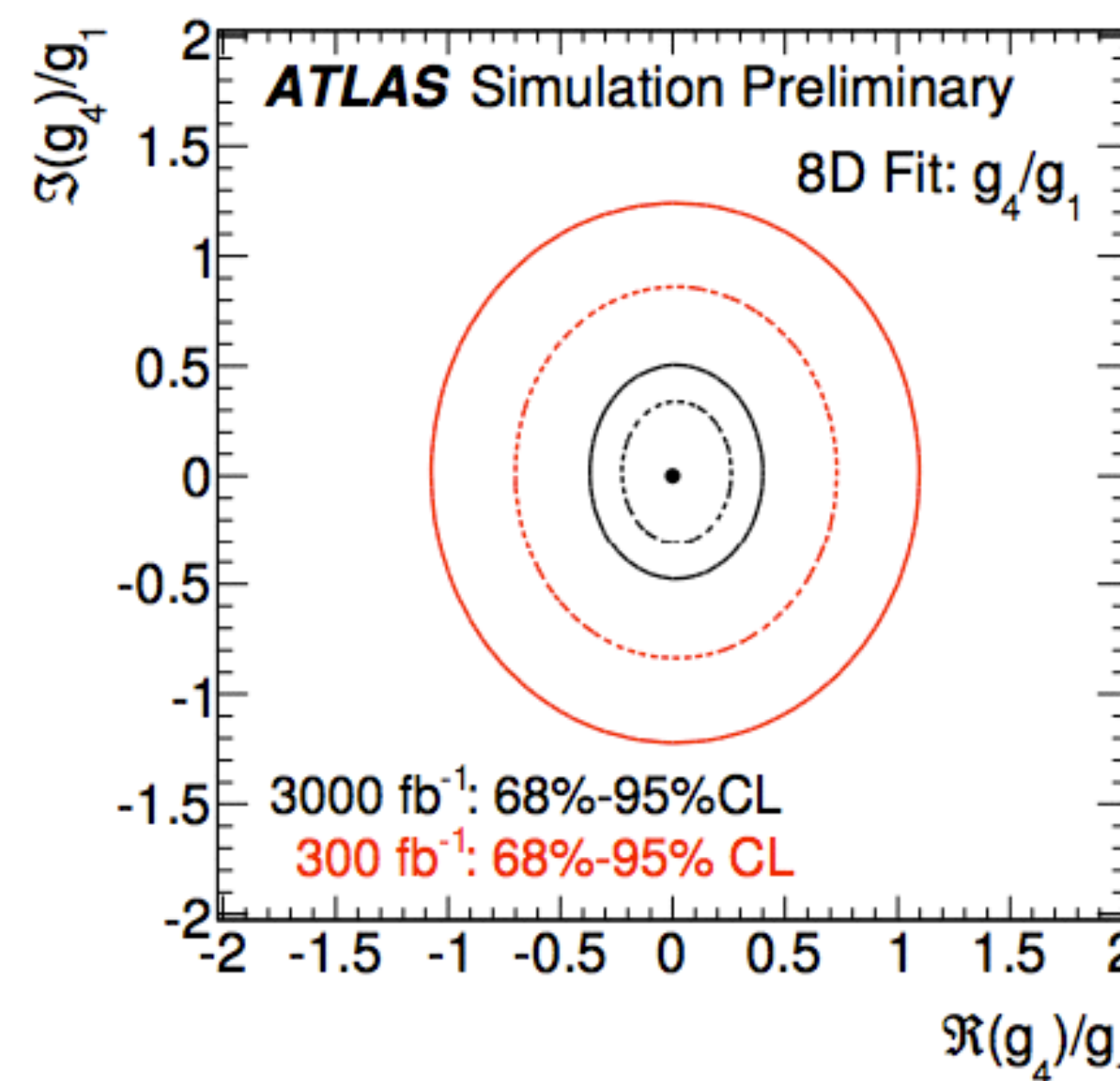
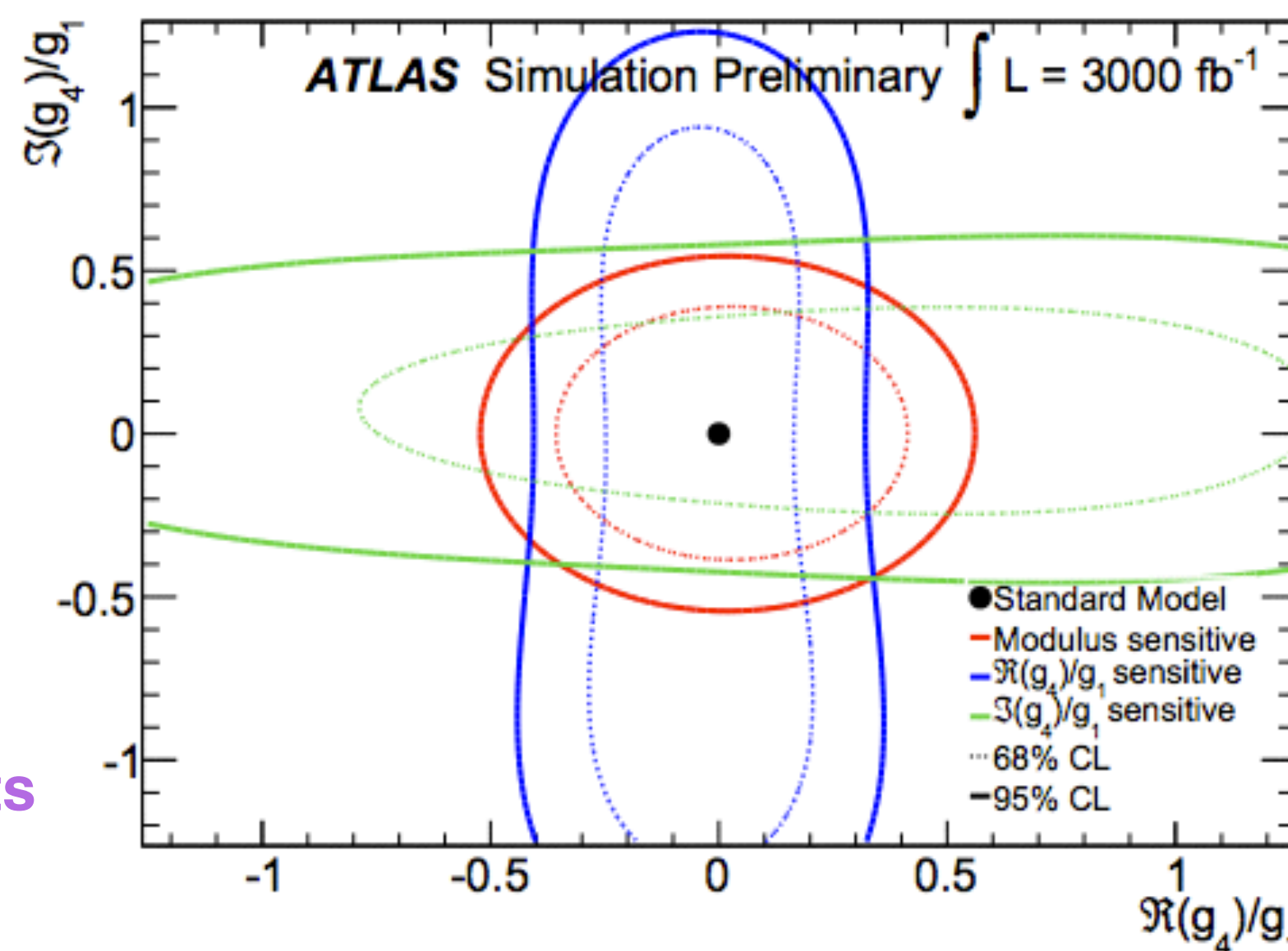


Constraints on CP violating interactions



- Likelihood scan in real and imaginary projection of g_4/g_1
 - KD fits show that phase of anomalous couplings can play an important role - *note this can also be parameterized by extending templates' dimensionality a la CMS*
 - 8D - correlation less between the real and imaginary projections - *simultaneously parameterizes kinematics in term of magnitude & phase of anomalous couplings*

ME sensitive fits



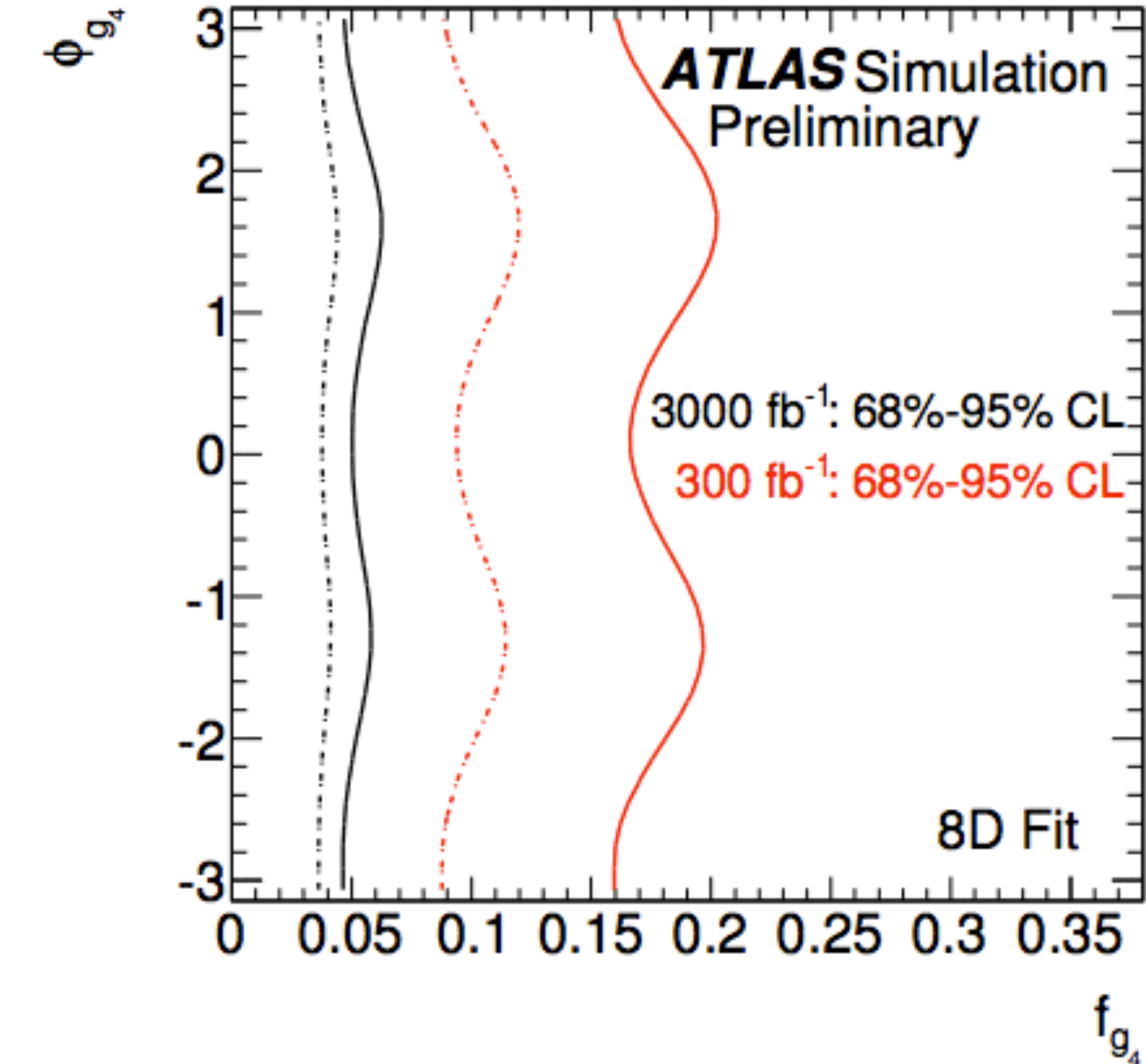
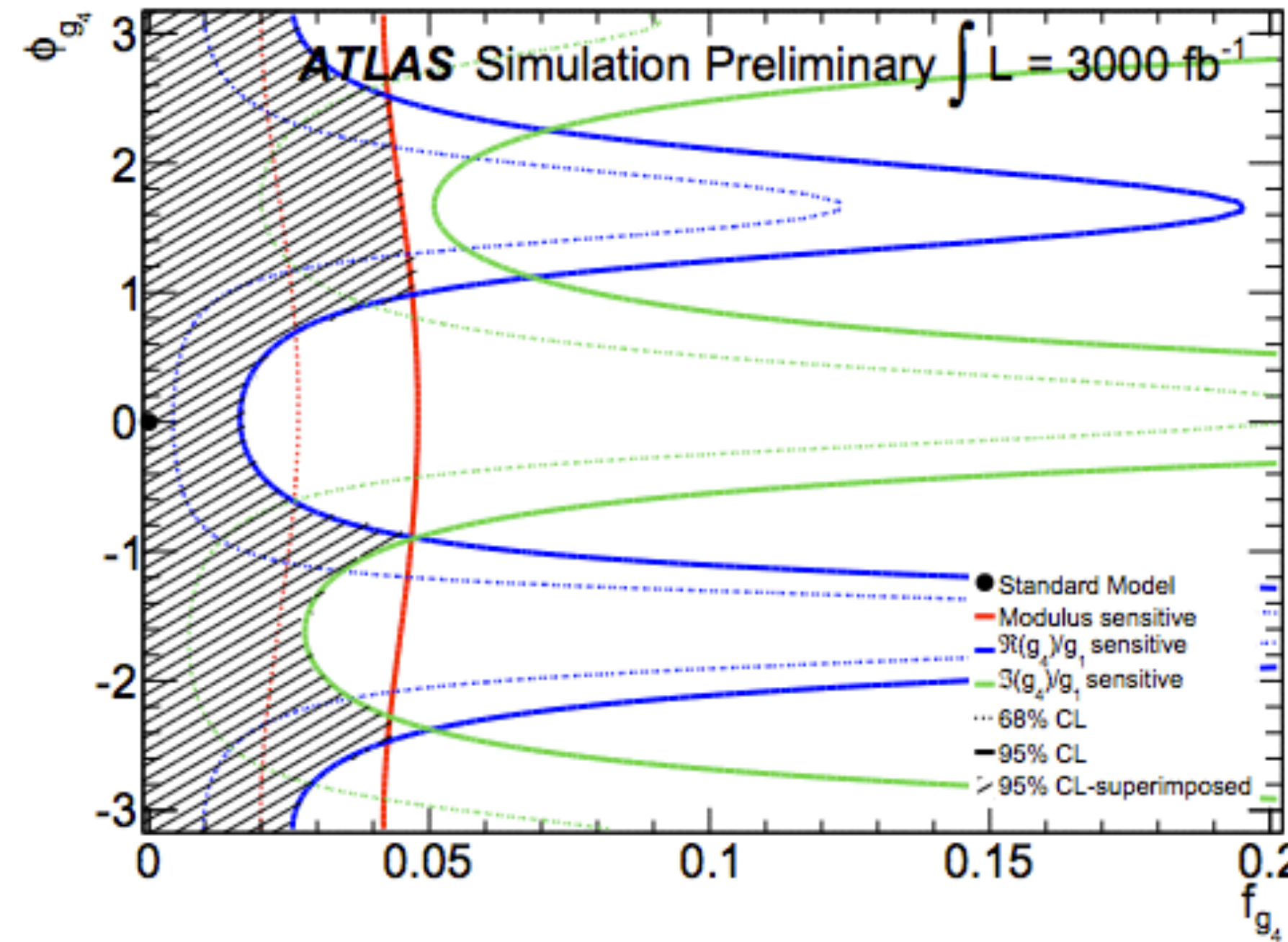
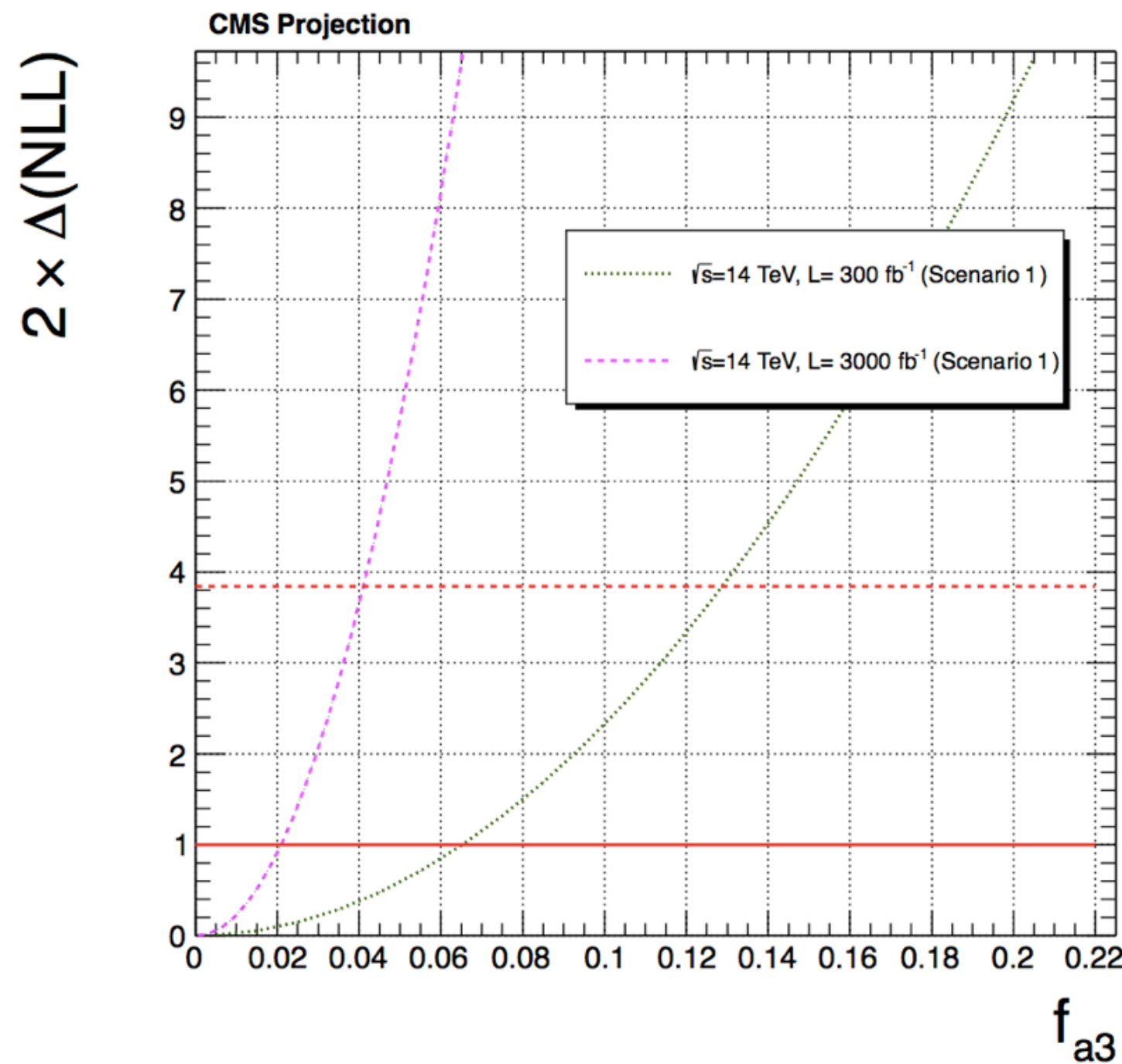
8D fits

Luminosity	$ g_4 /g_1$	$\Re(g_4)/g_1$	$\Im(g_4)/g_1$	$ g_2 /g_1$	$\Re(g_2)/g_1$	$\Im(g_2)/g_1$
300 fb^{-1}	1.03	(-1.01, 1.01)	(-1.02, 1.02)	1.39	(-0.88, 0.38)	(-1.13, 1.13)
3000 fb^{-1}	0.49	(-0.34, 0.26)	(-0.34, 0.48)	0.81	(-0.33, 0.11)	(-0.73, 0.75)

Luminosity	$ g_4 /g_1$	$\Re(g_4)/g_1$	$\Im(g_4)/g_1$	$ g_2 /g_1$	$\Re(g_2)/g_1$	$\Im(g_2)/g_1$
300	1.20	(-0.88, 0.91)	(-1.02, 1.05)	1.02	(-0.84, 0.44)	(-1.19, 1.18)
3000	0.60	(-0.30, 0.33)	(-0.39, 0.42)	0.60	(-0.30, 0.11)	(-0.71, 0.68)

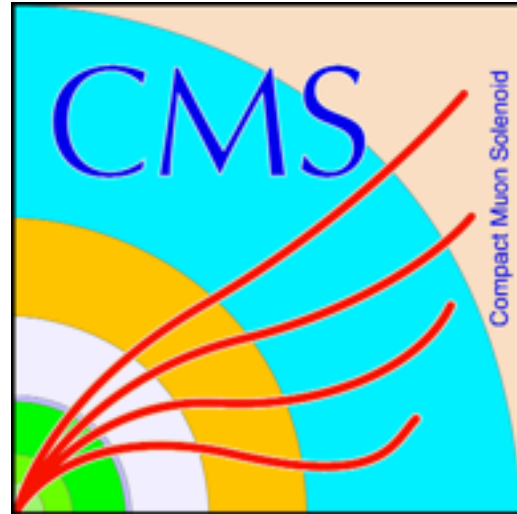
Comparison

- ATLAS analysis also produced likelihood scan as a function of f_{a_i} & ϕ_i
- $f_{a3} = f_{g4}$ — comparable results between ATLAS and CMS
($f_{a3} < 0.04$ @ 95% C.L.)



NOTE: the parameterization used in CMS is equivalent to the ATLAS *modulus sensitive* variable

Summary of results



Parameter	Observed	Expected
$H \rightarrow WW \rightarrow e\nu\mu\nu$		
$f_{a3}^{WW} \cos(\phi_{a3}^{WW})$	$-0.03^{+1.03}_{-0.97} [-1.00, 1.00]$	$0.00^{+1.00}_{-1.00} [-1.00, 1.00]$
$H \rightarrow ZZ \rightarrow 4\ell$		
$f_{a3} \cos(\phi_{a3})$	$0.00^{+0.14}_{-0.11} [-0.40, 0.43]$	$0.00^{+0.33}_{-0.33} [-0.70, 0.70]$
<i>Combined $H \rightarrow ZZ \rightarrow 4\ell$ & $H \rightarrow WW \rightarrow e\nu\mu\nu$</i>		
$f_{a3} \cos(\phi_{a3})$	$0.00^{+0.08}_{-0.08} [-0.27, 0.28]$	$0.00^{+0.23}_{-0.23} [-0.53, 0.53]$

Observed 95% C.L Limits



$H \rightarrow WW^* \rightarrow e\nu\mu\nu$	
$f_{g4} < 0.78$ for $\phi_{g4} = 0$	and $f_{g4} < 0.84$ for $\phi_{g4} = \pi$
$H \rightarrow ZZ^* \rightarrow 4\ell$	
$f_{g4} < 0.11$ for $\phi_{g4} = 0$	and $f_{g4} < 0.54$ for $\phi_{g4} = \pi$
<i>Combination of $H \rightarrow ZZ^* \rightarrow 4\ell$ and $H \rightarrow WW^* \rightarrow e\nu\mu\nu$</i>	
$f_{g4} < 0.090$ for $\phi_{g4} = 0$	and $f_{g4} < 0.41$ for $\phi_{g4} = \pi$

Conclusions

- CMS & ATLAS have mature programs for constraining the tensor structure of HVV interactions (HZZ, HWW, H $\gamma\gamma$, and HZ γ)
 - advanced techniques used to exploit all of the decay kinematics of the 4-body final states
 - measurements still limited by statistical uncertainties
 - constraining $a_2^{\gamma\gamma}$ ($a_2^{\gamma\gamma}$) from $a_3^{Z\gamma}$ ($a_3^{Z\gamma}$) couplings with high integrated luminosity is possible
- Projections show that ZZ anomalous coupling constraints can be improved by a factor of 10 with the full HL-LHC integrated luminosity

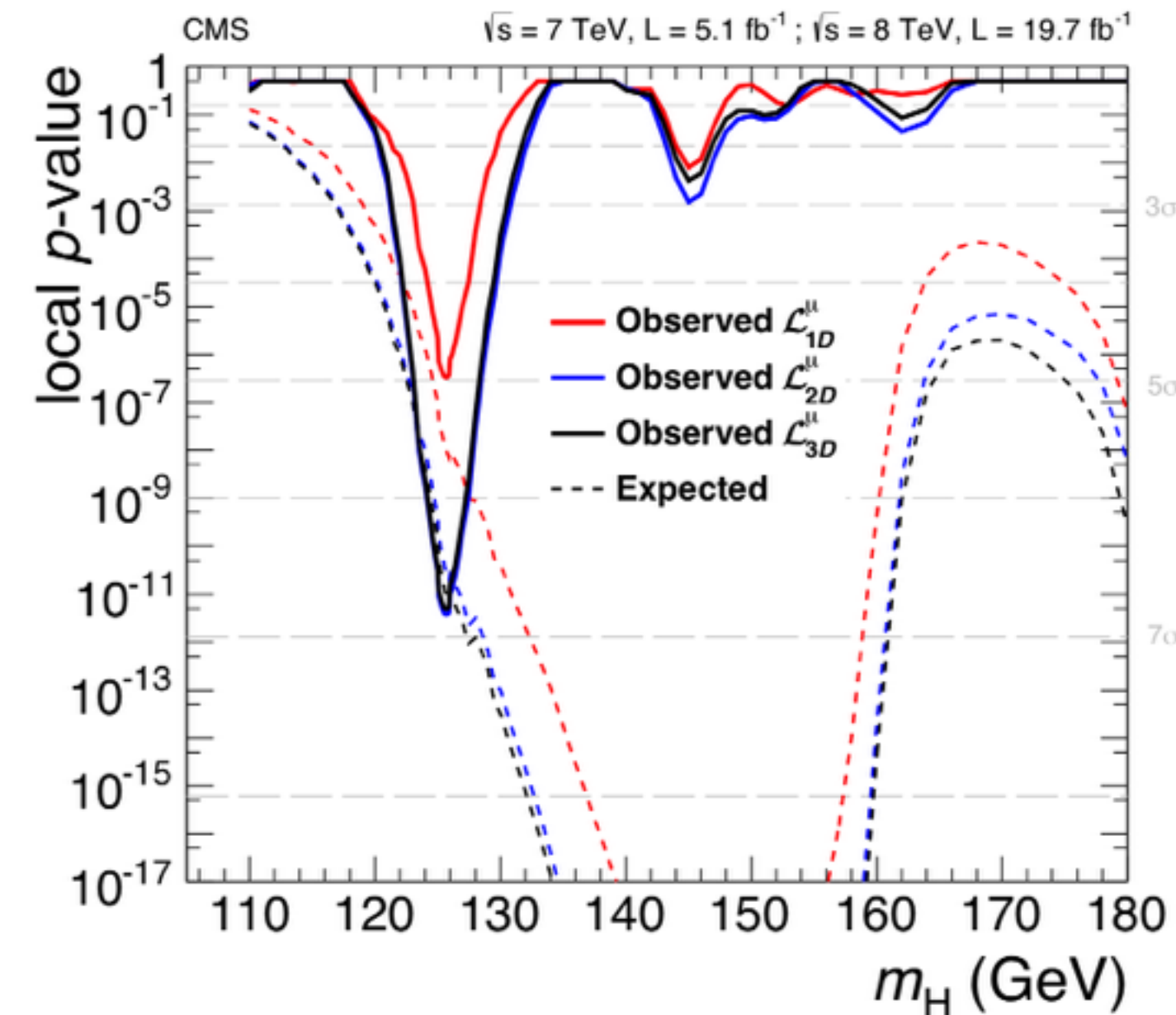
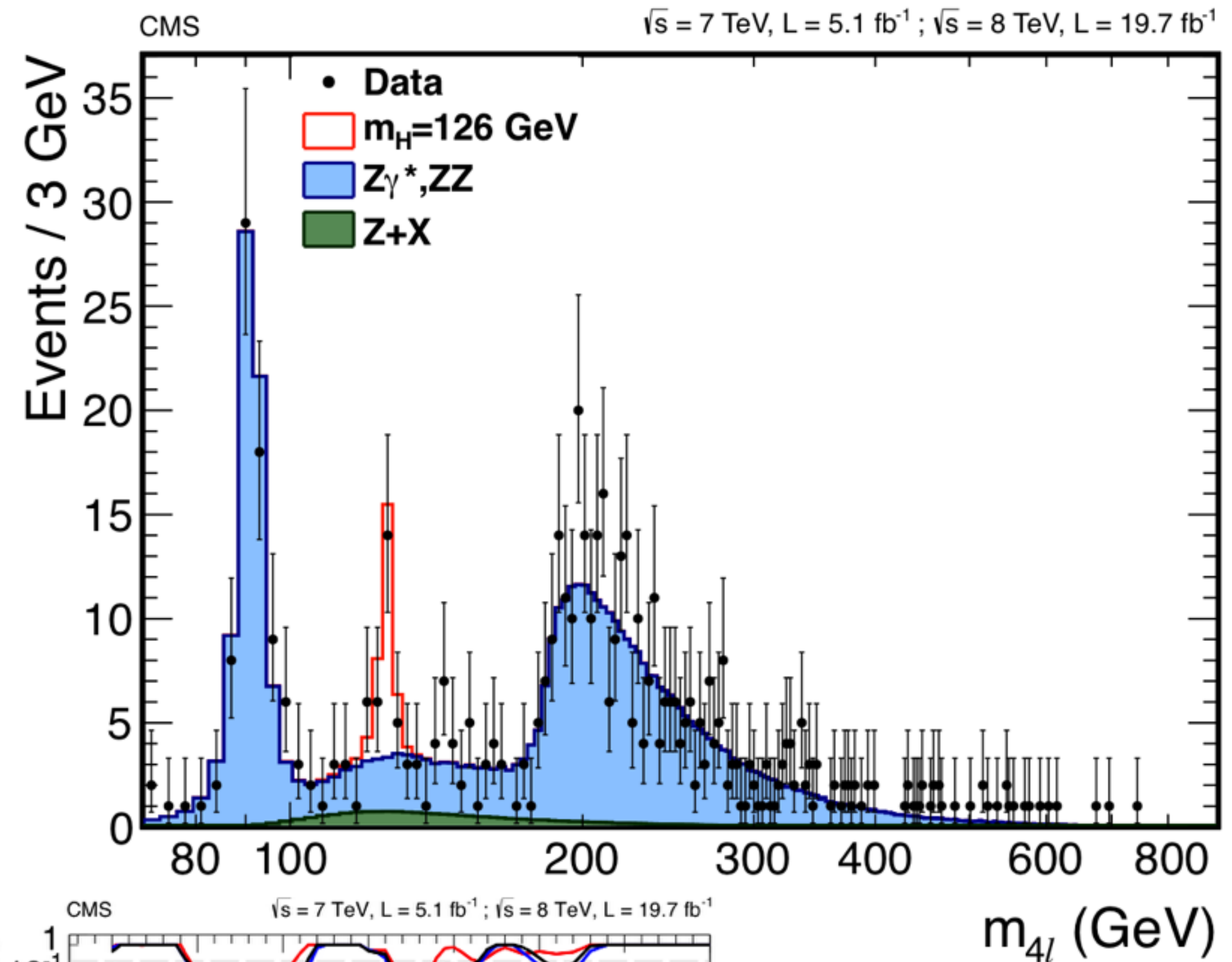
BACKUP

4l analyses CMS

- 2 pairs of opposite-sign same flavor leptons (e or μ)
(sub)leading lepton $p_T > 20$ (10 GeV)
 $m_{ll} \geq 4$ GeV (all combinations)
 $40 < m_1$ (closest to Z mass) < 120 GeV
 $12 < m_2 < 120$ GeV
- SM ZZ contributions taken from NLO MC sim.
- Z+X backgrounds estimated from a loose ID control region

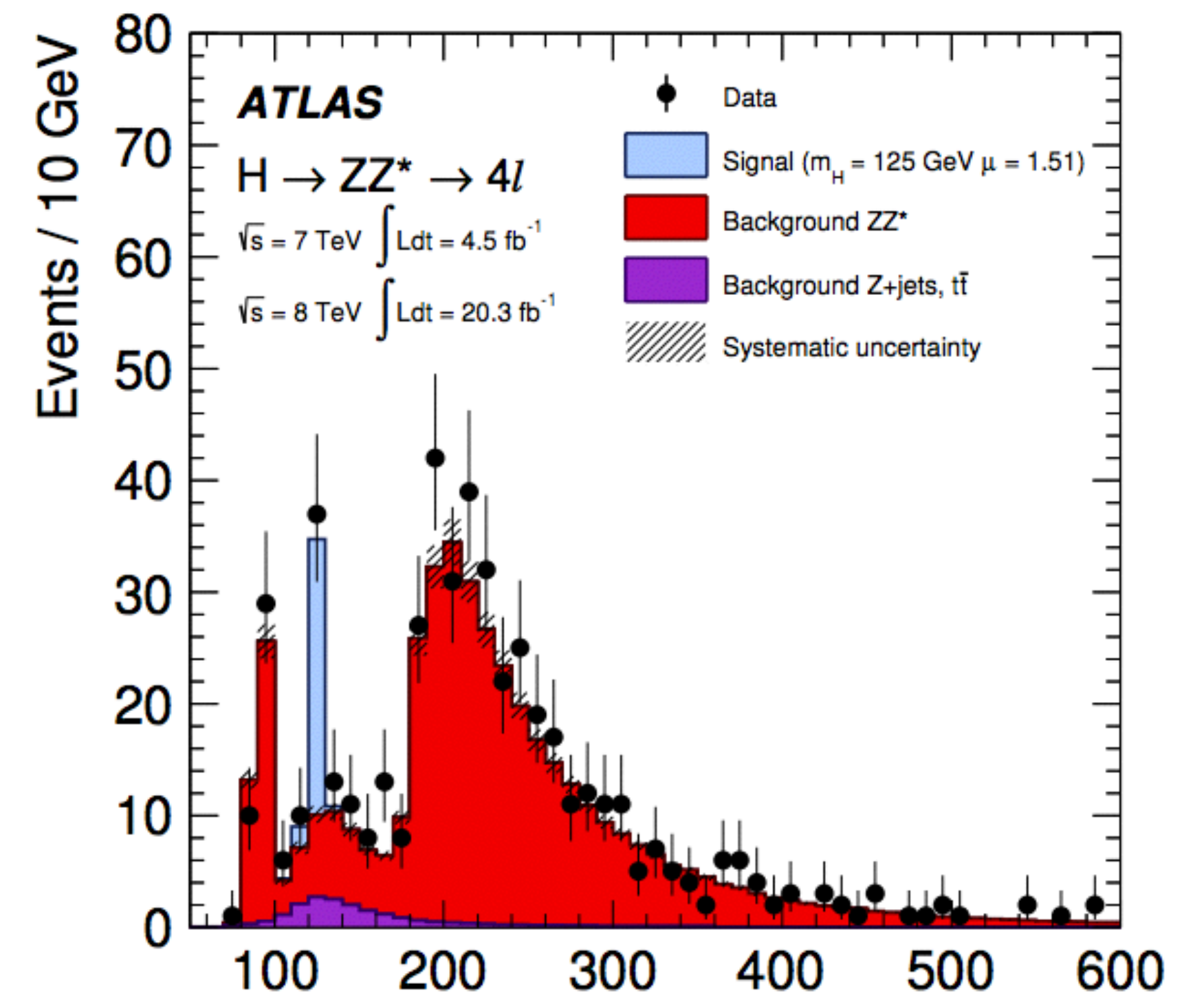
$$105.6 < m_{4\ell} < 140.6 \text{ GeV}$$

Channel	4e		4 μ		2e2 μ	
	7 TeV	8 TeV	7 TeV	8 TeV	7 TeV	8 TeV
$q\bar{q} \rightarrow ZZ$	0.84 ± 0.10	2.94 ± 0.33	1.80 ± 0.11	7.65 ± 0.49	2.24 ± 0.28	8.86 ± 0.68
$gg \rightarrow ZZ$	0.03 ± 0.01	0.20 ± 0.05	0.06 ± 0.02	0.41 ± 0.10	0.07 ± 0.02	0.50 ± 0.13
Z + X	0.62 ± 0.14	2.77 ± 0.62	0.22 ± 0.09	1.19 ± 0.48	1.06 ± 0.29	4.29 ± 1.10
Bkg.	1.49 ± 0.17	5.91 ± 0.71	2.08 ± 0.14	9.25 ± 0.69	3.37 ± 0.40	13.65 ± 1.30
Signal	0.70 ± 0.11	3.09 ± 0.47	1.24 ± 0.14	5.95 ± 0.71	1.67 ± 0.26	7.68 ± 0.98
Observed	1	9	3	15	6	16



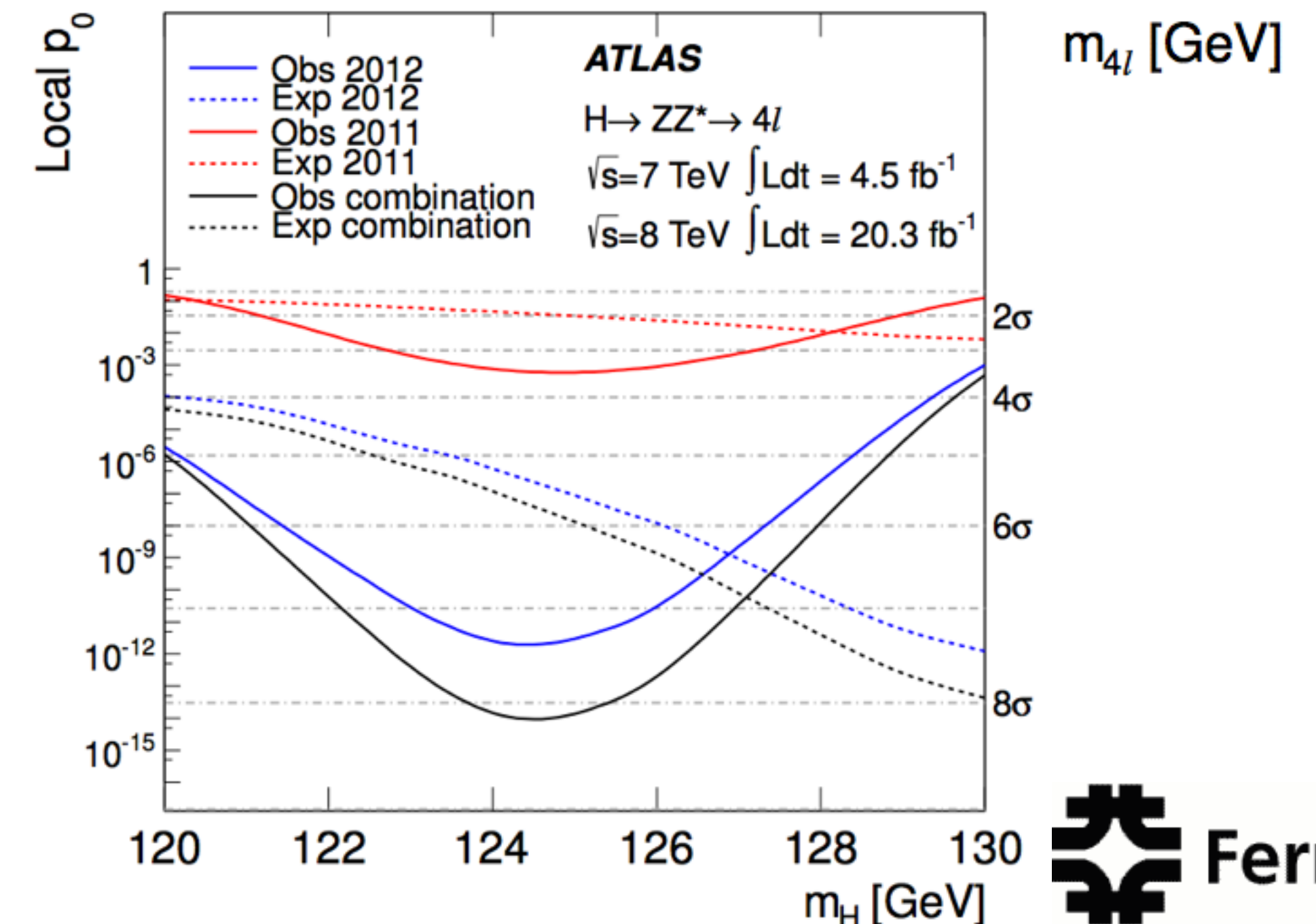
4l analyses ATLAS

- two pairs of opposite-sign same flavor leptons
three leading leptons $p_T > 20, 15, 10$ GeV
 $50 < m_{12}$ (closest to Z mass) < 106 GeV
 $12 < m_{34} < 115$ GeV
- SM ZZ contributions taken from NLO MC sim.
- Z+X backgrounds estimated from a loose ID control region



115 < m_{4l} < 130 GeV

	Signal	ZZ*	tt, Z + jets	Observed
$\sqrt{s} = 7$ TeV				
4 μ	1.02 \pm 0.10	0.65 \pm 0.03	0.14 \pm 0.06	3
2 μ 2e	0.47 \pm 0.05	0.29 \pm 0.02	0.53 \pm 0.12	1
2e2 μ	0.64 \pm 0.06	0.45 \pm 0.02	0.13 \pm 0.05	2
4e	0.45 \pm 0.04	0.26 \pm 0.02	0.59 \pm 0.12	2
Total	2.58 \pm 0.25	1.65 \pm 0.09	1.39 \pm 0.26	8
$\sqrt{s} = 8$ TeV				
4 μ	5.81 \pm 0.58	3.36 \pm 0.17	0.97 \pm 0.18	13
2 μ 2e	3.00 \pm 0.30	1.59 \pm 0.10	0.52 \pm 0.12	8
2e2 μ	3.72 \pm 0.37	2.33 \pm 0.11	0.84 \pm 0.14	9
4e	2.91 \pm 0.29	1.44 \pm 0.09	0.52 \pm 0.11	7
Total	15.4 \pm 1.5	8.72 \pm 0.47	2.85 \pm 0.39	37



Observables

- An example of a complete set of observables describing the higgs decay kinematics

$$m_{4l} = (q_{11} + q_{12} + q_{21} + q_{22})^2$$

$$m_{z1} = (q_{11} + q_{12})^2$$

$$m_{z2} = (q_{21} + q_{22})^2$$

$$\cos(\theta^*) = (q_{11} + q_{12}) \cdot z / |q_{11} + q_{12}|$$

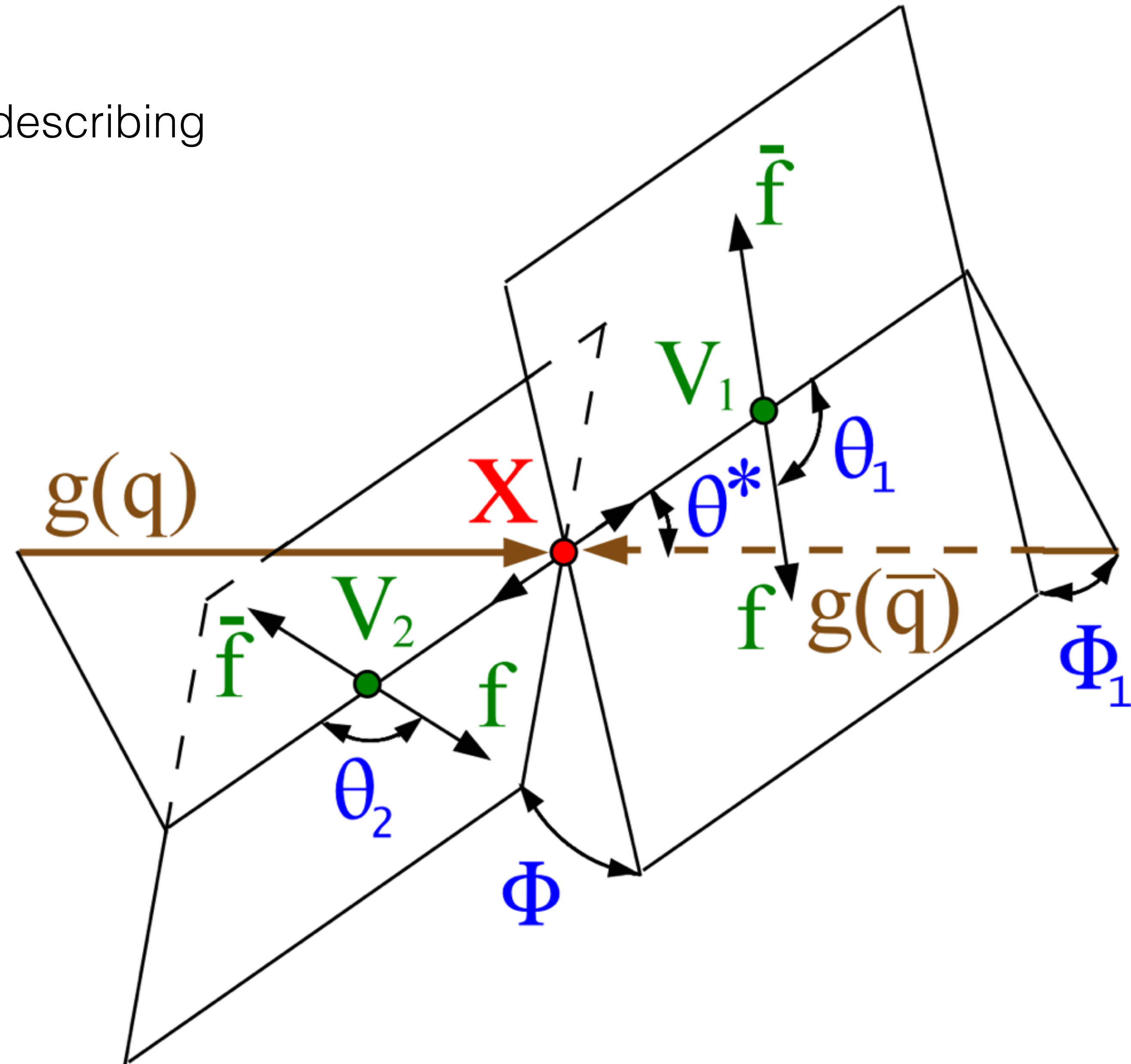
$$\Phi = \frac{\mathbf{q}_1 \cdot (\hat{\mathbf{n}}_1 \times \hat{\mathbf{n}}_2)}{|\mathbf{q}_1 \cdot (\hat{\mathbf{n}}_1 \times \hat{\mathbf{n}}_2)|} \times \cos^{-1}(-\hat{\mathbf{n}}_1 \cdot \hat{\mathbf{n}}_2)$$

$$\Phi_1 = \frac{\mathbf{q}_1 \cdot (\hat{\mathbf{n}}_1 \times \hat{\mathbf{n}}_{sc})}{|\mathbf{q}_1 \cdot (\hat{\mathbf{n}}_1 \times \hat{\mathbf{n}}_{sc})|} \times \cos^{-1}(\hat{\mathbf{n}}_1 \cdot \hat{\mathbf{n}}_{sc})$$

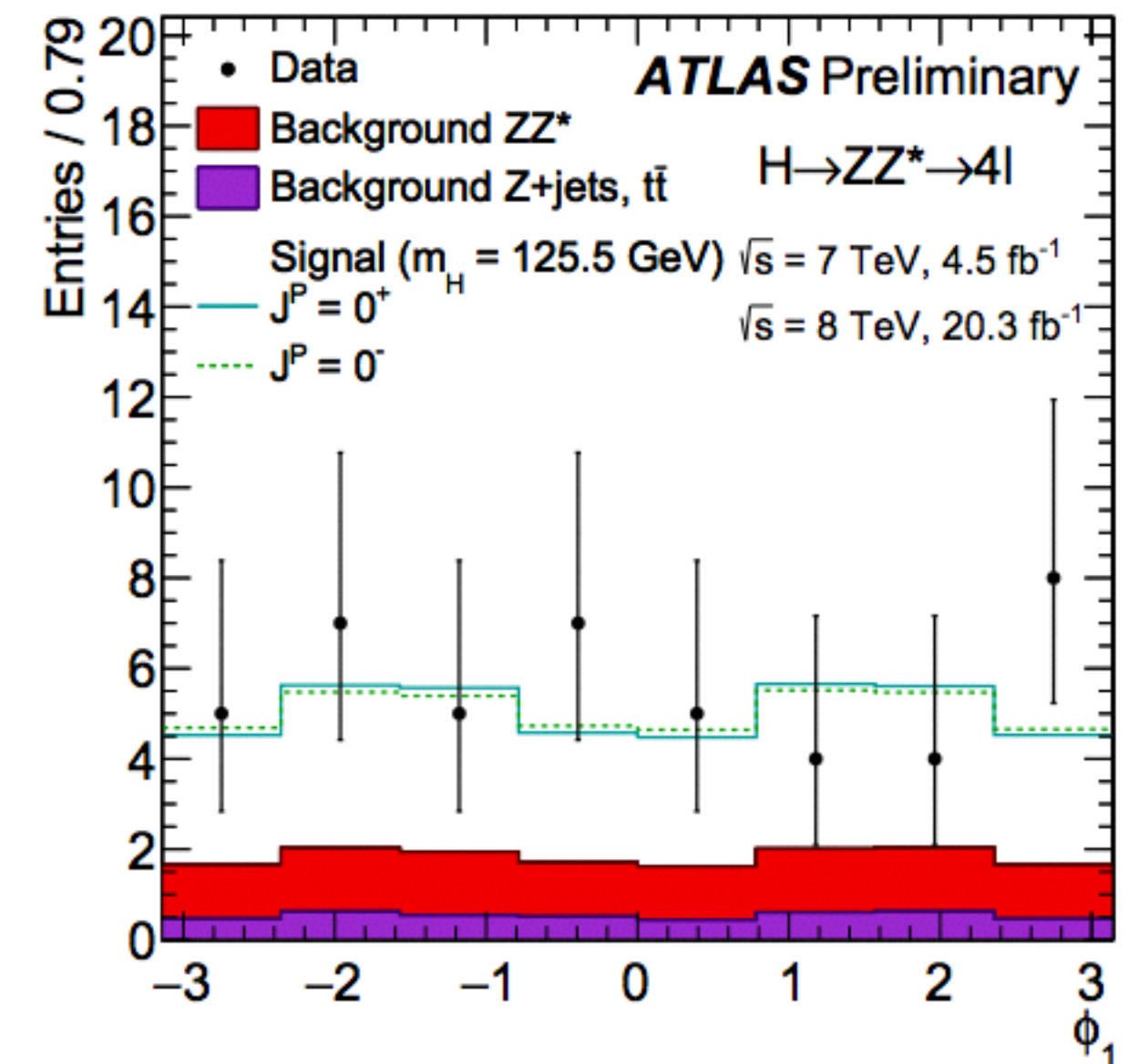
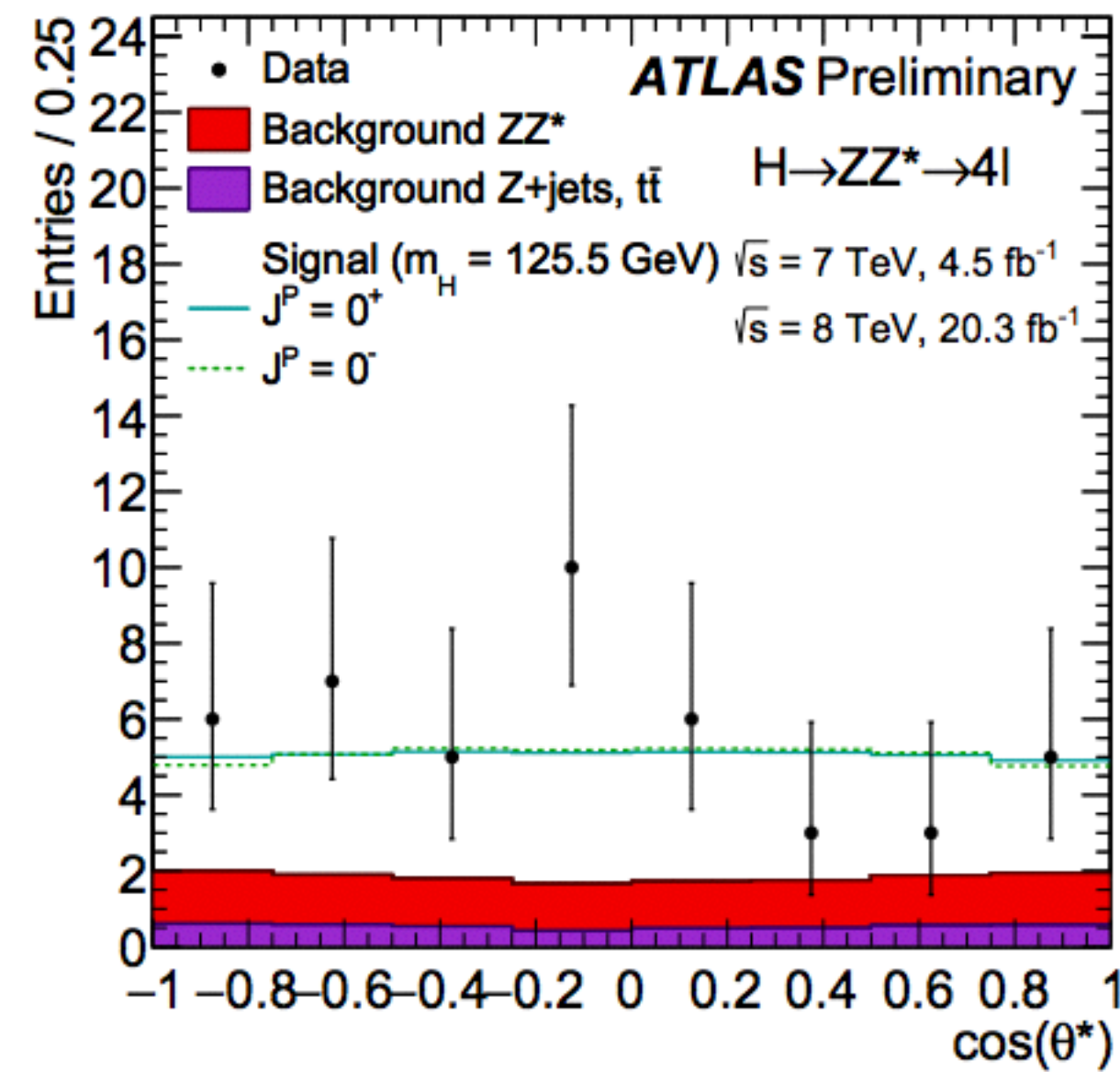
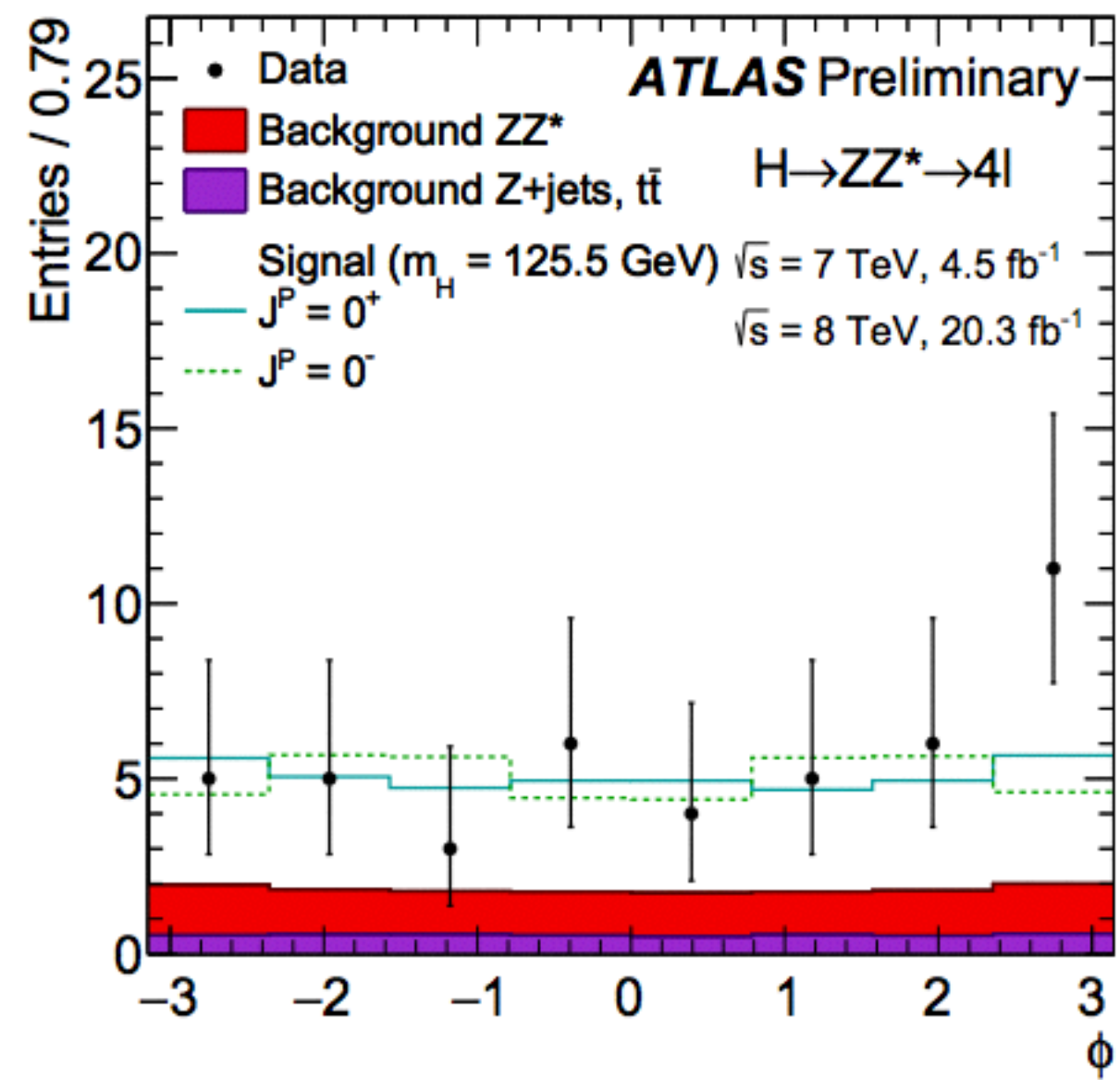
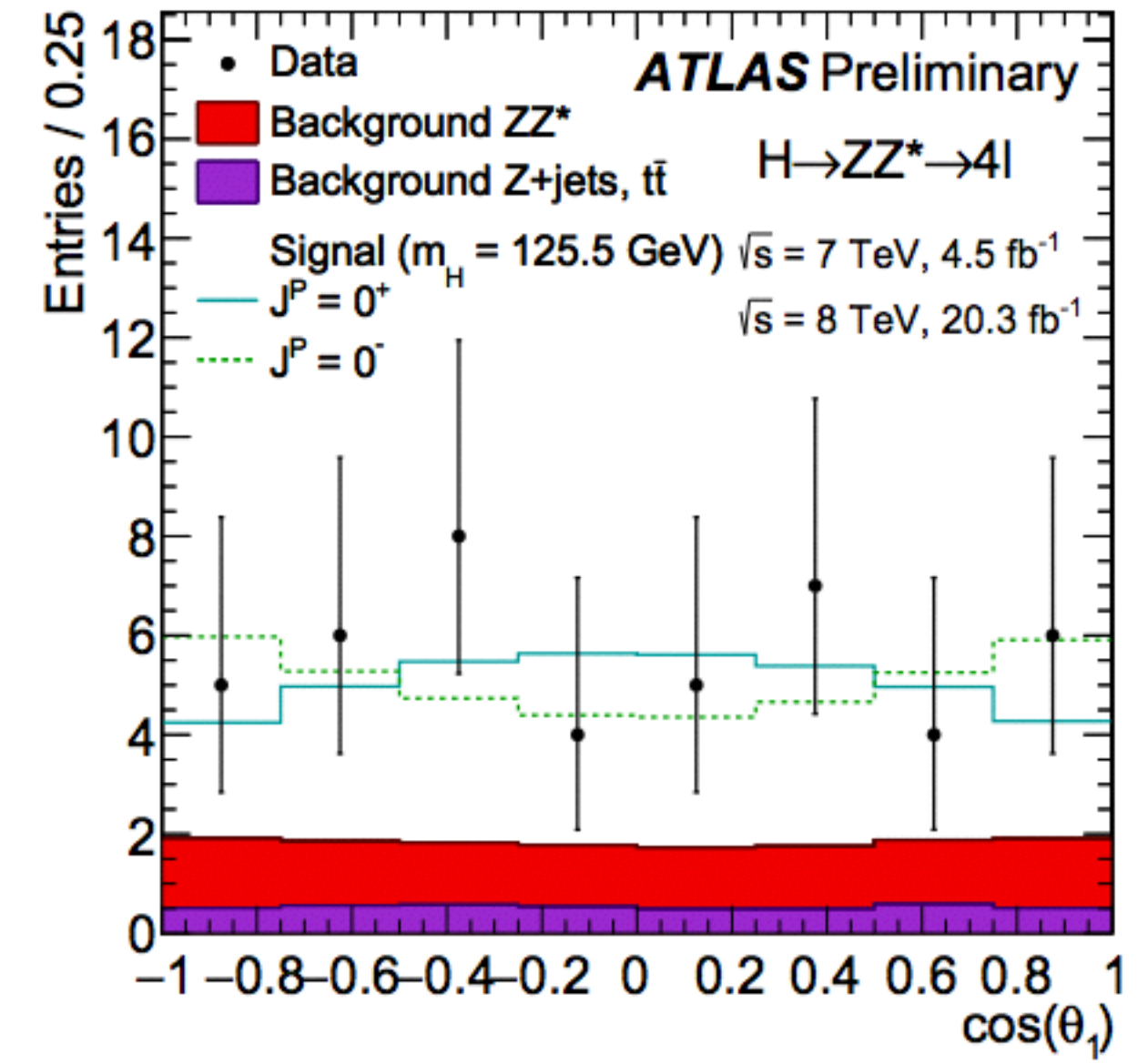
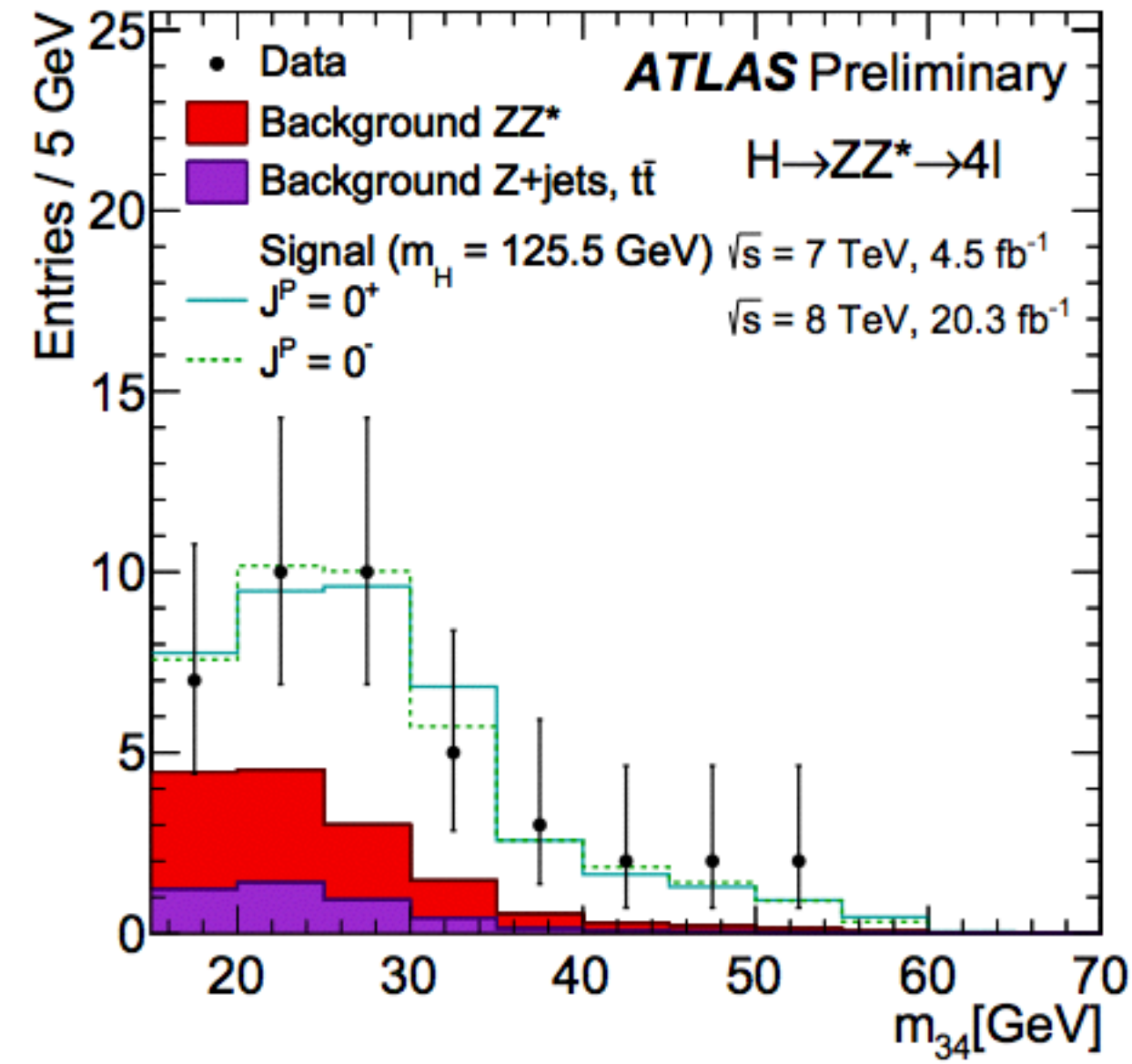
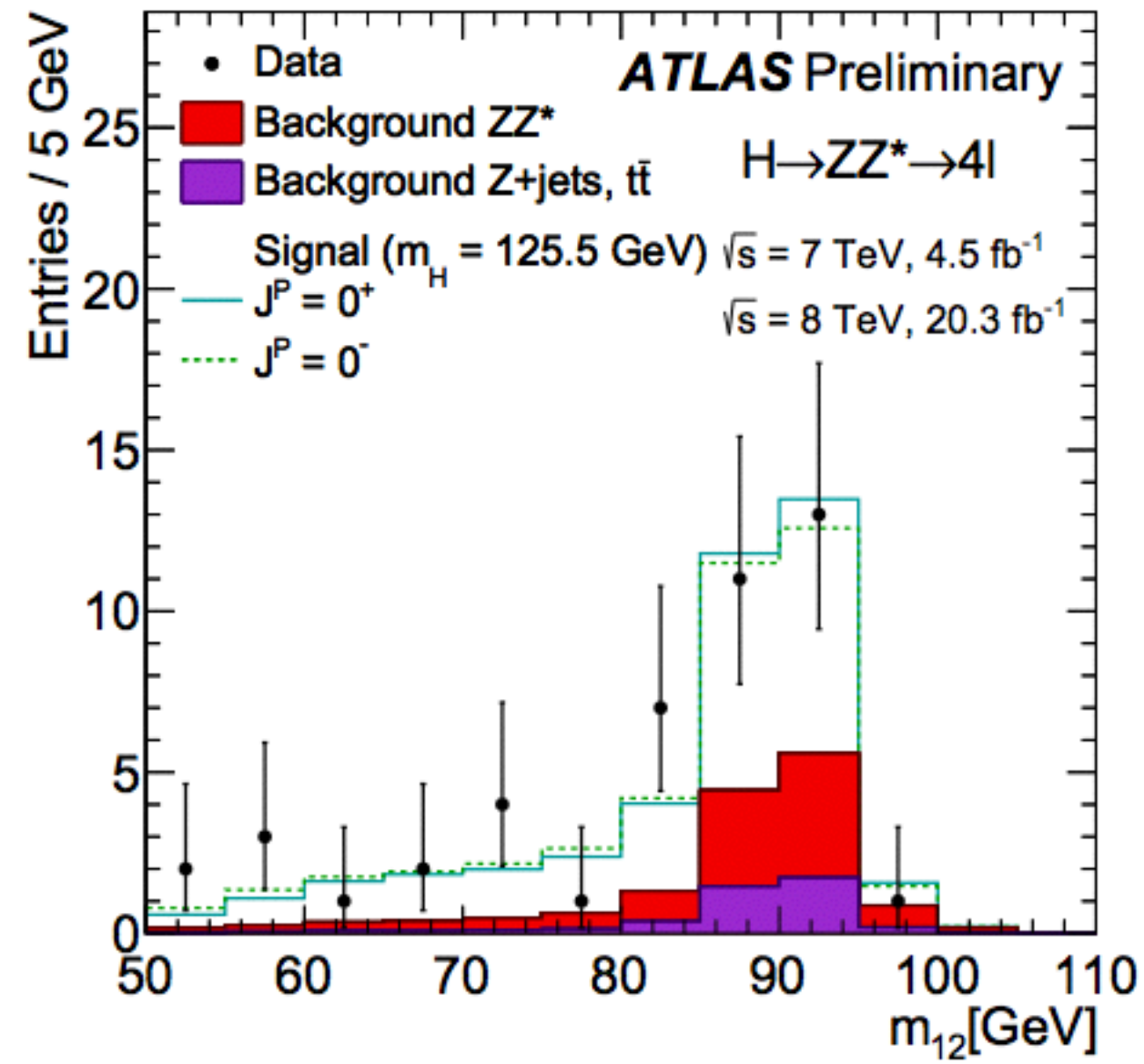
$$\theta_1 = \cos^{-1} \left(-\frac{\mathbf{q}_2 \cdot \mathbf{q}_{11}}{|\mathbf{q}_2| |\mathbf{q}_{11}|} \right)$$

$$\theta_2 = \cos^{-1} \left(-\frac{\mathbf{q}_1 \cdot \mathbf{q}_{21}}{|\mathbf{q}_1| |\mathbf{q}_{21}|} \right)$$

- other variables, p_T & η , are sensitive to NLO effects & insensitive to CP

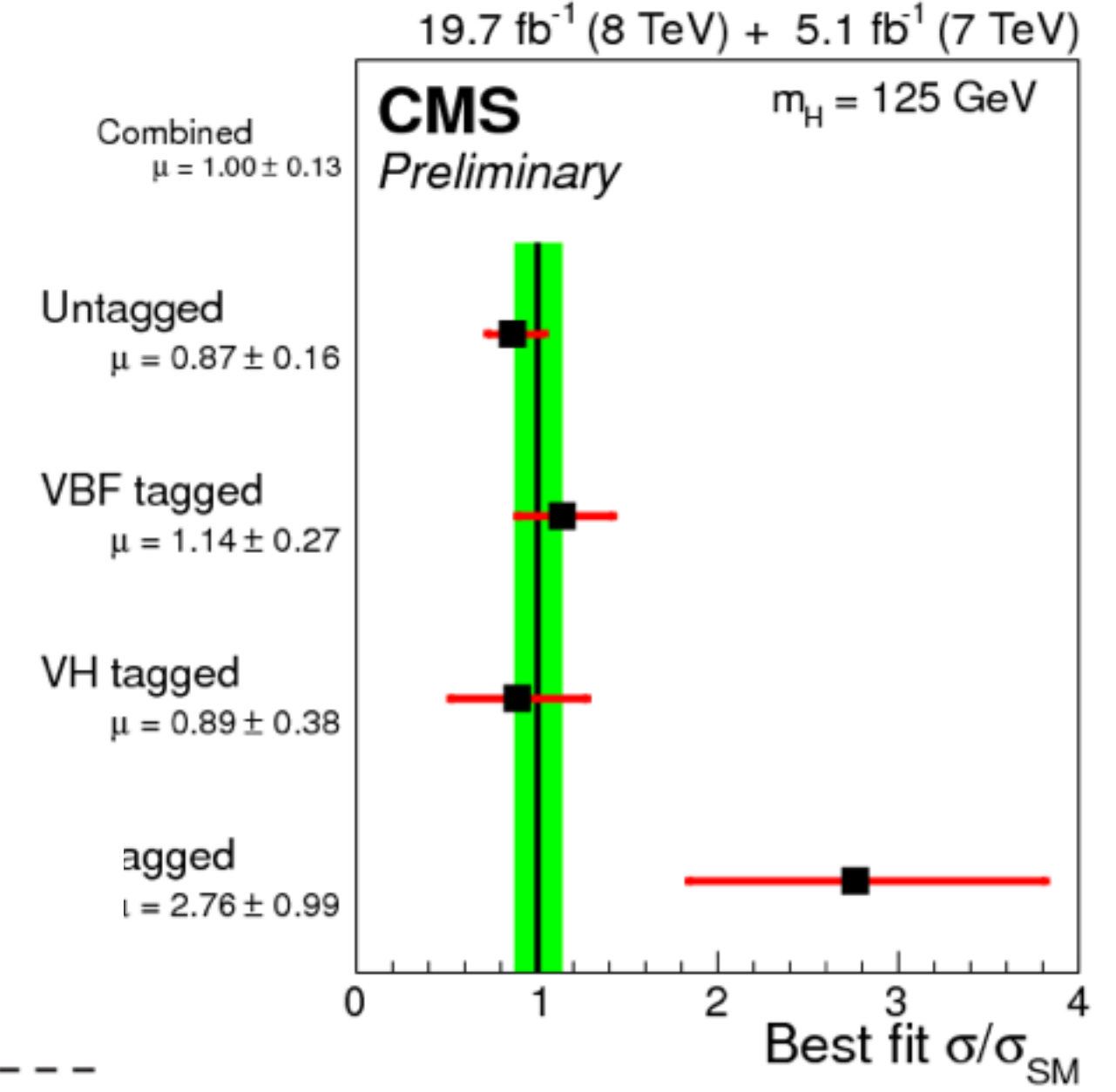
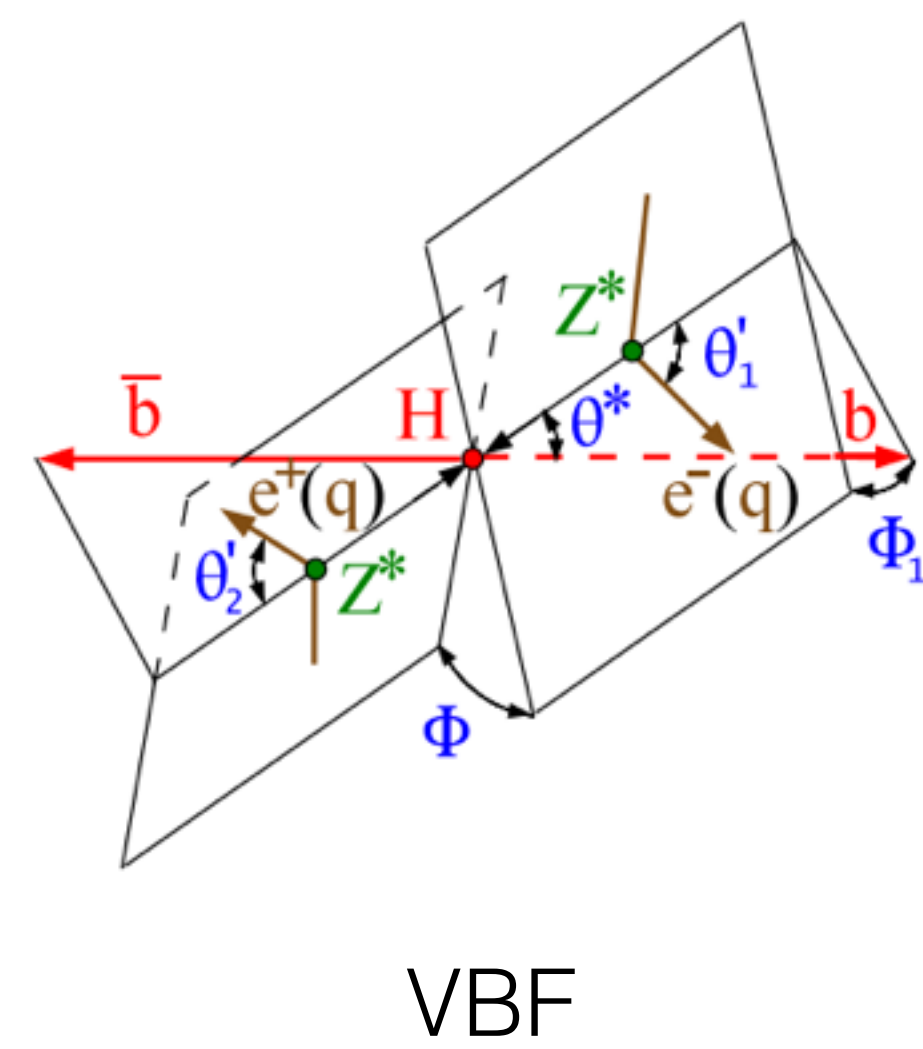
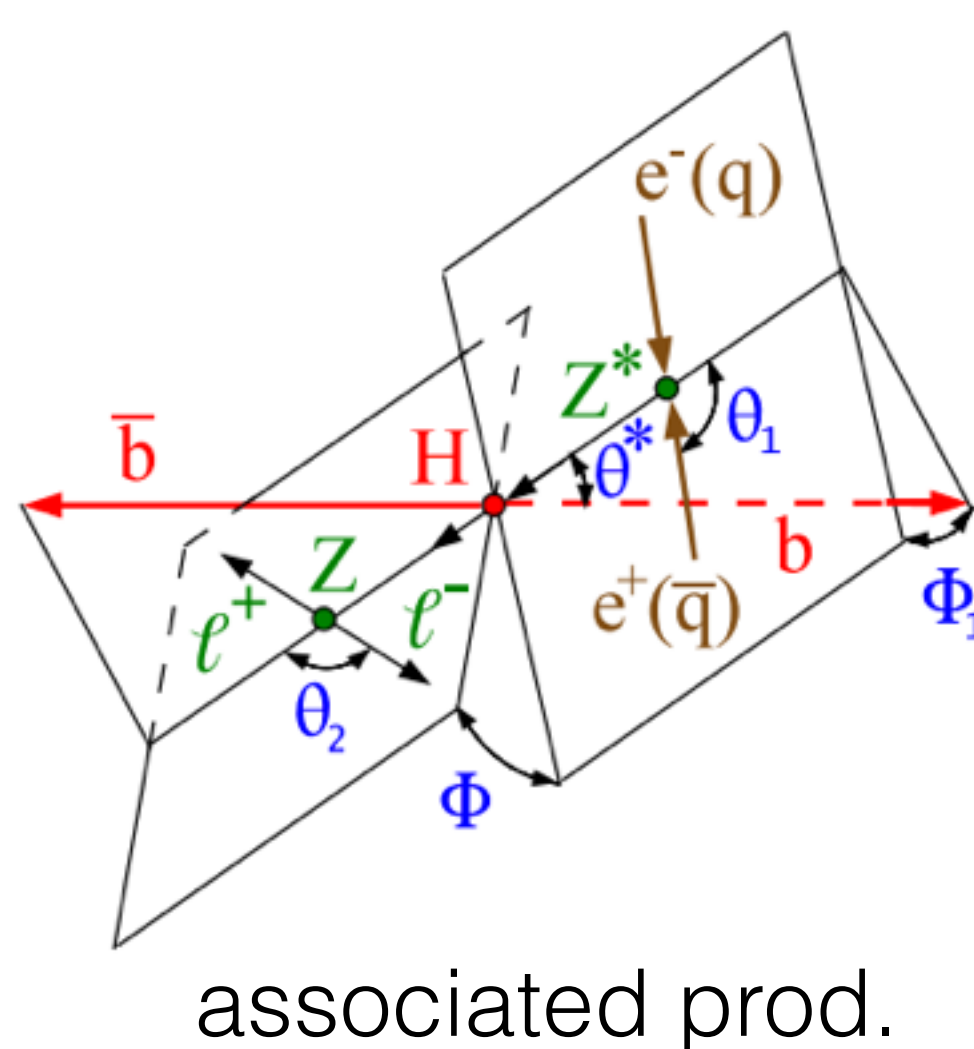
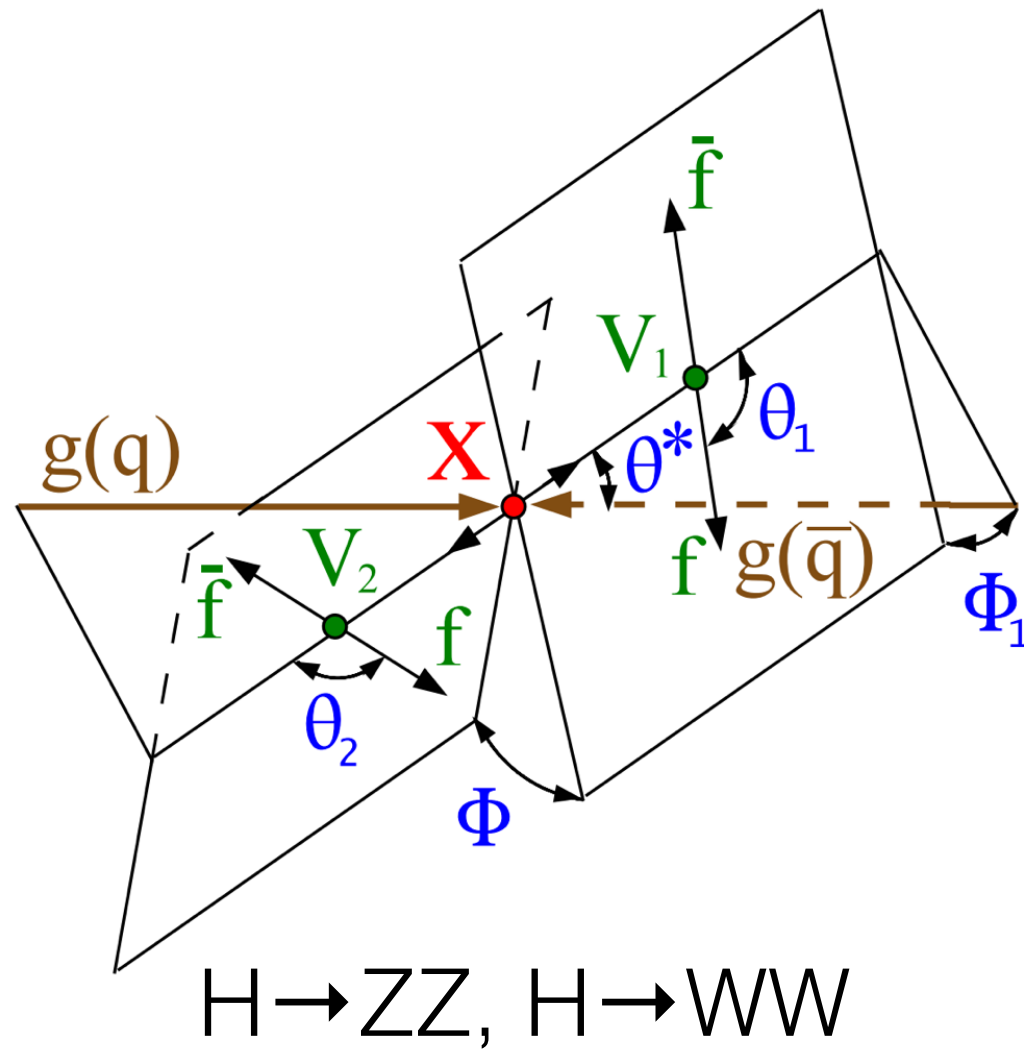
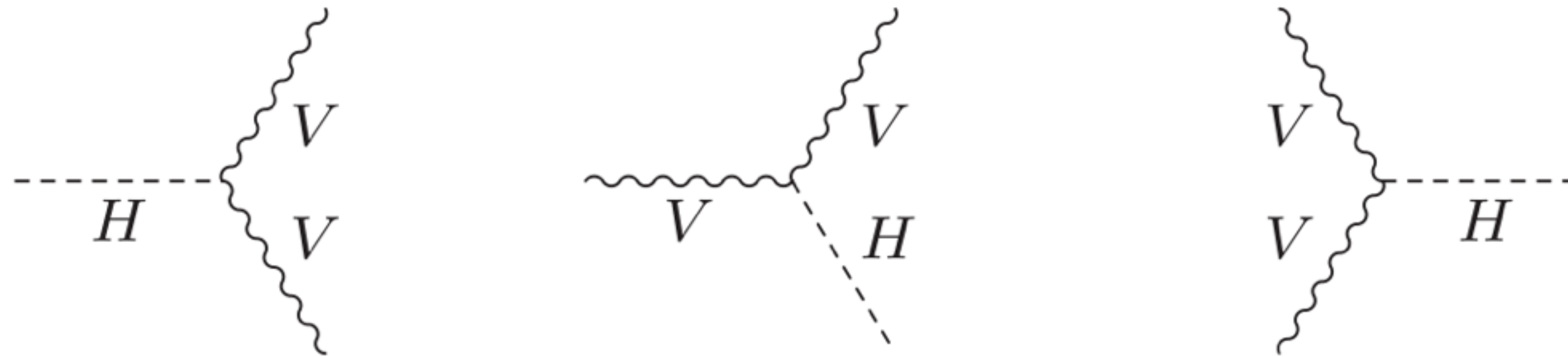


ATLAS distributions



Other production modes

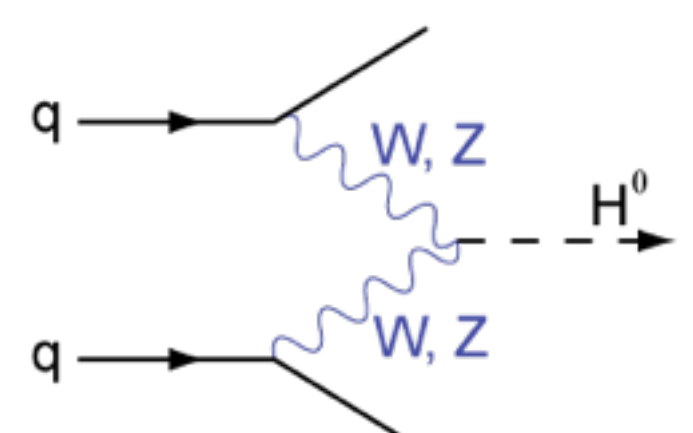
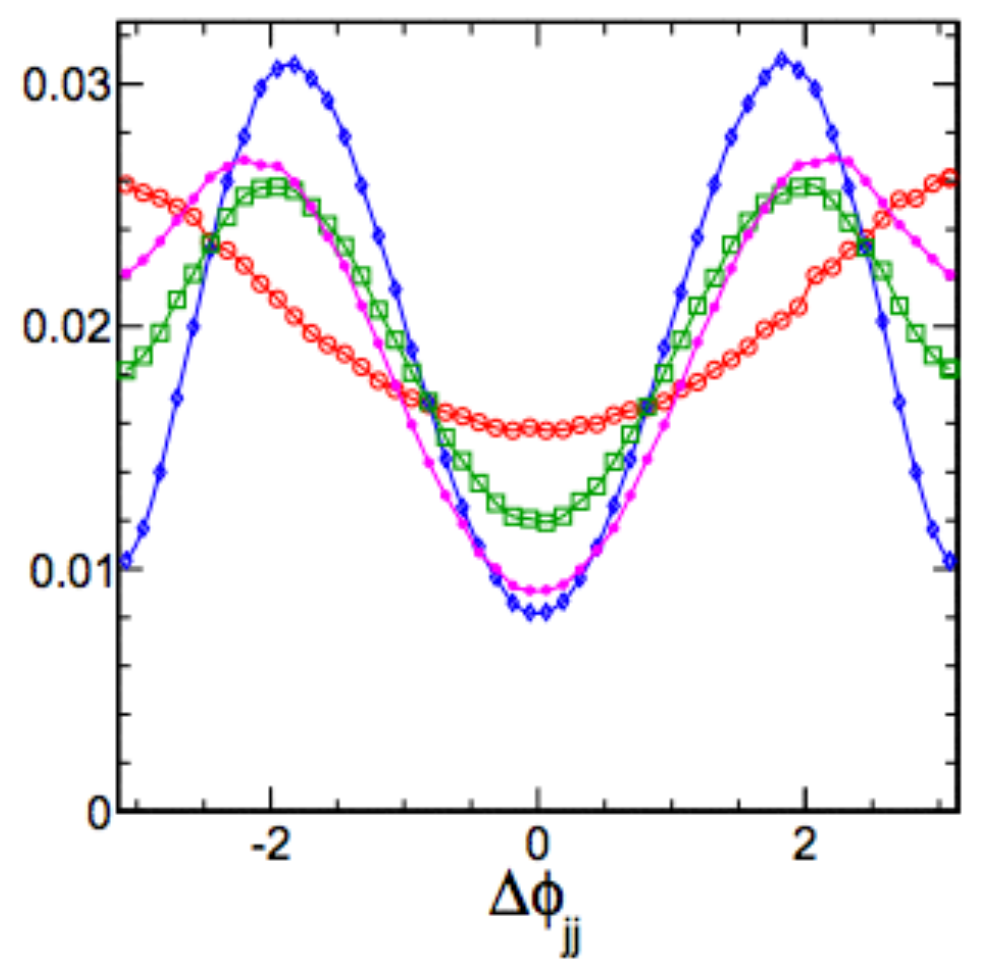
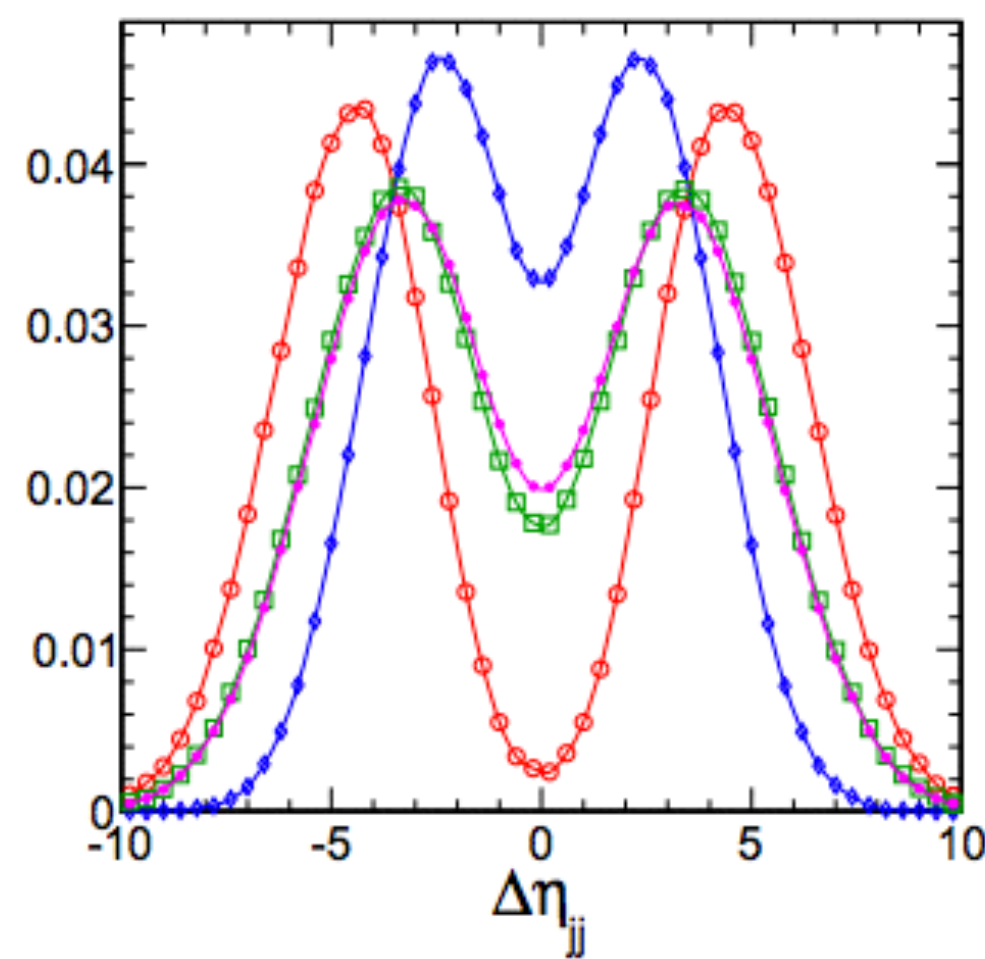
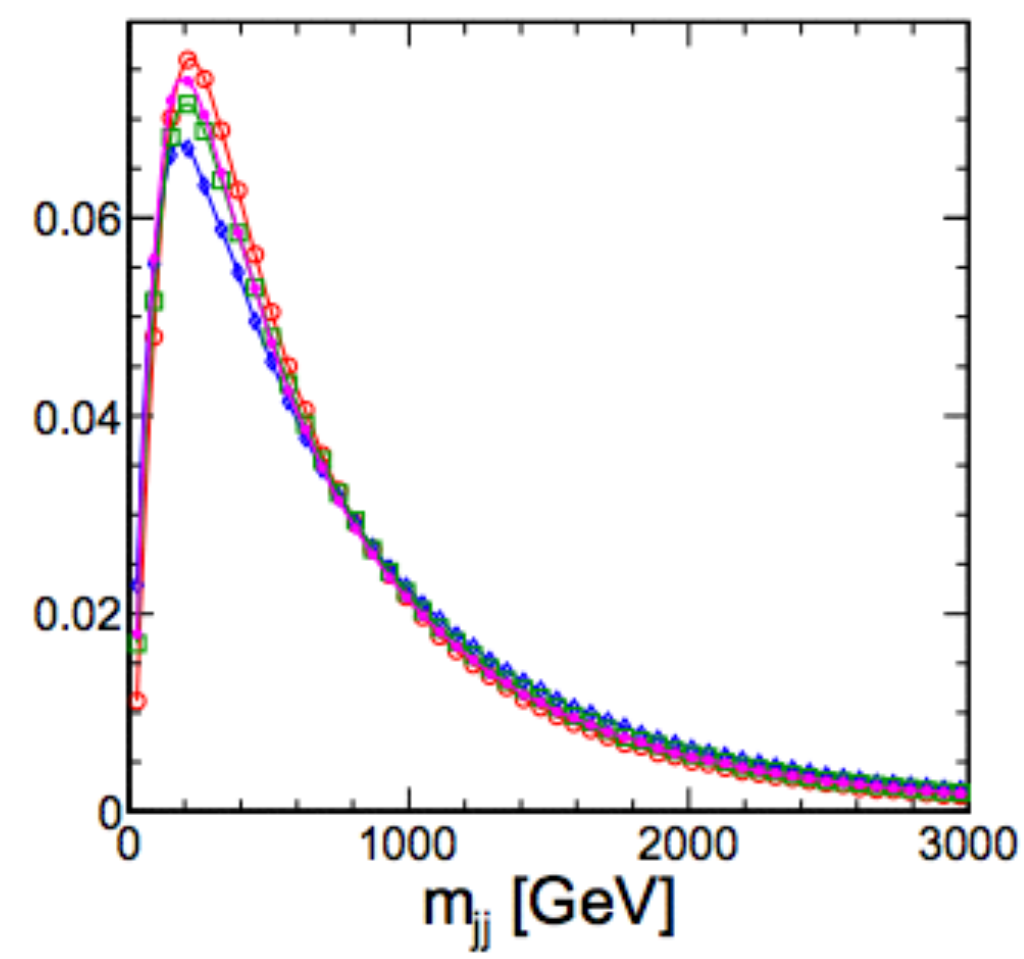
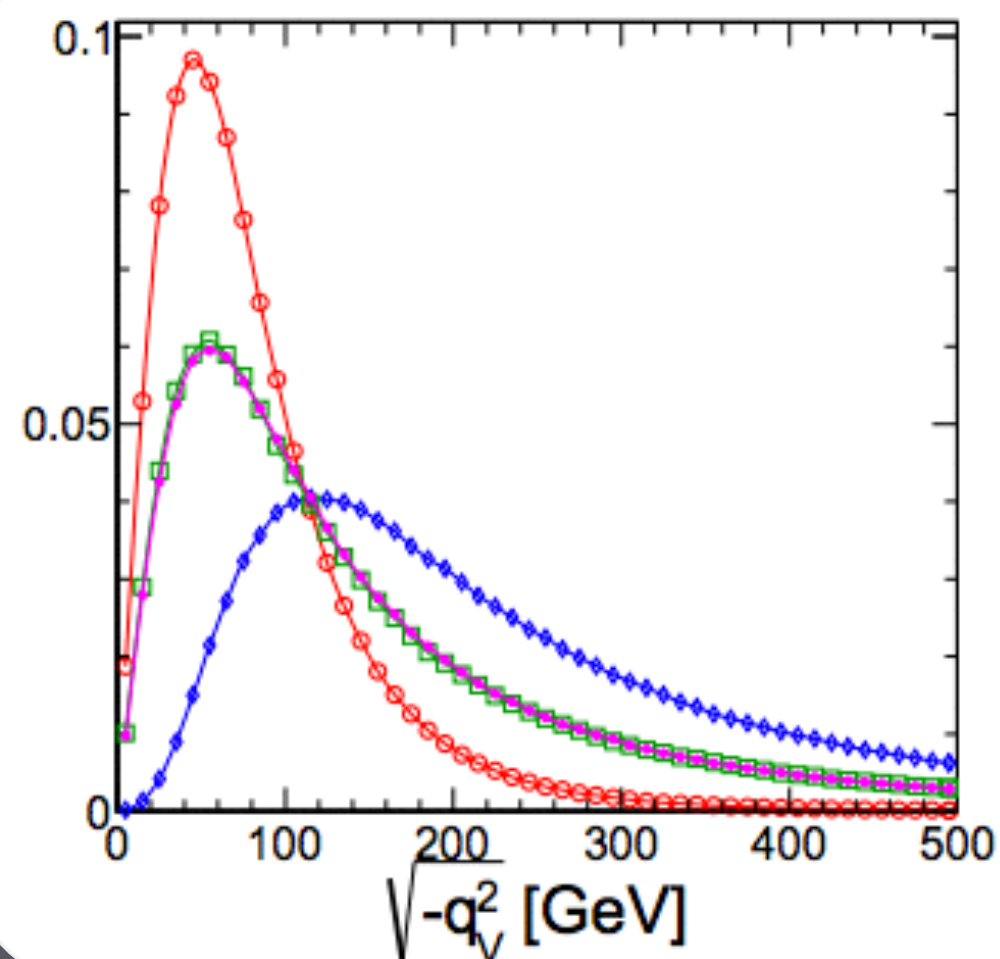
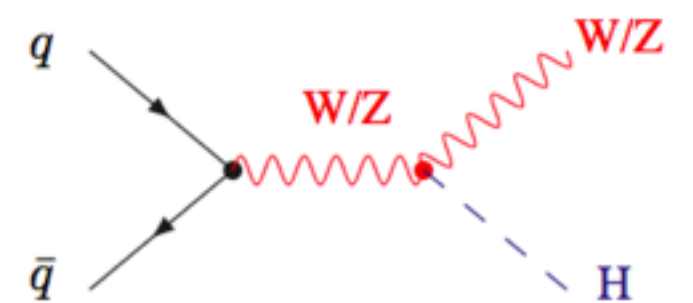
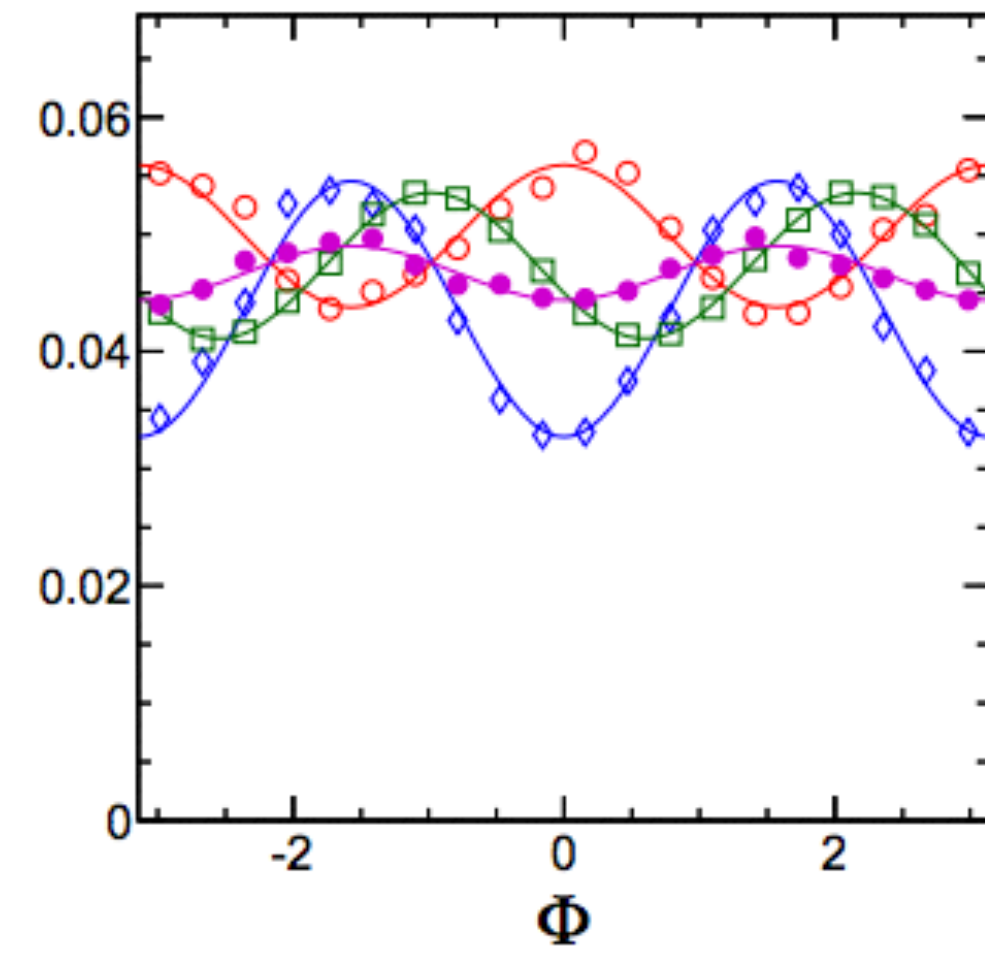
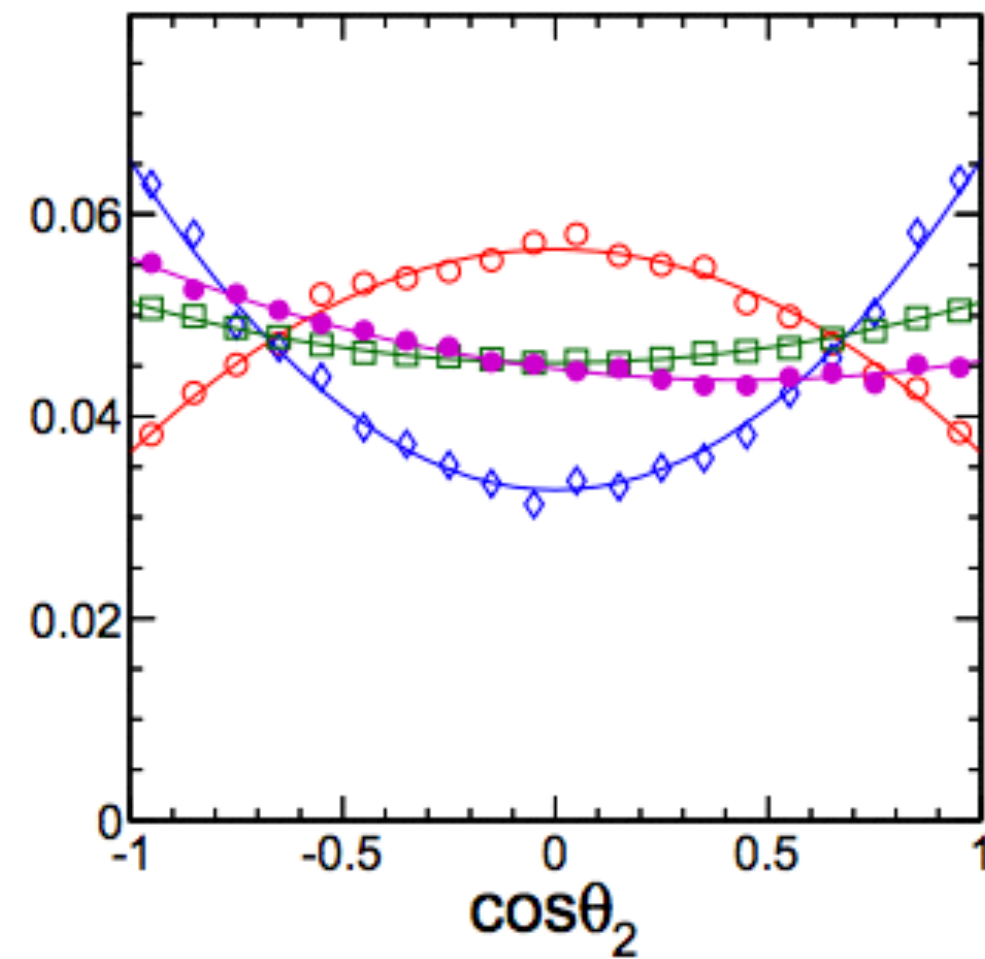
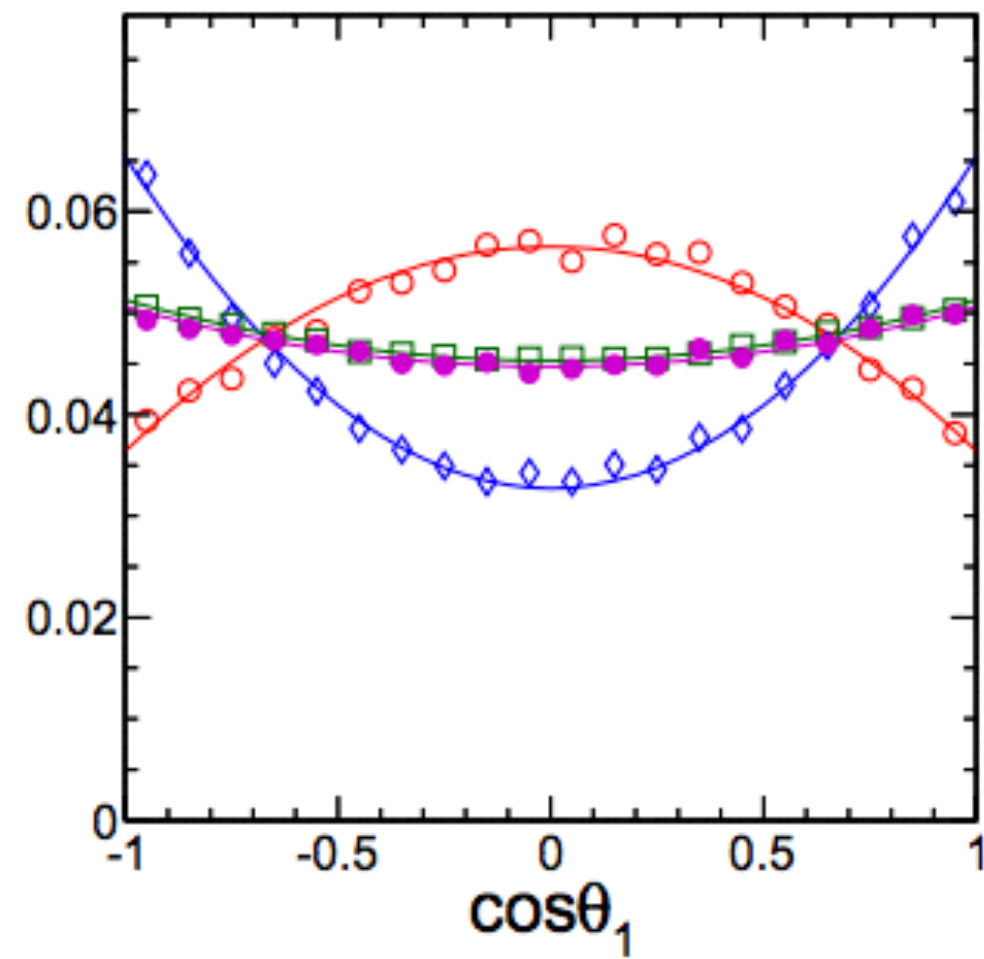
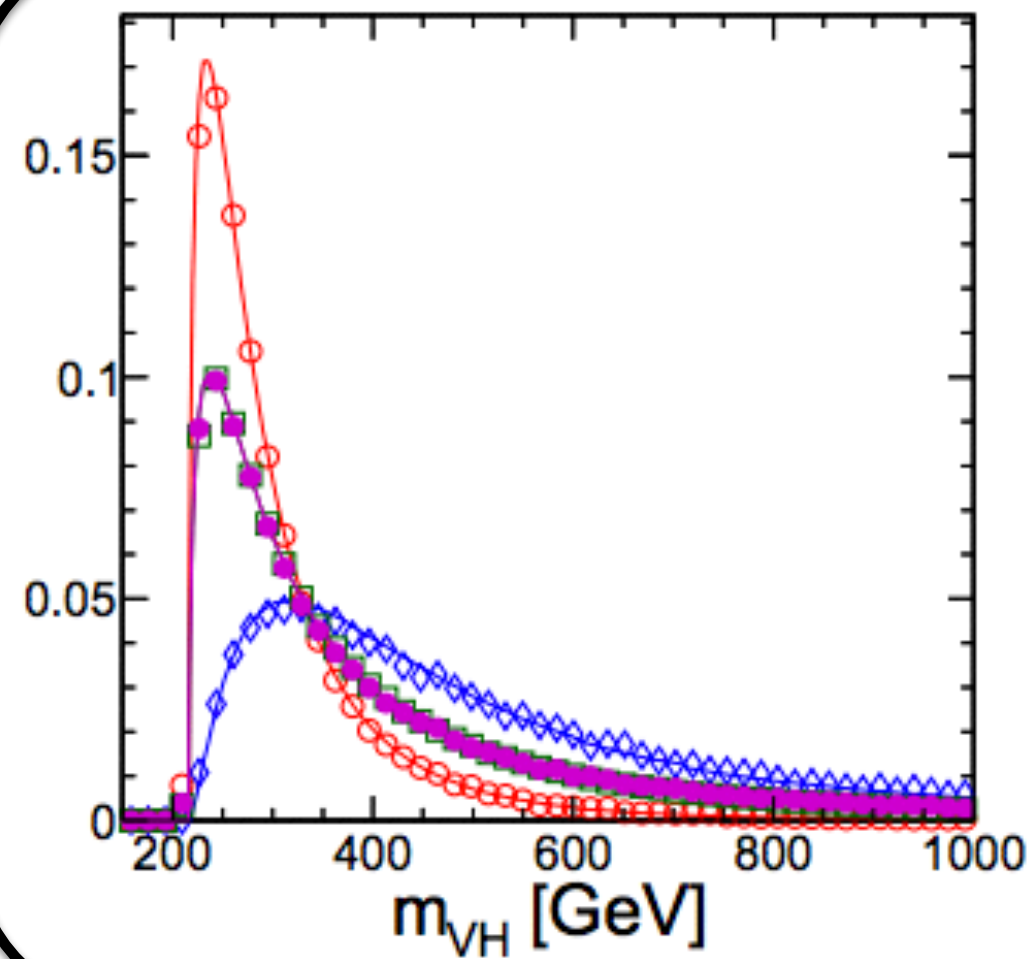
- As sensitivity to other production mechanisms become sensitive, new opportunities arise
- 0-like events have an enhancement at large \sqrt{s} - leading to increased sensitivity in associated production and VBF production channels



CMS-PAS-HIG-14-009

VH & VBF

- Kinematic distribution for $0+$ (red), $0-$ (blue), 50-50 mixture (green & magenta)



Prospects other production modes

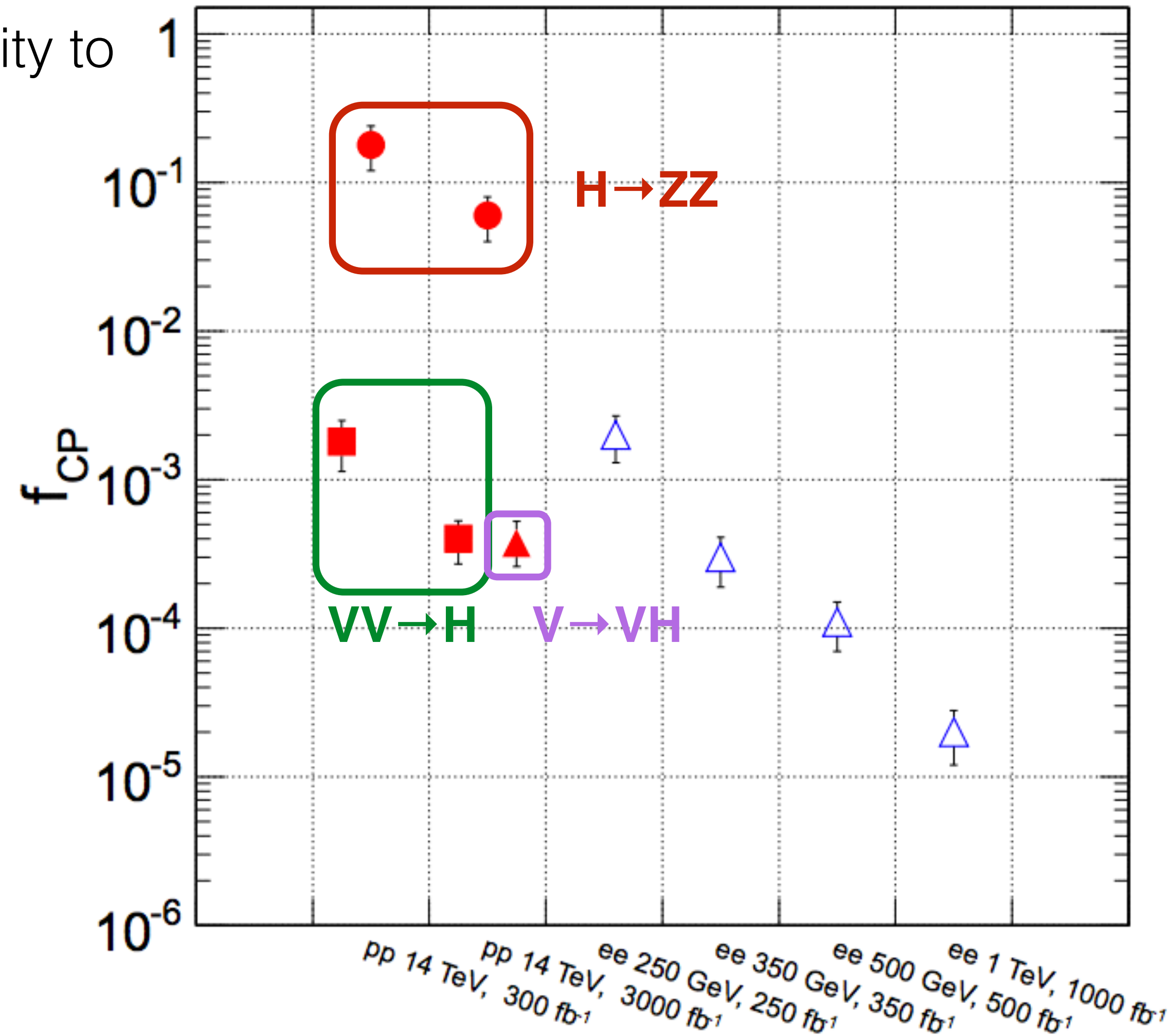
- Projections out to 300/fb & 3000/fb show strong sensitivity to CP-violating HVV interactions
- Associated production: assuming only $Z \rightarrow Z(\ell\ell)H(bb)$
- VBF production: assuming $H(\gamma\gamma) + H(ZZ)$

→ dominant sensitivity from associated and VBF production mechanisms!

$$f_{CP} = \frac{|a_3|^2 \sigma_3^{H \rightarrow ZZ}}{\sum |a_i|^2 \sigma_i^{H \rightarrow ZZ}}$$

e.g. $\frac{\sigma_3^{H \rightarrow ZZ}}{\sigma_1^{H \rightarrow ZZ}} \sim 0.160$

Summary of projections for discovery potential @ LHC & e+e collider



arXiv:1310.8361v2

ATLAS monte carlo

- signals modeled using *Powheg-Box Monte Carlo*
 - includes both ggF and VBF up to NLO in QCD
 - pT spectrum reweighted to NNLO and NNLL predictions

H → W(lν)W(lν)

Process	MC generator	$\sigma \cdot \mathcal{B}$ (pb)
Signal		
ggF $H \rightarrow WW^*$	POWHEG+PYTHIA8	0.435
VBF $H \rightarrow WW^*$	POWHEG+PYTHIA8	0.0356
VH $H \rightarrow WW^*$	PYTHIA8	0.0253
WW		
$q\bar{q} \rightarrow WW$ and $qg \rightarrow WW$	POWHEG+PYTHIA6	5.68
$gg \rightarrow WW$	GG2VV+HERWIG	0.196
$(q\bar{q} \rightarrow W) + (q\bar{q} \rightarrow W)$	PYTHIA8	0.480
$q\bar{q} \rightarrow WW$	SHERPA	5.68
VBS $WW + 2$ jets	SHERPA	0.0397
Top quarks		
$t\bar{t}$	POWHEG+PYTHIA6	26.6
Wt	POWHEG+PYTHIA6	2.35
$tq\bar{b}$	ACERMC+PYTHIA6	28.4
$t\bar{b}$	POWHEG+PYTHIA6	1.82
Other dibosons (VV)		
$W\gamma$ ($p_T^\gamma > 8$ GeV)	ALPGEN+HERWIG	369
$W\gamma^*$ ($m_{\ell\ell} \leq 7$ GeV)	SHERPA	12.2
WZ ($m_{\ell\ell} > 7$ GeV)	POWHEG+PYTHIA8	12.7
VBS $WZ + 2$ jets ($m_{\ell\ell} > 7$ GeV)	SHERPA	0.0126
$Z\gamma$ ($p_T^\gamma > 8$ GeV)	SHERPA	163
$Z\gamma^*$ (min. $m_{\ell\ell} \leq 4$ GeV)	SHERPA	7.31
ZZ ($m_{\ell\ell} > 4$ GeV)	POWHEG+PYTHIA8	0.733
$ZZ \rightarrow \ell\ell\nu\nu$ ($m_{\ell\ell} > 4$ GeV)	POWHEG+PYTHIA8	0.504
Drell-Yan		
Z ($m_{\ell\ell} > 10$ GeV)	ALPGEN+HERWIG	16500
VBF $Z + 2$ jets ($m_{\ell\ell} > 7$ GeV)	SHERPA	5.36

CMS discriminant details

$$\mathcal{P}_{\text{SM}} = \mathcal{P}_{\text{SM}}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{\text{sig}}^{\text{mass}}(m_{4\ell}|m_{\text{H}}),$$

$$\mathcal{P}_{J^P} = \mathcal{P}_{J^P}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{\text{sig}}^{\text{mass}}(m_{4\ell}|m_{\text{H}}),$$

$$\mathcal{P}_{J^P}^{\text{int}} = \left(\mathcal{P}_{\text{SM}+J^P}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) - \mathcal{P}_{J^P}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) - \mathcal{P}_{\text{SM}}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \right),$$

$$\mathcal{P}_{J^P}^{\text{int}\perp} = \left(\mathcal{P}_{\text{SM}+J^P\perp}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) - \mathcal{P}_{J^P}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) - \mathcal{P}_{\text{SM}}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \right),$$

$$\mathcal{P}_{q\bar{q}ZZ} = \mathcal{P}_{q\bar{q}ZZ}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{q\bar{q}ZZ}^{\text{mass}}(m_{4\ell}),$$

$$\mathcal{P}_{g\bar{g}ZZ} = \mathcal{P}_{g\bar{g}ZZ}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{g\bar{g}ZZ}^{\text{mass}}(m_{4\ell}),$$

$$\mathcal{D}_{\text{bkg}} = \frac{\mathcal{P}_{\text{SM}}}{\mathcal{P}_{\text{SM}} + c \times \mathcal{P}_{q\bar{q}ZZ}} = \left[1 + c(m_{4\ell}) \times \frac{\mathcal{P}_{q\bar{q}ZZ}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{q\bar{q}ZZ}^{\text{mass}}(m_{4\ell})}{\mathcal{P}_{\text{SM}}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell}) \times \mathcal{P}_{\text{sig}}^{\text{mass}}(m_{4\ell}|m_{\text{H}})} \right]^{-1},$$

$$\mathcal{D}_{J^P} = \frac{\mathcal{P}_{\text{SM}}}{\mathcal{P}_{\text{SM}} + \mathcal{P}_{J^P}} = \left[1 + \frac{\mathcal{P}_{J^P}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell})}{\mathcal{P}_{\text{SM}}^{\text{kin}}(m_1, m_2, \vec{\Omega}|m_{4\ell})} \right]^{-1},$$

$$\mathcal{D}_{\text{int}} = \frac{\mathcal{P}_{J^P}^{\text{int}}(m_1, m_2, \vec{\Omega}|m_{4\ell})}{\mathcal{P}_{\text{SM}}^{\text{kin}} + \mathcal{P}_{J^P}^{\text{kin}}}.$$

ATLAS discriminant details

$$O_1(\kappa_{HV V}) = \frac{2\Re(ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)^* \cdot ME(\kappa_{HV V} \neq 0; \kappa_{SM}, \kappa_{AV V} = 0; \alpha = 0))}{|ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)|^2},$$

$$O_2(\kappa_{HV V}) = \frac{|ME(\kappa_{HV V} \neq 0; \kappa_{SM}, \kappa_{AV V} = 0; \alpha = 0)|^2}{|ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)|^2},$$

$$O_1(\kappa_{AV V}, \alpha) = \frac{2\Re(ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)^* \cdot ME(\kappa_{AV V} \neq 0; \kappa_{SM}, \kappa_{HV V} = 0; \alpha = \pi/2))}{|ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)|^2},$$

$$O_2(\kappa_{AV V}, \alpha) = \frac{|ME(\kappa_{AV V} \neq 0; \kappa_{SM}, \kappa_{HV V} = 0; \alpha = \pi/2)|^2}{|ME(\kappa_{SM} \neq 0; \kappa_{HV V}, \kappa_{AV V} = 0; \alpha = 0)|^2}.$$

CMS summary

Parameter	Observed	Expected	$f_{ai}^{VV} = 1$
$f_{\Lambda 1} \cos(\phi_{\Lambda 1})$	$0.22^{+0.10}_{-0.16} [-0.25, 0.37]$	$0.00^{+0.16}_{-0.87} [-1.00, 0.27]$ $\cup [0.92, 1.00]$	1.1% (16%)
$f_{a2} \cos(\phi_{a2})$	$0.00^{+0.41}_{-0.06} [-0.66, -0.57]$ $\cup [-0.15, 1.00]$	$0.00^{+0.38}_{-0.08} [-0.18, 1.00]$	5.2% (5.0%)
$f_{a3} \cos(\phi_{a3})$	$0.00^{+0.14}_{-0.11} [-0.40, 0.43]$	$0.00^{+0.33}_{-0.33} [-0.70, 0.70]$	0.02% (0.41%)
$f_{\Lambda 1}^{WW} \cos(\phi_{\Lambda 1}^{WW})$	$0.21^{+0.18}_{-1.21} [-1.00, 1.00]$	$0.00^{+0.34}_{-1.00} [-1.00, 0.41]$ $\cup [0.49, 1.00]$	78% (67%)
$f_{a2}^{WW} \cos(\phi_{a2}^{WW})$	$-0.02^{+1.02}_{-0.16} [-1.00, -0.54]$ $\cup [-0.29, 1.00]$	$0.00^{+1.00}_{-0.12} [-1.00, -0.58]$ $\cup [-0.22, 1.00]$	42% (46%)
$f_{a3}^{WW} \cos(\phi_{a3}^{WW})$	$-0.03^{+1.03}_{-0.97} [-1.00, 1.00]$	$0.00^{+1.00}_{-1.00} [-1.00, 1.00]$	34% (49%)
$f_{\Lambda 1}^{Z\gamma} \cos(\phi_{\Lambda 1}^{Z\gamma})$	$-0.27^{+0.34}_{-0.49} [-1.00, 1.00]$	$0.00^{+0.83}_{-0.53} [-1.00, 1.00]$	26% (16%)
$f_{a2}^{Z\gamma} \cos(\phi_{a2}^{Z\gamma})$	$0.00^{+0.14}_{-0.20} [-0.49, 0.46]$	$0.00^{+0.51}_{-0.51} [-0.78, 0.79]$	<0.01% (0.01%)
$f_{a3}^{Z\gamma} \cos(\phi_{a3}^{Z\gamma})$	$0.02^{+0.21}_{-0.13} [-0.40, 0.51]$	$0.00^{+0.51}_{-0.51} [-0.75, 0.75]$	<0.01% (<0.01%)
$f_{a2}^{\gamma\gamma} \cos(\phi_{a2}^{\gamma\gamma})$	$0.12^{+0.20}_{-0.11} [-0.04, +0.51]$	$0.00^{+0.11}_{-0.09} [-0.32, 0.34]$	<0.01% (<0.01%)
$f_{a3}^{\gamma\gamma} \cos(\phi_{a3}^{\gamma\gamma})$	$-0.02^{+0.06}_{-0.13} [-0.35, 0.32]$	$0.00^{+0.15}_{-0.11} [-0.37, 0.40]$	<0.01% (<0.01%)

Cross sections for f_{ai} Conversion

Interaction	Anomalous Coupling	Coupling Phase	Effective Fraction	Translation Constant
HZZ	Λ_1	$\phi_{\Lambda 1}$	$f_{\Lambda 1}$	$\sigma_1 / \tilde{\sigma}_{\Lambda 1} = 1.45 \times 10^4 \text{ TeV}^{-4}$
	a_2	$\phi_{a 2}$	$f_{a 2}$	$\sigma_1 / \sigma_2 = 2.68$
	a_3	$\phi_{a 3}$	$f_{a 3}$	$\sigma_1 / \sigma_3 = 6.36$
HWW	Λ_1^{WW}	$\phi_{\Lambda 1}^{WW}$	$f_{\Lambda 1}^{WW}$	$\sigma_1^{WW} / \tilde{\sigma}_{\Lambda 1}^{WW} = 1.87 \times 10^4 \text{ TeV}^{-4}$
	a_2^{WW}	$\phi_{a 2}^{WW}$	$f_{a 2}^{WW}$	$\sigma_1^{WW} / \sigma_2^{WW} = 1.25$
	a_3^{WW}	$\phi_{a 3}^{WW}$	$f_{a 3}^{WW}$	$\sigma_1^{WW} / \sigma_3^{WW} = 3.01$
	$\Lambda_1^{Z\gamma}$	$\phi_{\Lambda 1}^{Z\gamma}$	$f_{\Lambda 1}^{Z\gamma}$	$\sigma_1' / \tilde{\sigma}_{\Lambda 1}^{Z\gamma} = 5.76 \times 10^3 \text{ TeV}^{-4}$
HZ γ	$a_2^{Z\gamma}$	$\phi_{a 2}^{Z\gamma}$	$f_{a 2}^{Z\gamma}$	$\sigma_1' / \sigma_2^{Z\gamma} = 22.4 \times 10^{-4}$
	$a_3^{Z\gamma}$	$\phi_{a 3}^{Z\gamma}$	$f_{a 3}^{Z\gamma}$	$\sigma_1' / \sigma_3^{Z\gamma} = 27.2 \times 10^{-4}$
H $\gamma\gamma$	$a_2^{\gamma\gamma}$	$\phi_{a 2}^{\gamma\gamma}$	$f_{a 2}^{\gamma\gamma}$	$\sigma_1' / \sigma_2^{\gamma\gamma} = 28.2 \times 10^{-4}$
	$a_3^{\gamma\gamma}$	$\phi_{a 3}^{\gamma\gamma}$	$f_{a 3}^{\gamma\gamma}$	$\sigma_1' / \sigma_3^{\gamma\gamma} = 28.8 \times 10^{-4}$