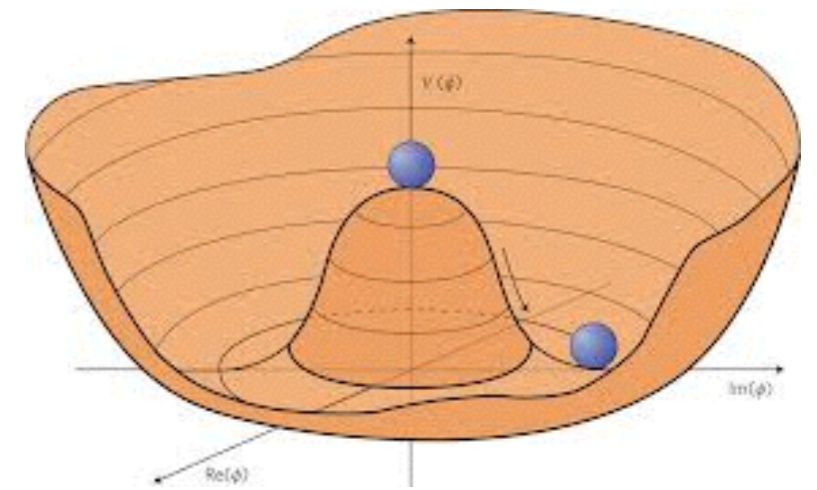


# di-Higgs at the LHC

Michael Spannowsky

IPPP, Durham University

- We have the remnant of elw. symmetry breaking
- Now we want to finally study the mechanism/potential



$$\mathcal{L} = (D_\mu \Phi)^\dagger (D^\mu \Phi) - V(\Phi^\dagger \Phi) \quad V(\Phi^\dagger \Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

- Studying quartic impossible at envisioned FCs
- One of best reasons for phase 2 upgrade:  
if ILC does not go to 1 TeV might not outperform LHC
- Jose talked about HH in Higgs portal context
- To contrast BSM measurement SM needs to be understood first. This talk SM HH and non-HP BSM

# Higgs self-coupling measurements in the Standard Model

$$\begin{aligned} -\mathcal{L} \supset & \frac{1}{2}m_h^2 h^2 + \sqrt{\frac{\eta}{2}}m_h h^3 + \frac{\eta}{4}h^4 \\ & - gm_V V^2 h - \frac{m_f}{v} \bar{f} f h \\ & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

# Higgs self-coupling measurements in the Standard Model

$$-\mathcal{L} \supset \frac{1}{2}m_h^2 h^2 + \sqrt{\frac{\eta}{2}}m_h h^3 + \frac{\eta}{4}h^4$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

$$- gm_V V^2 h - \frac{m_f}{v} \bar{f} f h$$
$$- \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v)$$

# Higgs self-coupling measurements in the Standard Model

$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2}m_h^2 h^2 + \sqrt{\frac{\eta}{2}}m_h h^3 + \frac{\eta}{4}h^4 \longrightarrow \text{Potential needs at least} \\
 & \text{dihiggs production!} \\
 & - gm_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v)
 \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

# Higgs self-coupling measurements in the Standard Model

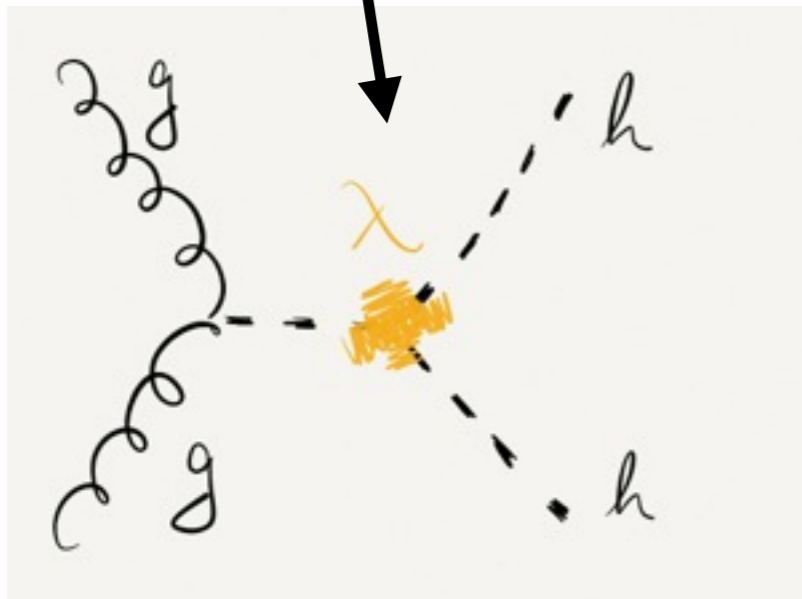
$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4 \longrightarrow \text{Potential needs at least} \\
 & \text{dihiggs production!} \\
 & - g m_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \\
 & = - \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} h + \frac{\alpha_s}{24\pi v^2} G_{\mu\nu}^a G^{a\mu\nu} h^2 + \dots
 \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

# Higgs self-coupling measurements in the Standard Model

$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4 \longrightarrow \text{Potential needs at least} \\
 & \text{dihiggs production!} \\
 & - gm_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \\
 & = - \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} h + \frac{\alpha_s}{24\pi v^2} G_{\mu\nu}^a G^{a\mu\nu} h^2 + \dots
 \end{aligned}$$

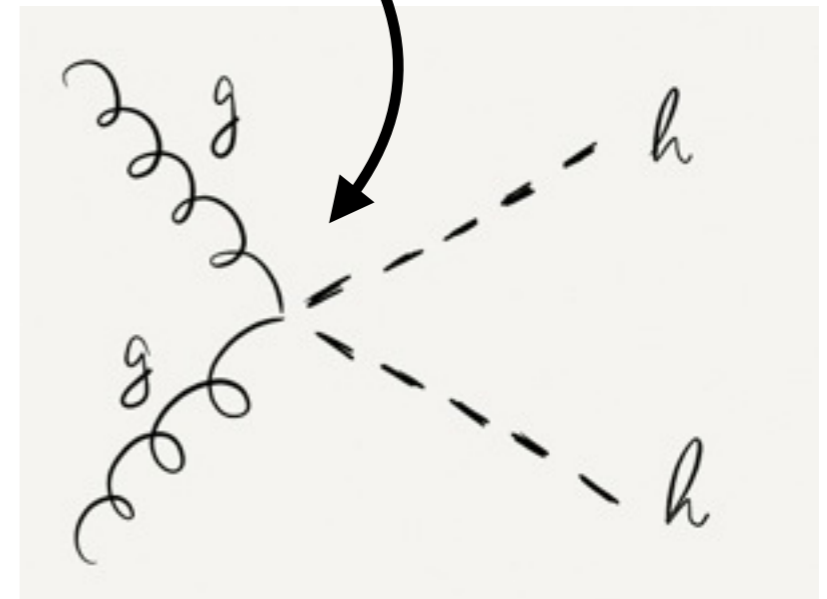
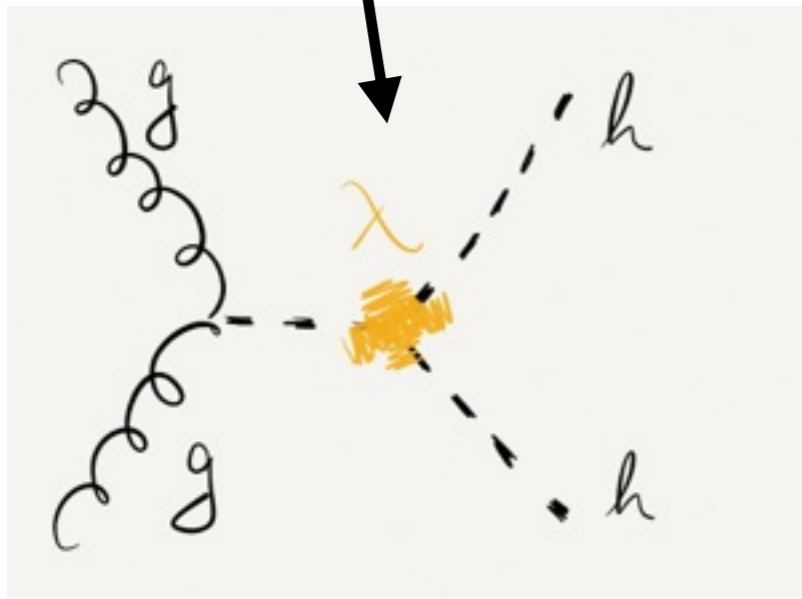
$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$



# Higgs self-coupling measurements in the Standard Model

$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4 \longrightarrow \text{Potential needs at least} \\
 & \text{dihiggs production!} \\
 & - gm_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \\
 & = - \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} h + \frac{\alpha_s}{24\pi v^2} G_{\mu\nu}^a G^{a\mu\nu} h^2 + \dots
 \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

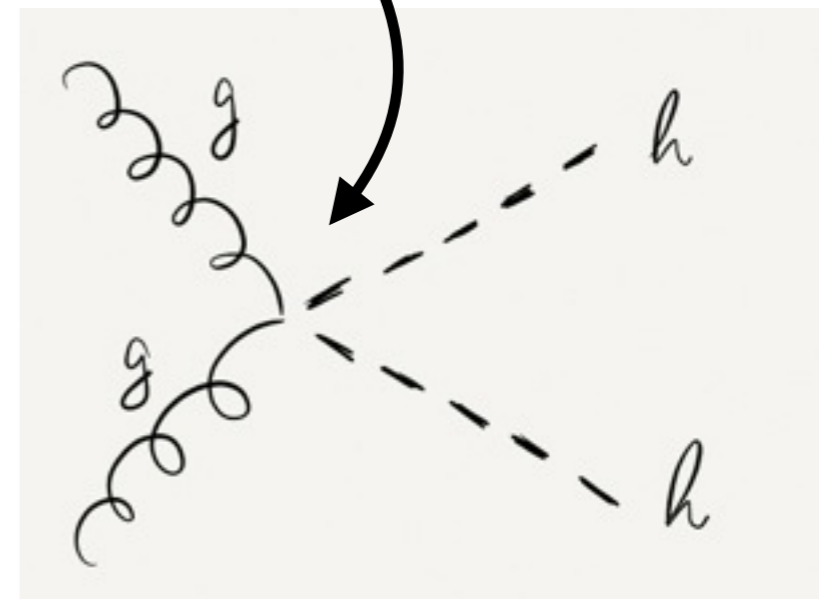
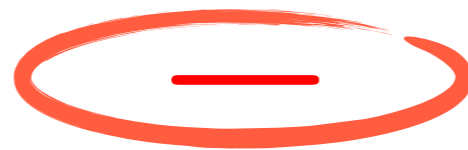
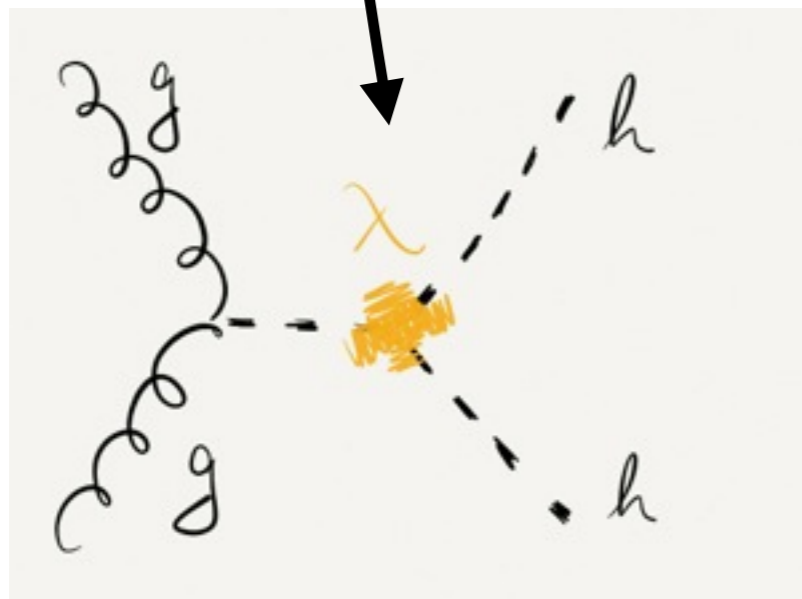




# Higgs self-coupling measurements in the Standard Model

$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4 \longrightarrow \text{Potential needs at least} \\
 & \text{dihiggs production!} \\
 & - g m_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \\
 & = - \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} h + \frac{\alpha_s}{24\pi v^2} G_{\mu\nu}^a G^{a\mu\nu} h^2 + \dots
 \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

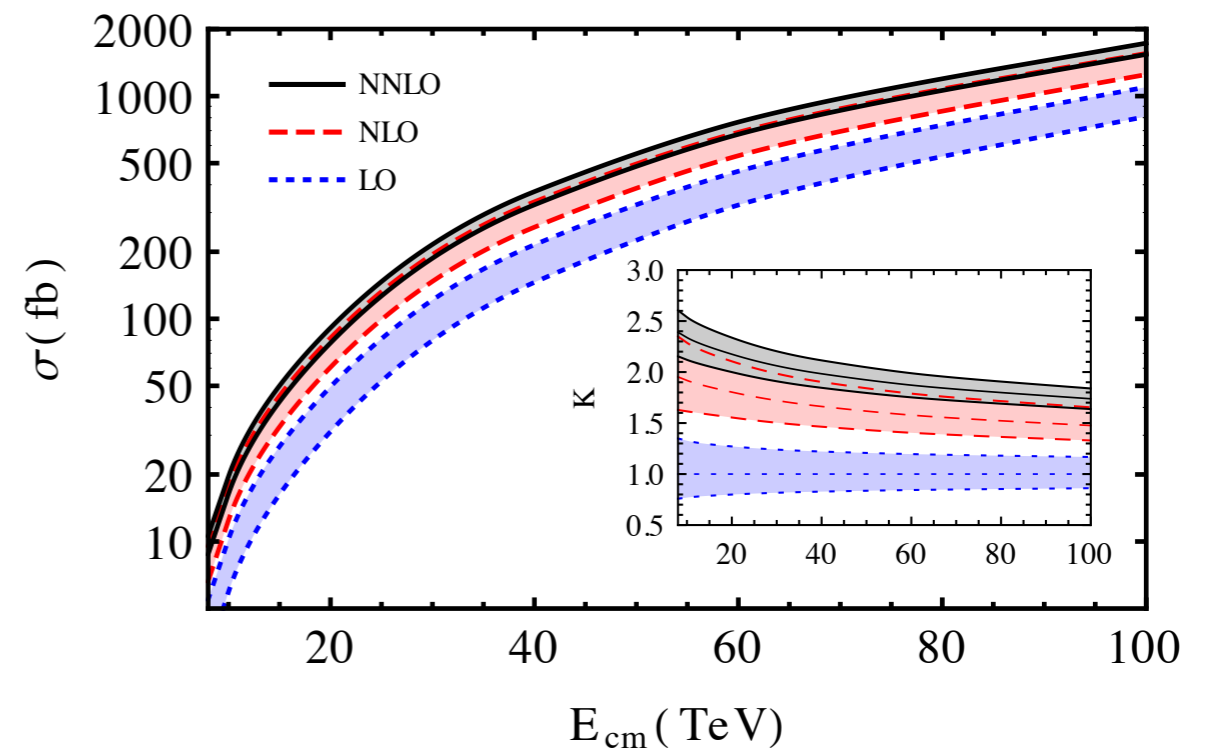
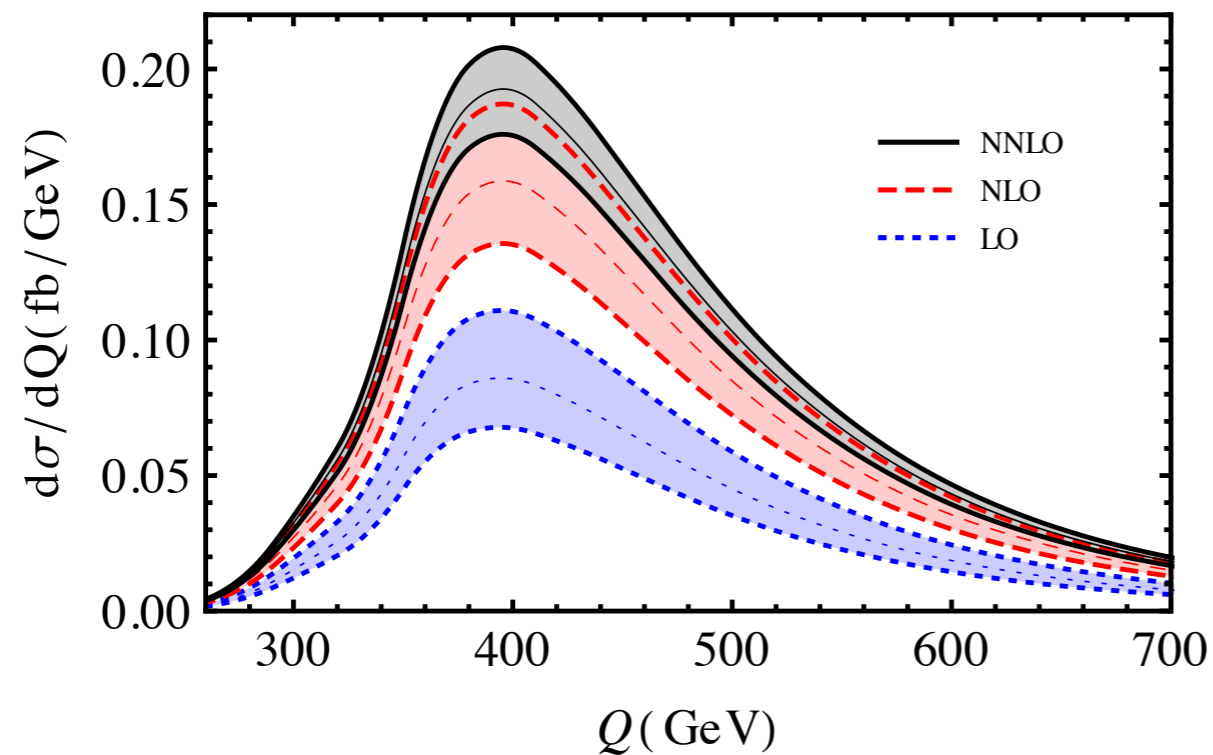


# Recent progress in HH cross section calculations

<b>HH</b>	LO full mt	[Plehn, Spira, Zerwas Nucl. Phys. B479, hep-ph 9603205] and now many others...
	NLO eff. mt	[Dawson, Dittmaier, Spira PRD 58 1998] implemented in HPair <a href="http://people.web.psi.ch/spira/hpair/">http://people.web.psi.ch/spira/hpair/</a>
	NLO full mt	[Grigo, Hoff, Melnikov, Steinhauser 1311.7425]
	NNLO eff. mt	[De Florian, Mazzitelli PRL 111, 1309.6594]
<b>HHj</b>	LO full mt	[Dolan, Englert, MS JHEP 1210, 1206.5001] [Li, Yan, Zhao 1312.3830] [Maierhoefer, Papaefstathiou 1401.0007]
	LO full mt (reweighted)	[Dolan, Englert, Greiner, MS 1310.1084]

# Higgs selfcoupling in HH+X

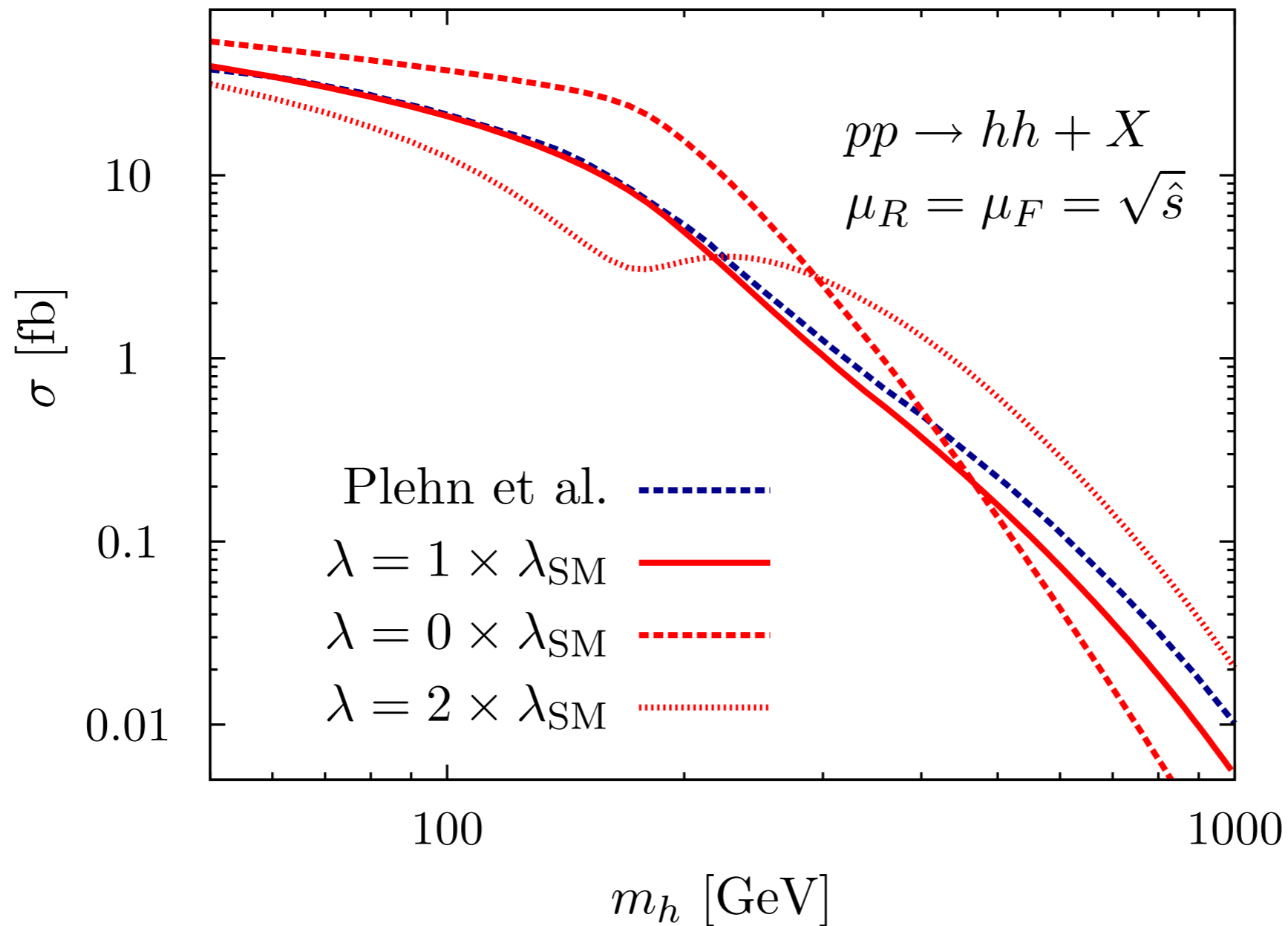
[De Florian, Mazzitelli PRL 111, 1309.6594]



$E_{cm}$	8 TeV	14 TeV	33 TeV	100 TeV
$\sigma_{\text{NNLO}}$	9.76 fb	40.2 fb	243 fb	1638 fb
Scale [%]	+9.0 – 9.8	+8.0 – 8.7	+7.0 – 7.4	+5.9 – 5.8
PDF [%]	+6.0 – 6.1	+4.0 – 4.0	+2.5 – 2.6	+2.3 – 2.6
PDF+ $\alpha_s$ [%]	+9.3 – 8.8	+7.2 – 7.1	+6.0 – 6.0	+5.8 – 6.0

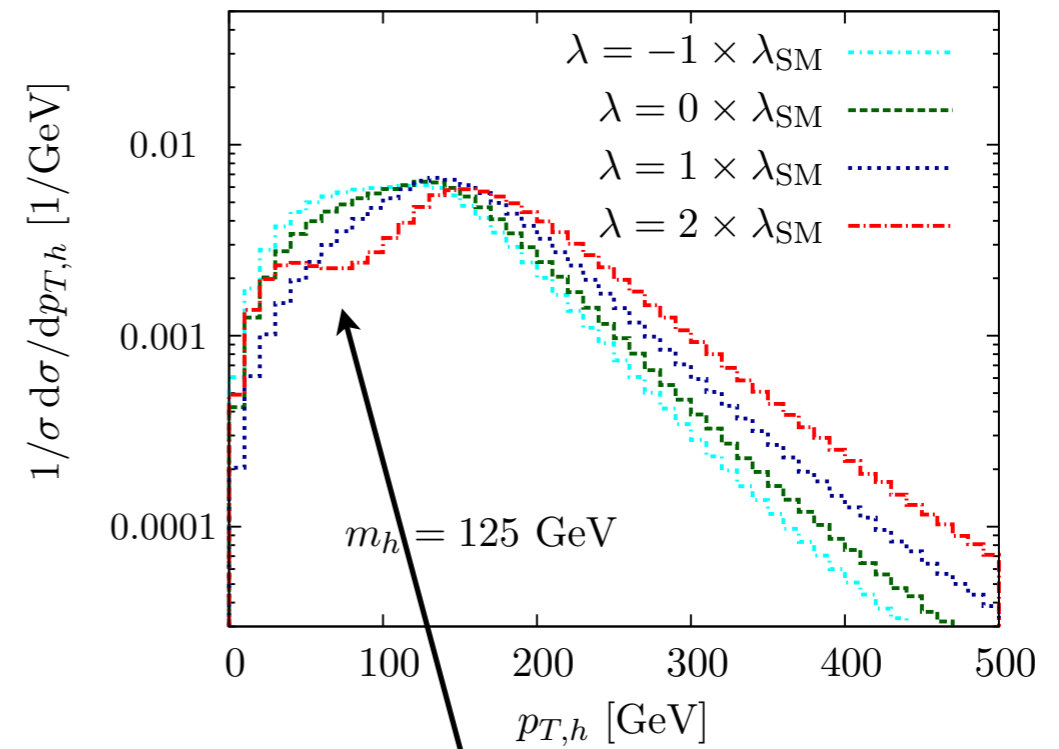
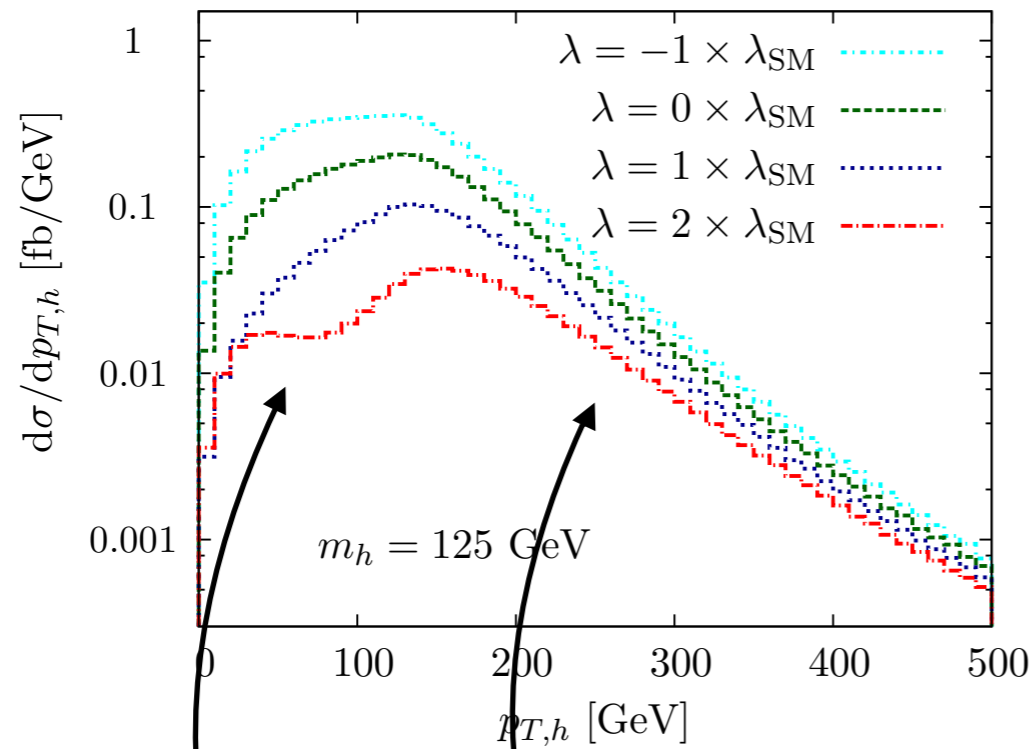
- NNLO in effective ggH
- Very large k-factors

# Higgs selfcoupling in $HH+X$



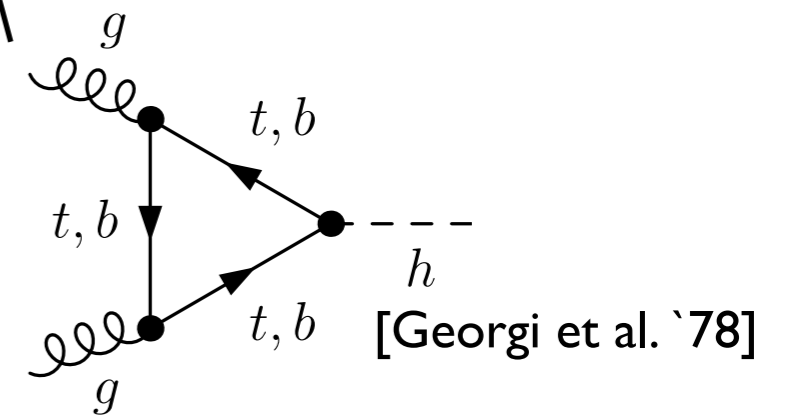
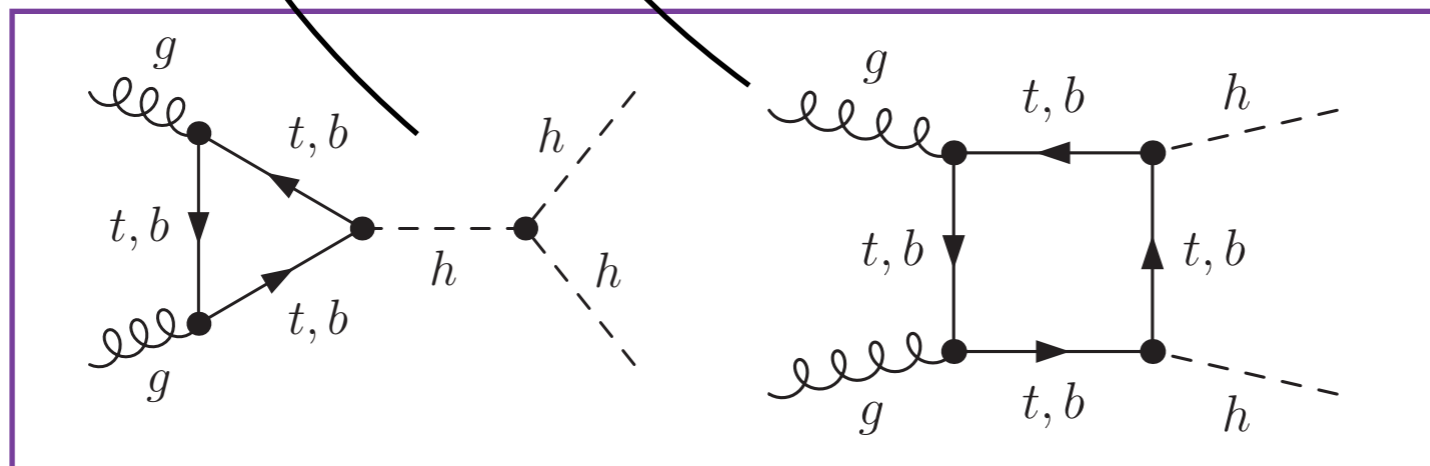
A priori good sensitivity for  $m_h = 125$  GeV

# Higgs selfcoupling in HH+X



has maximum contribution for

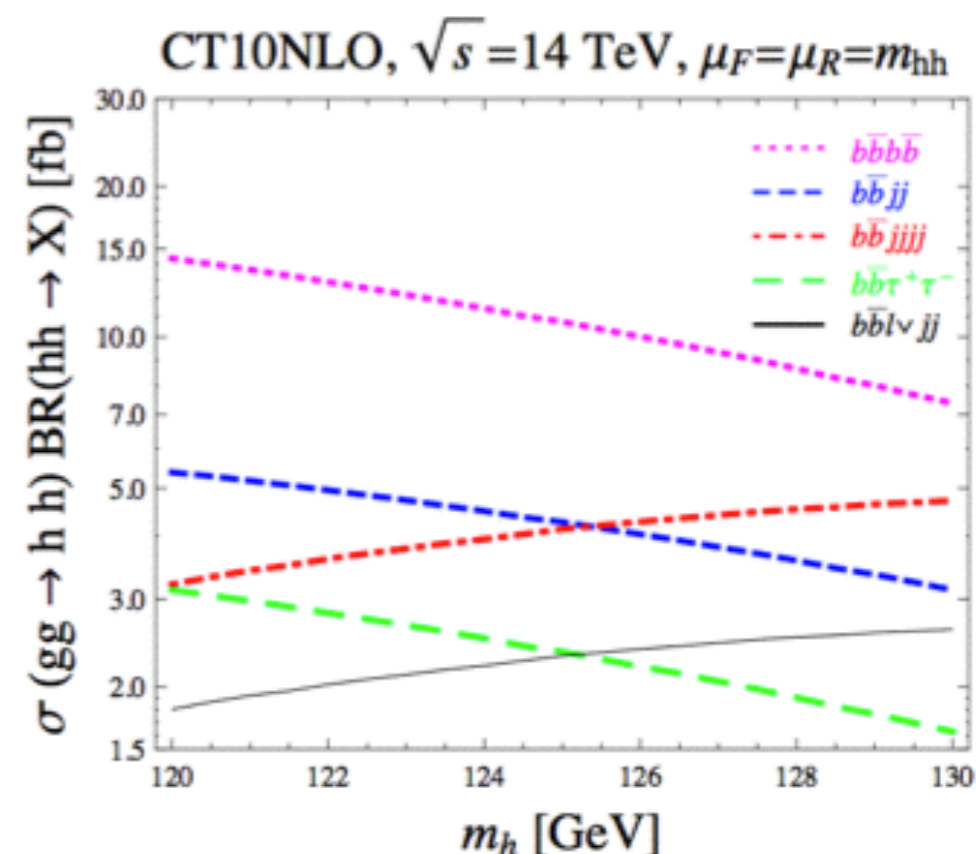
$$s = (p_{h,1} + p_{h,2})^2 = 4m_t^2$$



## Where is sensitivity located?

Measuring this small cross section in an inclusive search is very challenging at the HL-LHC: compromise between branching ratio and cleanliness of the signal

Channel	BR (%)	Events/3 ab
$bbWW$	24.7	30000
$bb\tau\tau$	7.3	9000
$WWWW$	4.3	5200
$bb\gamma\gamma$	0.27	330
$bbZZ(\rightarrow e^+e^-\mu^+\mu^-)$	0.015	19
$\gamma\gamma\gamma\gamma$	0.00052	1

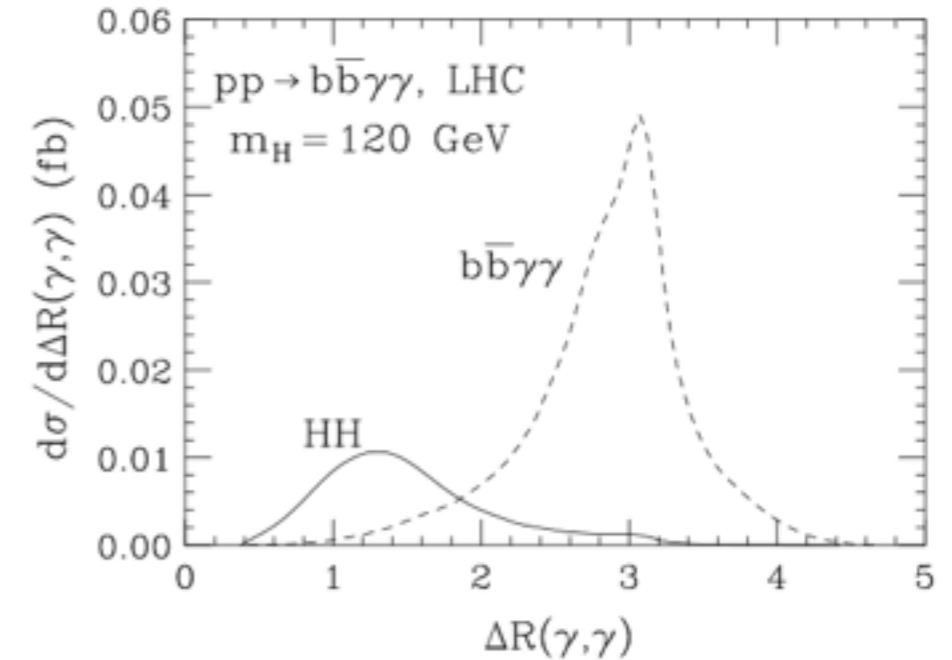


Several channels are currently under study by the collaborations

# Possible final states:

$\bar{b}b\gamma\gamma$  :

- Early on considered most sensitive
- Problem, difficult to simulate backgrounds:  
large reducible backgrounds (jets→photons)  
multi-jets have not been included



- Possibly side-band analysis might help
  - Still best analysis but parton level
- [Baur, Pehn, Rainwater '03]

cuts:

$$\begin{aligned}
 p_T(b) &> 45 \text{ GeV}, & |\eta(b)| < 2.5, & \Delta R(b, b) > 0.4, \\
 m_H - 20 \text{ GeV} &< m_{b\bar{b}} < m_H + 20 \text{ GeV}, \\
 p_T(\gamma) &> 20 \text{ GeV}, & |\eta(\gamma)| < 2.5, & \Delta R(\gamma, \gamma) > 0.4, \\
 m_H - 2.3 \text{ GeV} &< m_{\gamma\gamma} < m_H + 2.3 \text{ GeV}, \\
 \Delta R(\gamma, b) &> 0.4,
 \end{aligned}$$

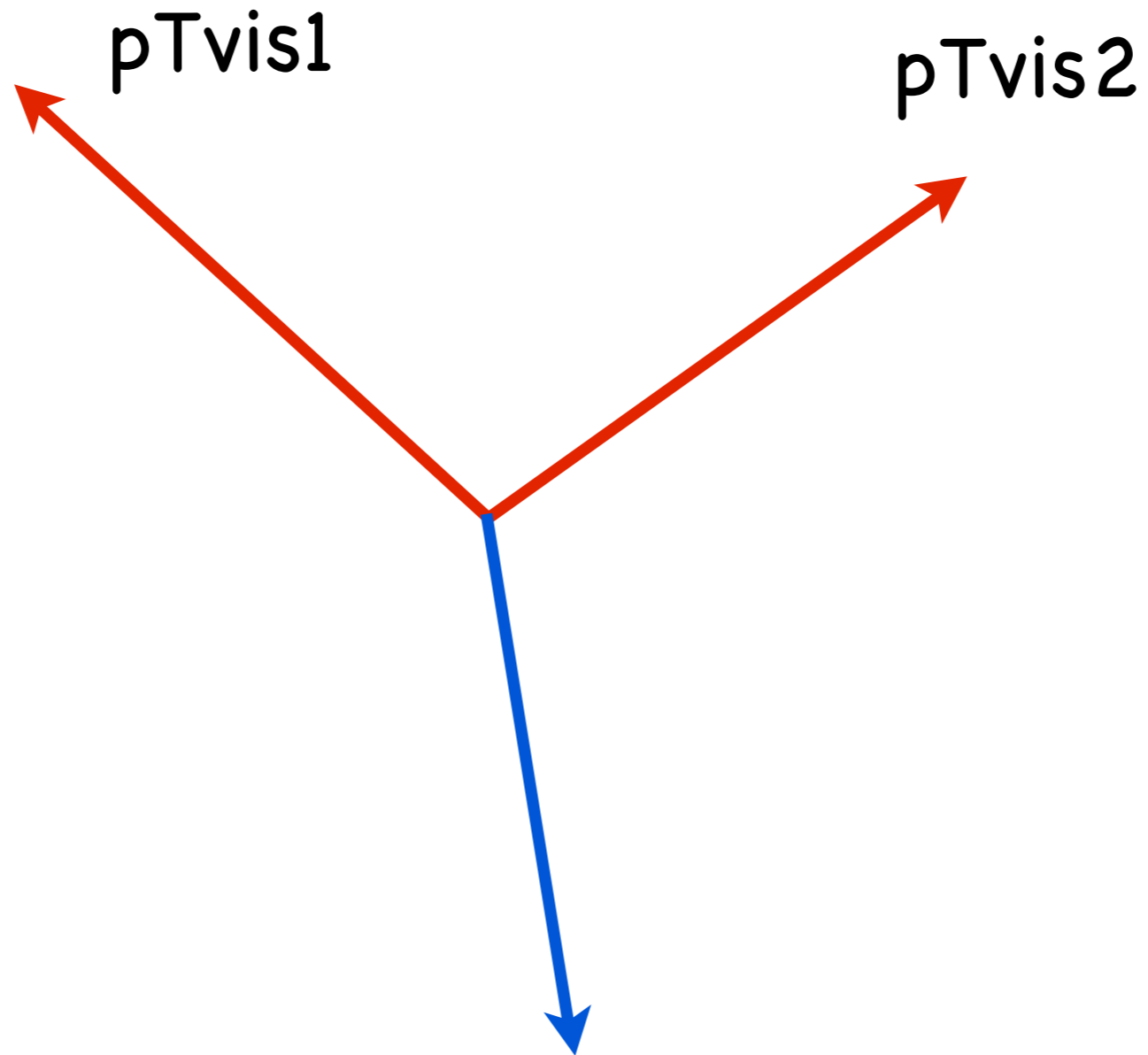
analysis stage	$HH$	$b\bar{b}\gamma\gamma$	$c\bar{c}\gamma\gamma$	$b\bar{b}\gamma j$	$c\bar{c}\gamma j$	$jj\gamma\gamma$	$b\bar{b}jj$	$c\bar{c}jj$	$\gamma jjj$	$jjjj$	$\Sigma(\text{bkg})$
before cuts	0.15	-	-	-	-	-	-	-	-	-	-
+ Eq. (3)	0.043	0.056	0.42	65	250	11	$2.5 \times 10^4$	$2.5 \times 10^4$	7700	$5 \times 10^6$	$5 \times 10^6$
+ Eq. (4)	0.035	0.0060	0.0215	8.28	17.0	0.84	4520	4520	364	$4 \times 10^5$	$4 \times 10^5$
$\times \epsilon \cdot P_{LHC}^{hi}$	0.0106	0.0029	0.0020	0.0031	0.0013	0.0077	0.0013	0.0003	0.0030	0.0022	0.0233
$N_{LHC} (hi)$	6	2	1	2	1	5	1	0	2	1	14
$\times \epsilon \cdot P_{LHC}^{lo}$	0.0106	0.0029	0.0020	0.0020	0.0008	0.0077	0.0005	0.0001	0.0017	0.0009	0.0186
$N_{LHC} (lo)$	6	2	1	1	0	5	0	0	1	1	11

## Possible final states:

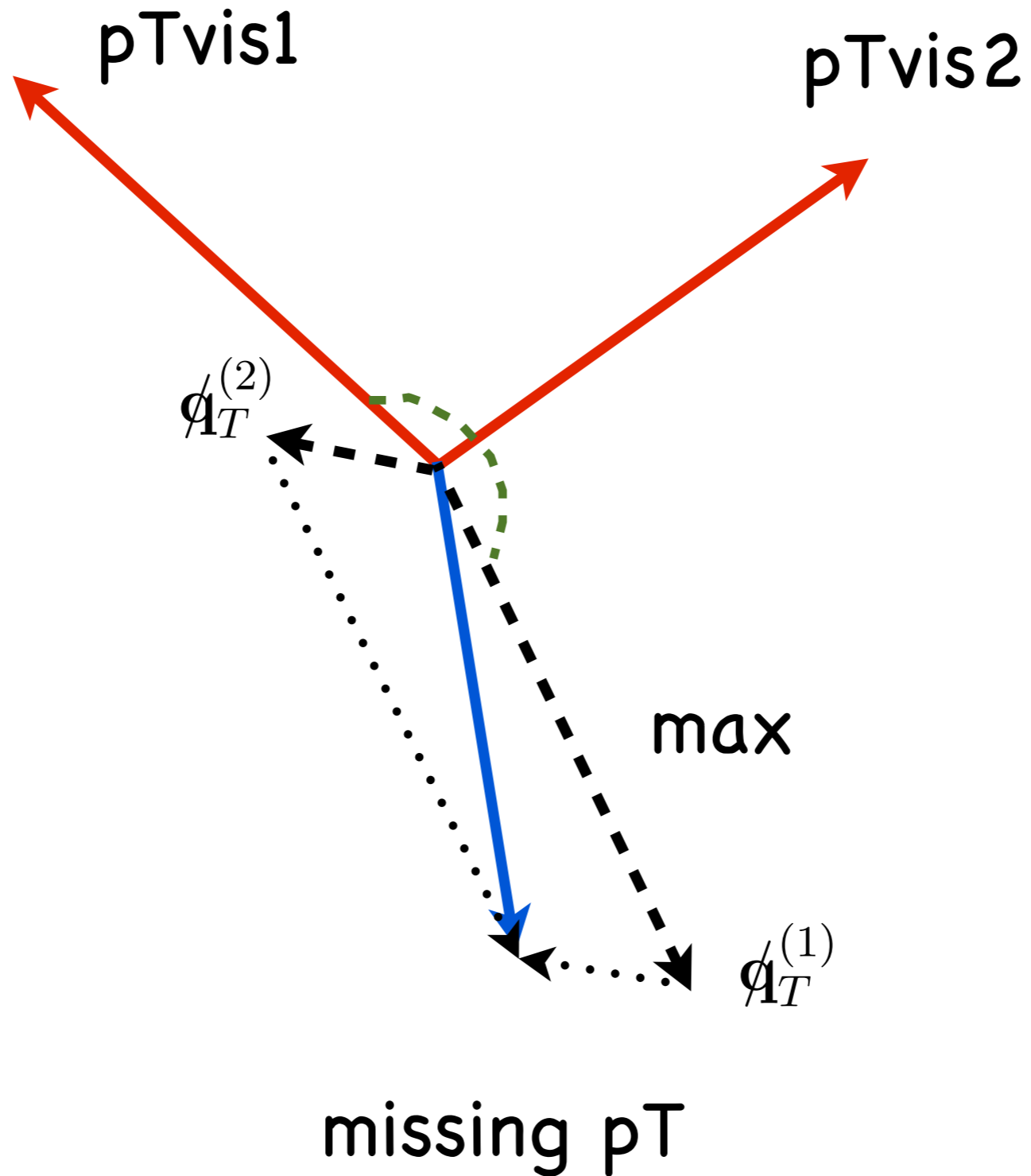
$b\bar{b}\tau^+\tau^-$  :

- Rate improvement
- Though reconstruction of taus difficult
- Plenty and sizable elw. backgrounds, e.g Zbb, HZ, WH, ttbar
- Best strategy:
  - require boost (fatjet) + jet substructure
  - reject ttbar with mT2

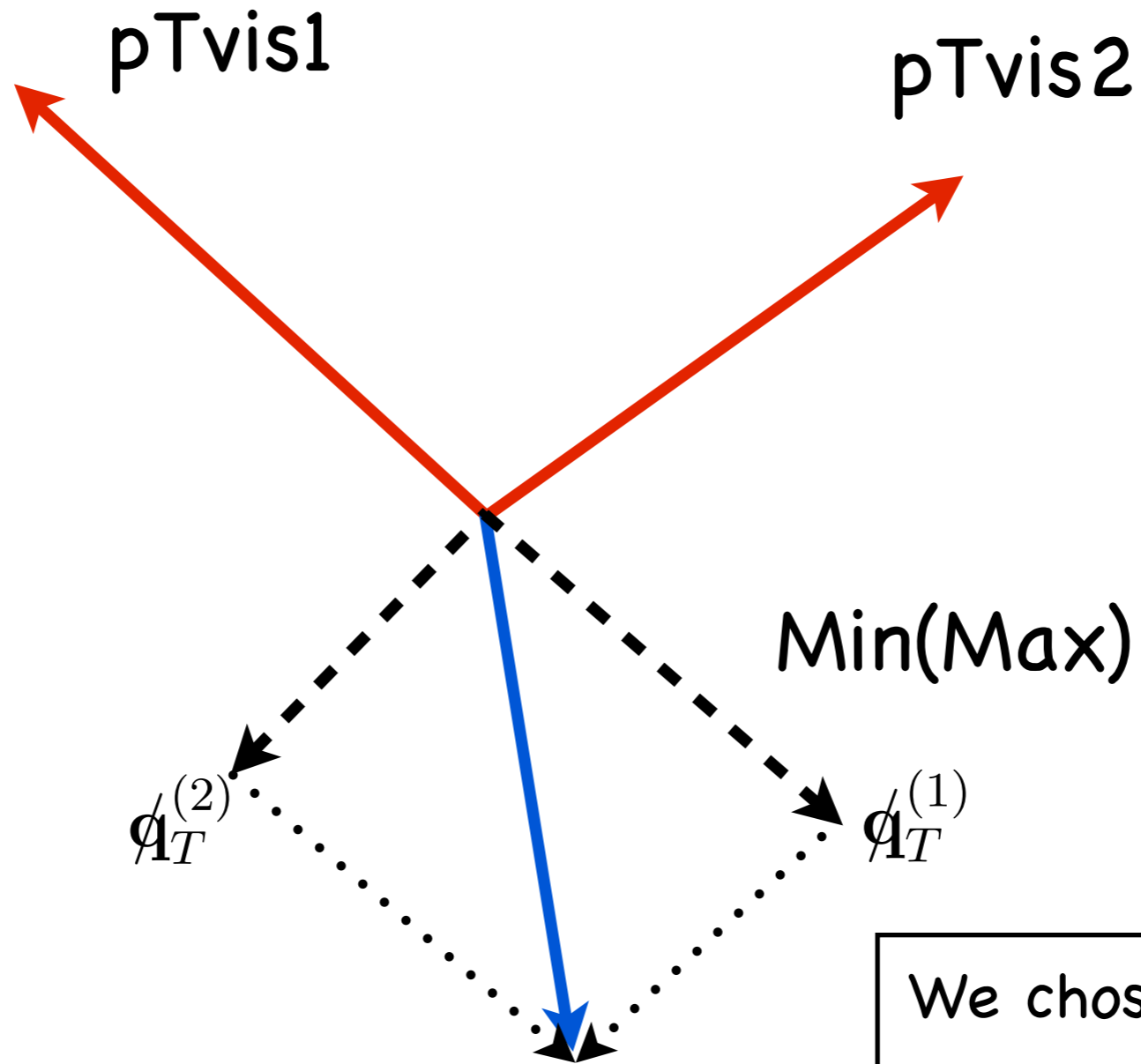




$$m_{T2}^2(\chi) \equiv \min_{\mathbf{q}_T^{(1)} + \mathbf{q}_T^{(2)} = \mathbf{p}_T} \left[ \max \left\{ m_T^2(\mathbf{p}_T^{\pi^{(1)}}; \mathbf{q}_T^{(1)}; \chi), m_T^2(\mathbf{p}_T^{\pi^{(2)}}; \mathbf{q}_T^{(2)}; \chi) \right\} \right]$$



$$m_{T2}^2(\chi) \equiv \min_{\mathbf{q}_T^{(1)} + \mathbf{q}_T^{(2)} = \mathbf{p}_T} \left[ \max \left\{ m_T^2(\mathbf{p}_T^{\pi^{(1)}}, \mathbf{q}_T^{(1)}; \chi), m_T^2(\mathbf{p}_T^{\pi^{(2)}}, \mathbf{q}_T^{(2)}; \chi) \right\} \right]$$



We chose:  
 pTvis2=b1 and pTvis1=b2  
 → vetos ttbar...  
 more prob. HZ, Zbb

$$m_{T2}^2(\chi) \equiv \min_{\mathbf{q}_T^{(1)} + \mathbf{q}_T^{(2)} = \mathbf{p}_T} \left[ \max \left\{ m_T^2(\mathbf{p}_T^{\pi^{(1)}}, \mathbf{q}_T^{(1)}; \chi), m_T^2(\mathbf{p}_T^{\pi^{(2)}}, \mathbf{q}_T^{(2)}; \chi) \right\} \right]$$

# Possible final states:

$b\bar{b}\tau^+\tau^-$  :

- Rate improvement
- Though reconstruction of taus difficult
- Plenty and sizable elw. backgrounds, e.g Zbb, HZ, WH, ttbar
- Best strategy: - require boost (fatjet) + jet substructure  
- reject ttbar with mT2

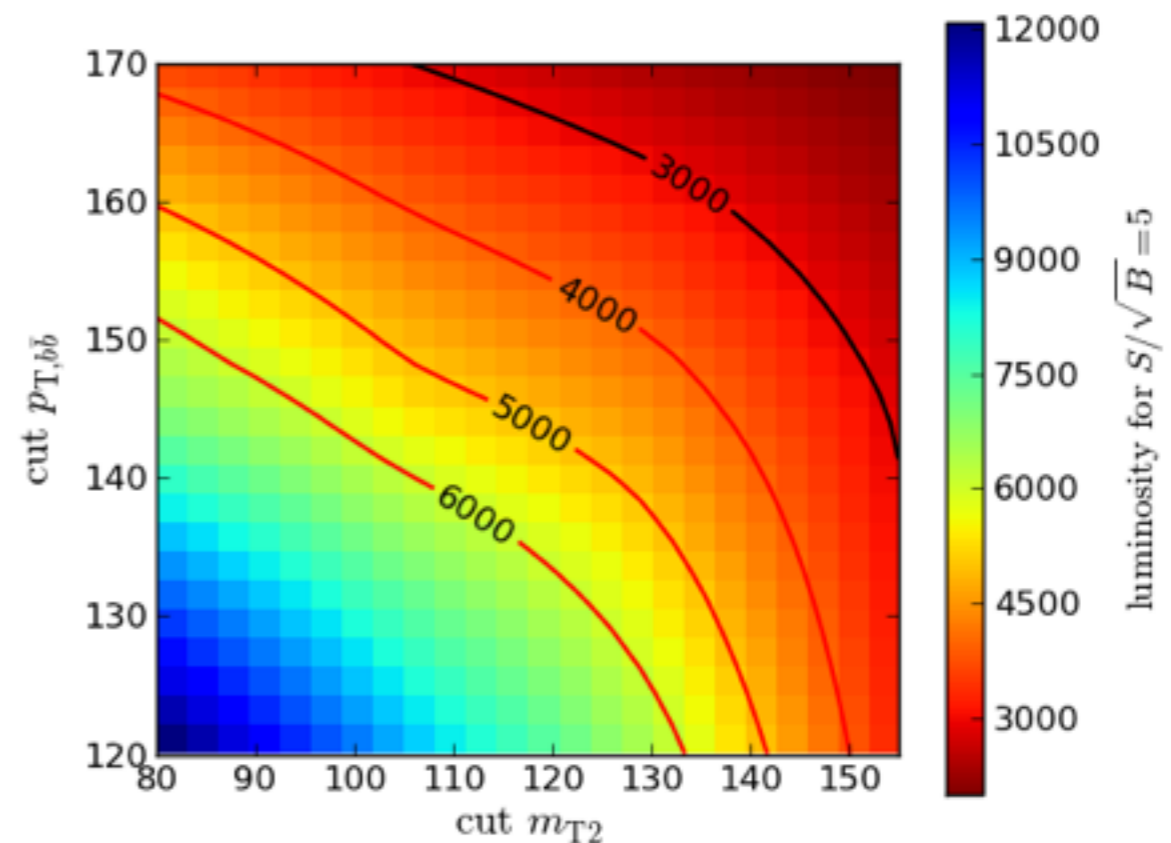
'straightforward'  
(without jet substructure)  
to obtain S/B ~ 1/5

Exclusion at 95% CL:

$$\lambda > \lambda_{95\% \text{ CL}}^{3000/\text{fb}} \simeq 3.0 \times \lambda_{\text{SM}}$$

[Dolan, Englert, MS]

[Barr, Dolan, Englert, MS]



# Possible final states:

[Papaefstathiou, Yang, Zurita]

$$\bar{b}bW^+W^- :$$

$$hh \rightarrow \bar{b}bW^+W^- \rightarrow \bar{b}b\ell\nu jj$$

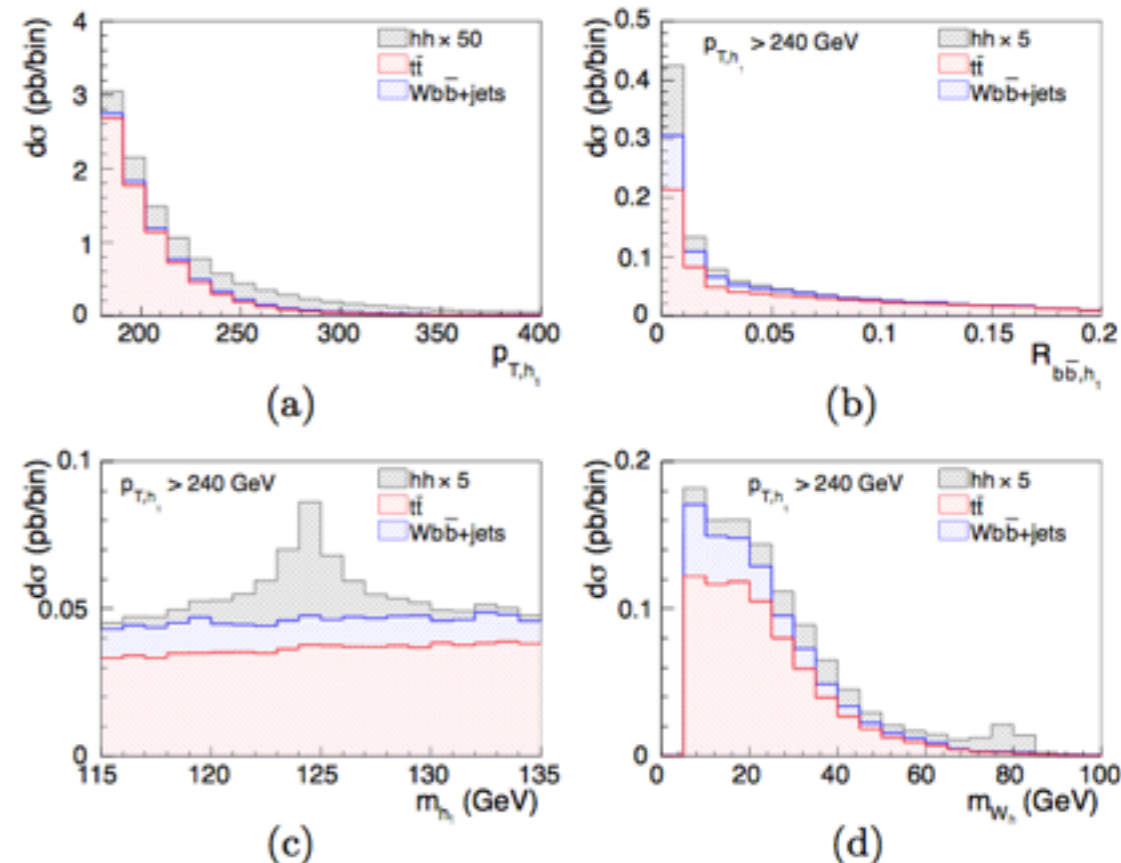
- Fully reconstructable final state
- Triggering easy due to lepton
- But looks like  $t\bar{t}$ ...

Process	$\sigma_{\text{initial}}$ (fb)	$\sigma_{\text{basic}}$ (fb)
$hh \rightarrow \bar{b}b\ell\nu jj$	2.34	0.134
$t\bar{t} \rightarrow \bar{b}b\ell\nu jj$	$240 \times 10^3$	15.5
$W(\rightarrow \ell\nu)\bar{b}b+\text{jets}$	$2.17 \times 10^3$	0.97
$W(\rightarrow \ell\nu)+\text{jets}$	$2.636 \times 10^6$	$\mathcal{O}(0.01)$
$h(\rightarrow \ell\nu jj)+\text{jets}$	36.11	$\mathcal{O}(0.0001)$
$h(\rightarrow \ell\nu jj)\bar{b}b$	6.22	$\mathcal{O}(0.001)$
$h(\rightarrow \bar{b}b) + WW(\rightarrow \ell\nu jj)$	0.0252	-

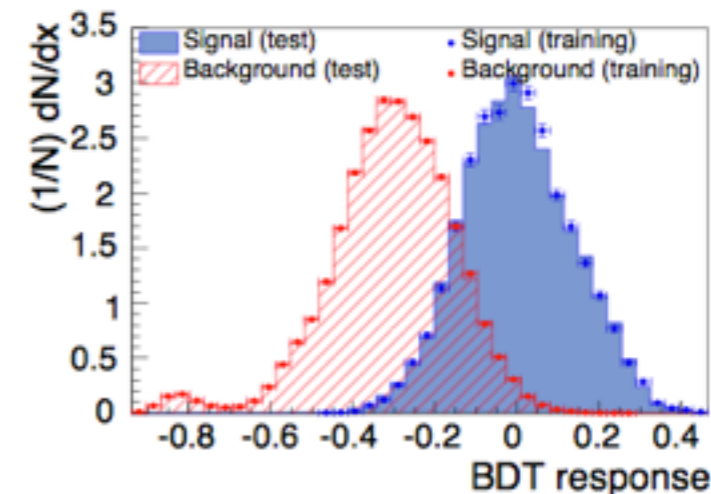
For SM coupling

$$S/\sqrt{S+B} \sim 2.4 \quad 3.1\sigma$$

with  $S=9$  and  $B=6$  after 600 fb



BDT



# Possible final states:

[Baur, Plehn, Rainwater]

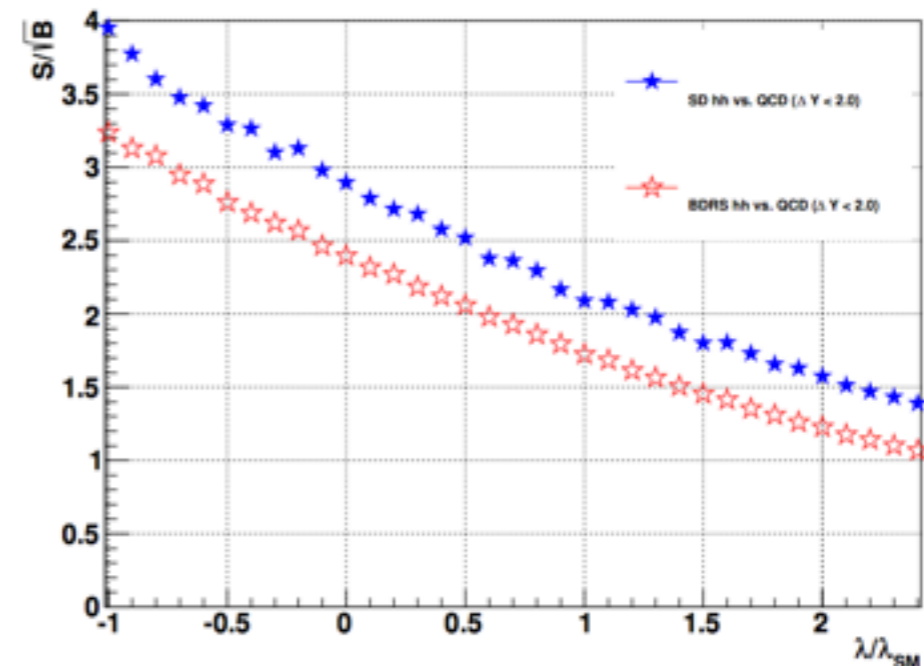
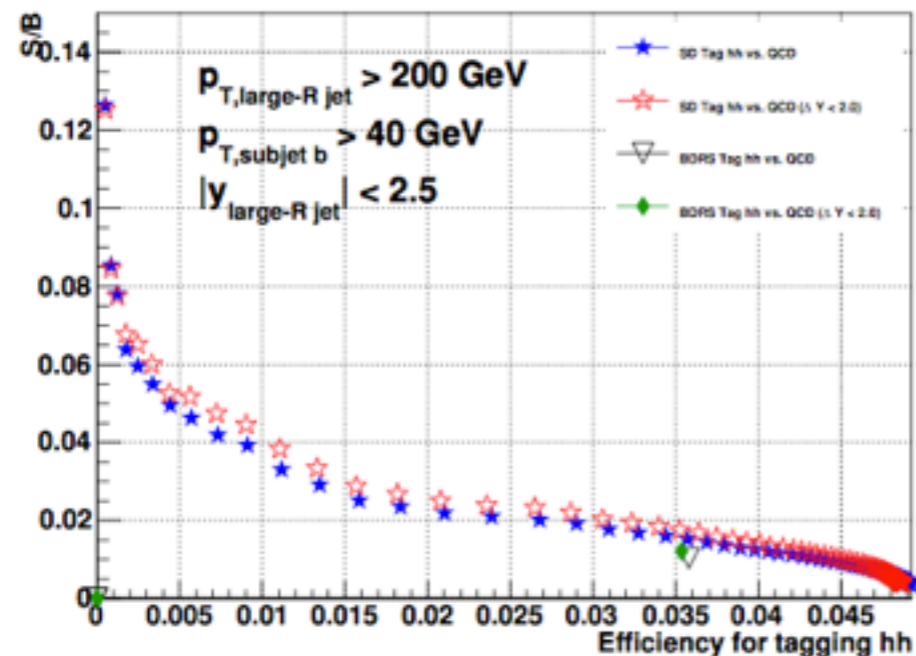
[Dolan, Englert, MS]

[Papaefstathiou, Ferreira, MS]

$\bar{b}b\bar{b}b$  :

- Difficult to trigger (requires large pT cuts or fat jet)
- Huge QCD backgrounds
- Can try to use jet substructure techniques to overcome large backgrounds
- Maybe sideband possible?
- After reconstruction and 3000 fb:

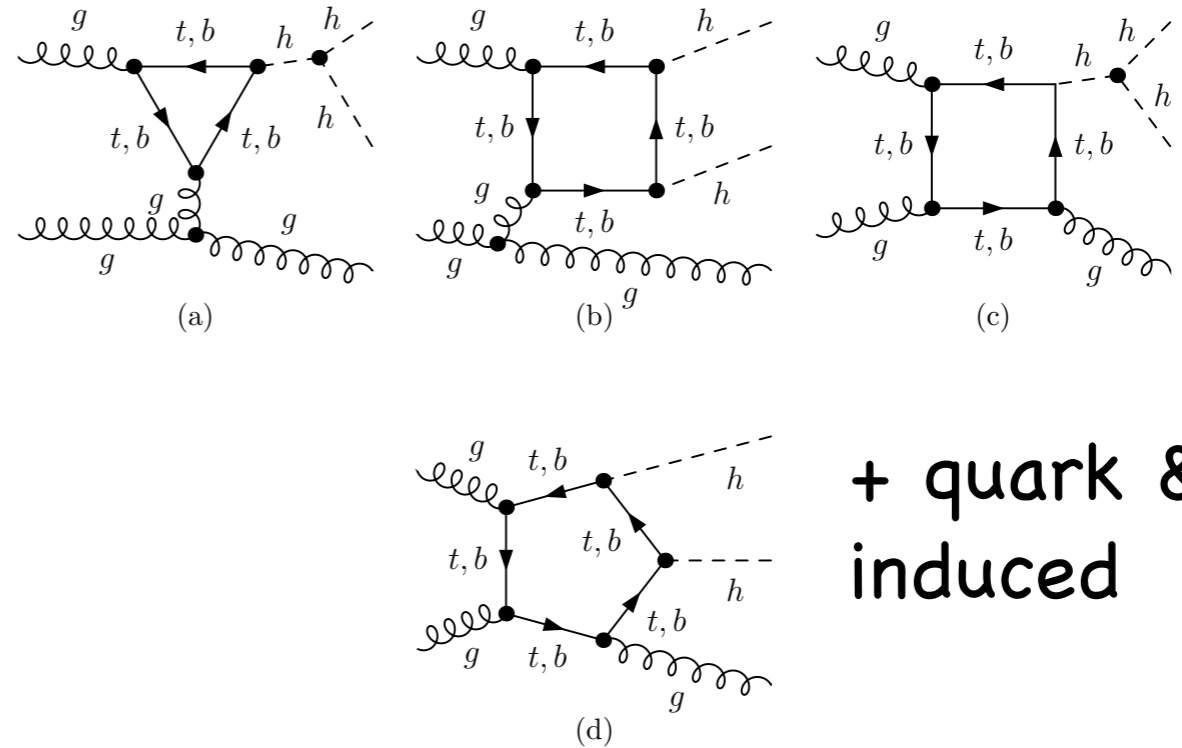
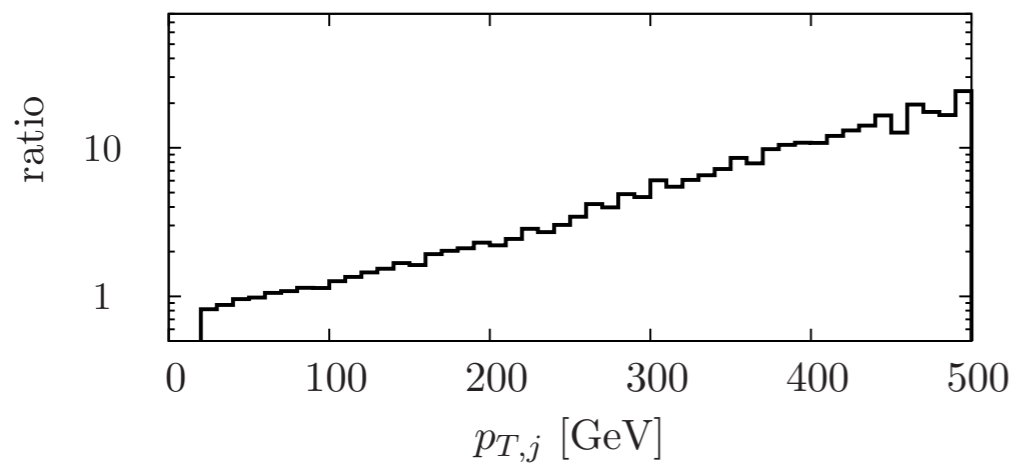
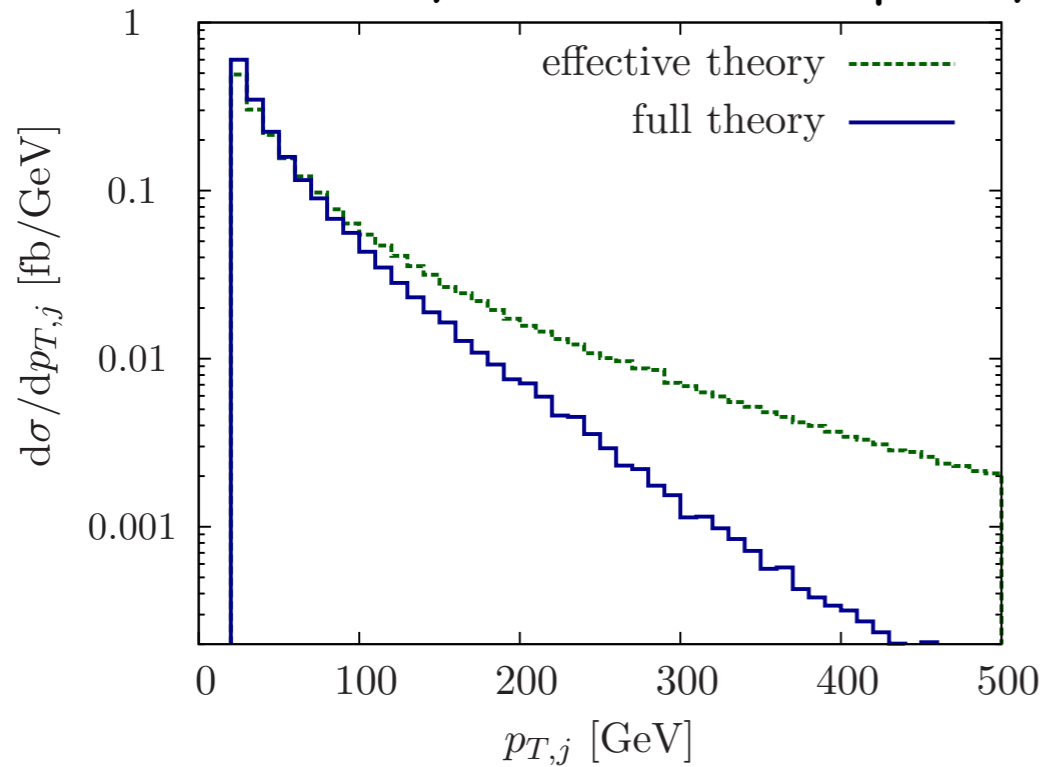
sample	$\sigma_{\text{initial}}$ (fb)
$hh, h \rightarrow b\bar{b}$ (SM)	10.7
QCD ( $b\bar{b}$ )( $b\bar{b}$ )	$151.1 \times 10^3$
$Zb\bar{b}, Z \rightarrow b\bar{b}$	$8.8 \times 10^3$
$hZ, h \rightarrow b\bar{b}, Z \rightarrow b\bar{b}$	70.0
$hW, h \rightarrow b\bar{b}, W \rightarrow c\bar{b}(\bar{c}b)$	96.4



# More jets more fun

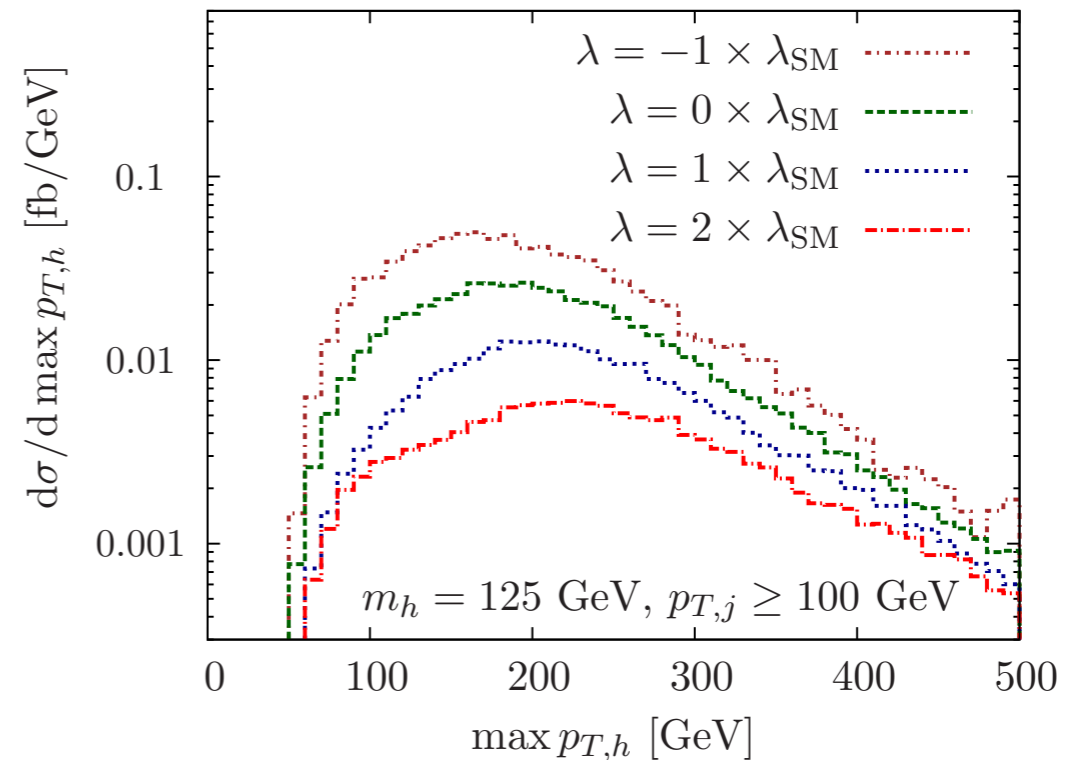
- need to work a little harder

Eff. theory breaks down quickly



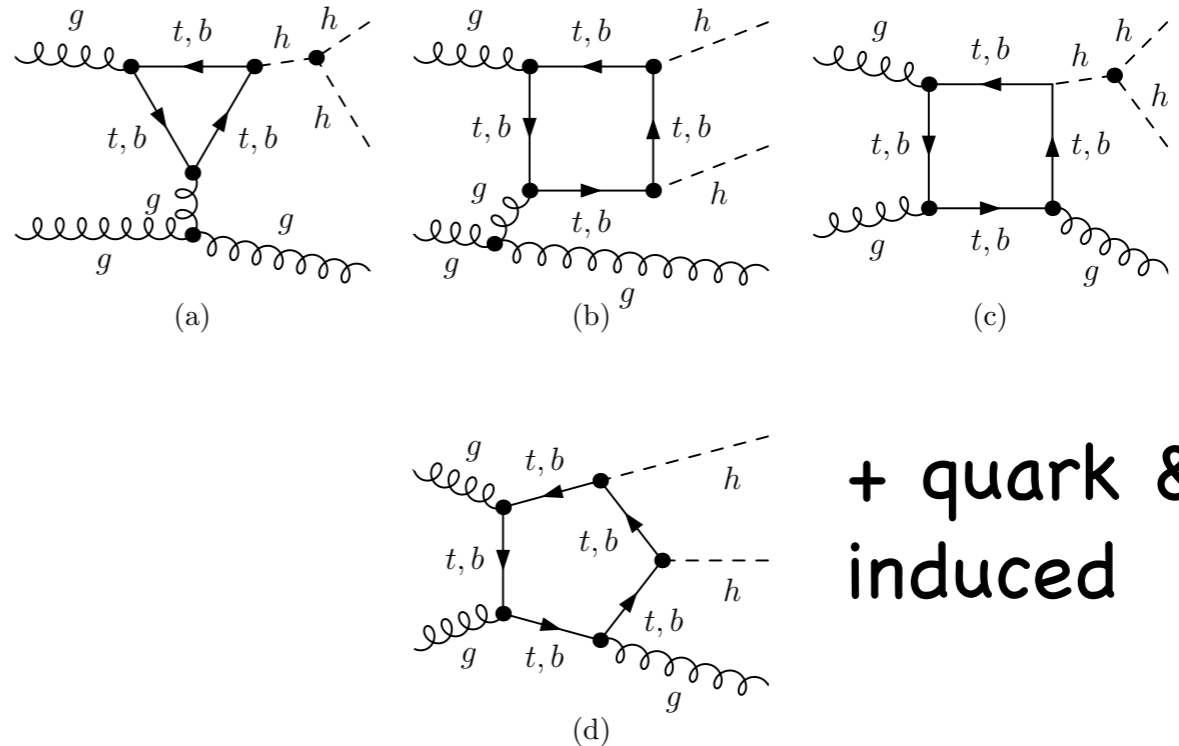
+ quark & gluon induced

retain sensitivity for boosted Higgs

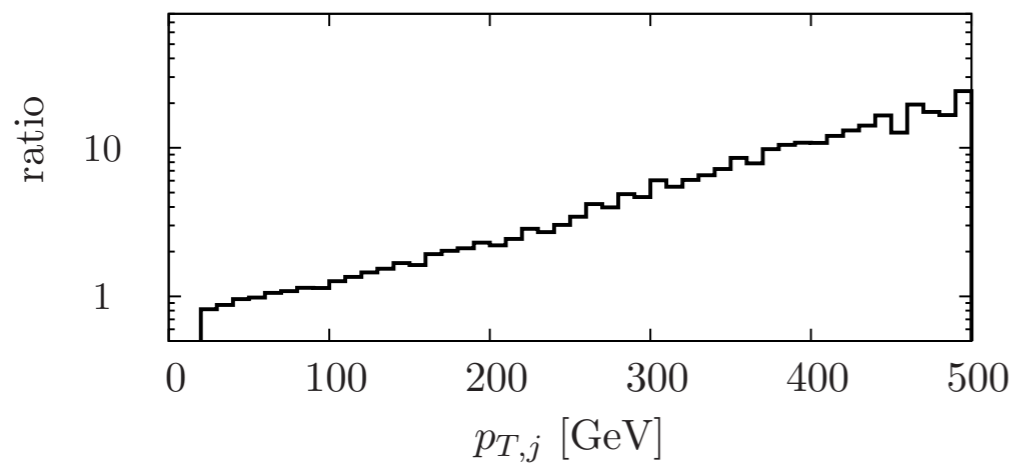
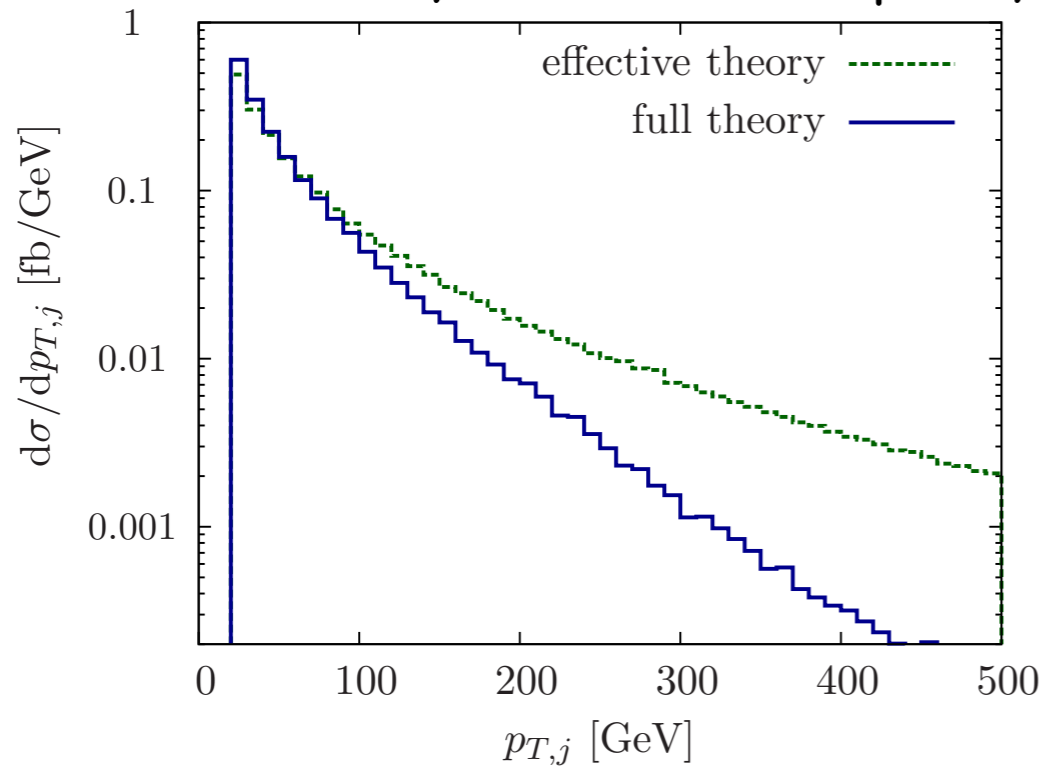


# More jets more fun

- need to work a little harder

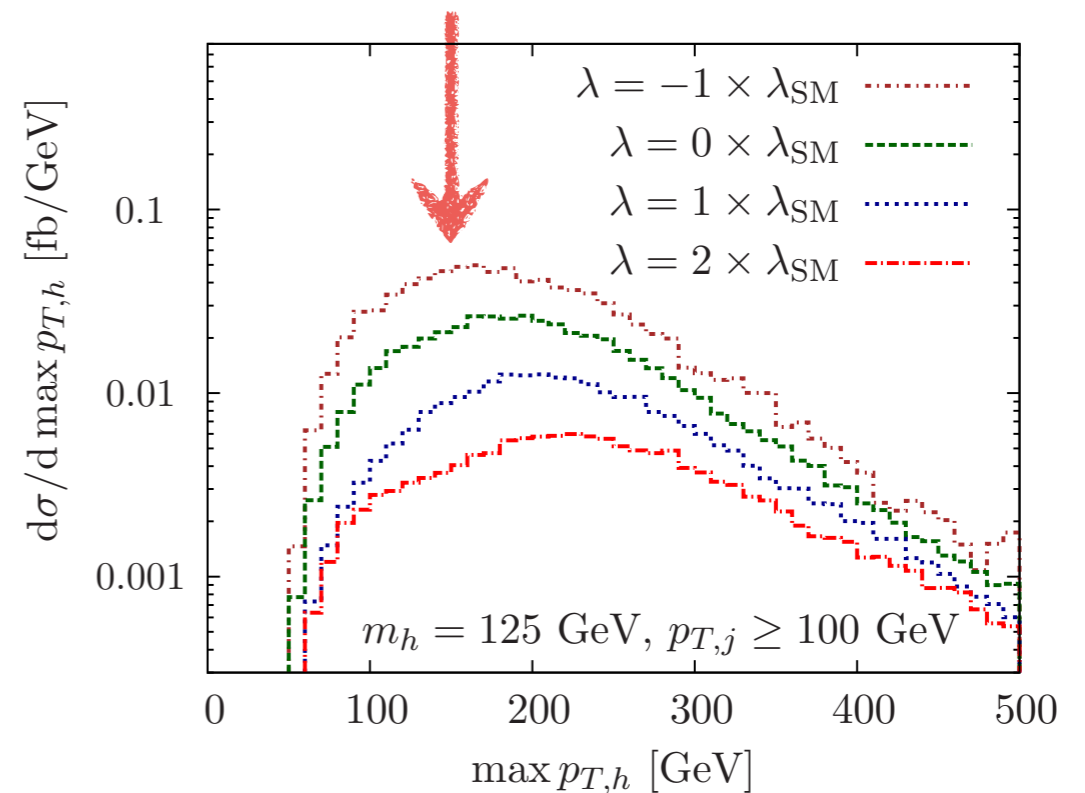


Eff. theory breaks down quickly



+ quark & gluon induced

retain sensitivity for boosted Higgs





# Higgs selfcoupling in HHj+X

- Additional jet ameliorates  $1/m_{hh}^2$  suppression
- Jet adds handle to suppress backgrounds -> improvement S/B
- But cross section very small

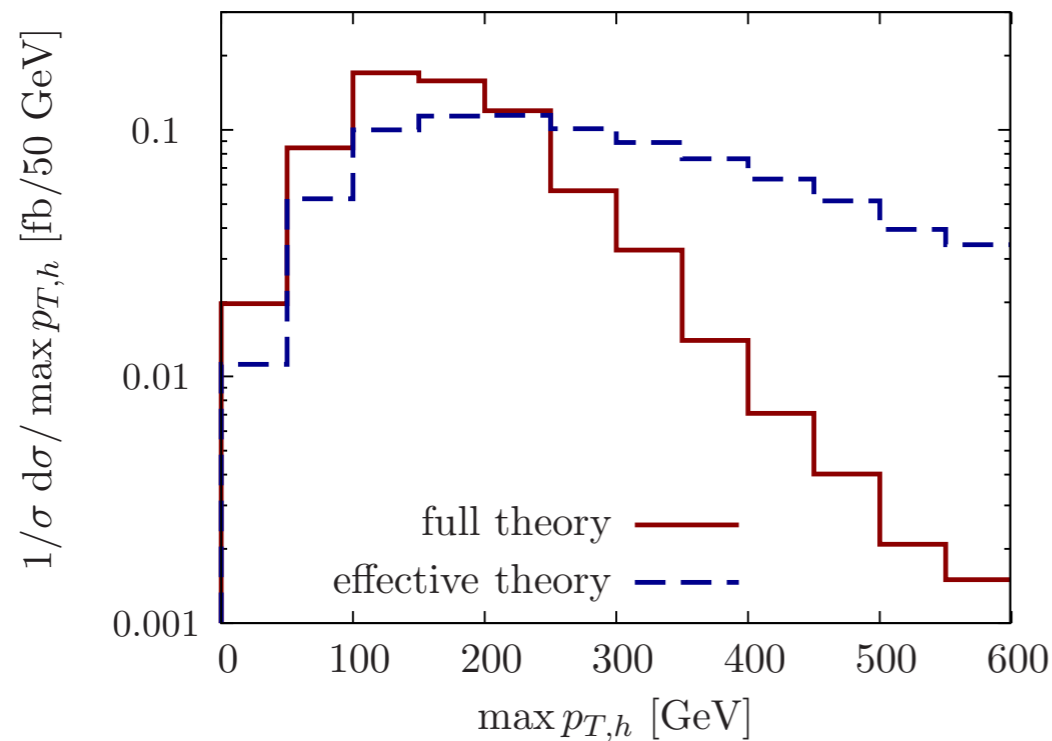
$b\bar{b}\tau^+\tau^-$  (assuming small tau fake rate)

	$\xi = 0$	$\xi = 1$	$\xi = 2$	$b\bar{b}\tau\tau$	$b\bar{b}\tau\tau$ [ELW]	$b\bar{b}W^+W^-$	ratio to $\xi = 1$
cross section before cuts	59.48	28.34	13.36	67.48	8.73	873000	$3.2 \cdot 10^{-5}$
reconstructed Higgs from $\tau$ s	4.05	1.94	0.91	2.51	1.10	1507.99	$1.9 \cdot 10^{-3}$
fatjet cuts	2.27	1.09	0.65	1.29	0.84	223.21	$4.8 \cdot 10^{-3}$
kinematic Higgs reconstruction ( $m_{b\bar{b}}$ )	0.41	0.26	0.15	0.104	0.047	9.50	$2.3 \cdot 10^{-2}$
Higgs with double $b$ -tag	0.148	0.095	0.053	0.028	0.020	0.15	0.48

$b\bar{b}\tau^+\tau^-j$  (assuming small tau fake rate)

	$\xi = 0$	$\xi = 1$	$\xi = 2$	$b\bar{b}\tau^+\tau^-j$	$b\bar{b}\tau^+\tau^-j$ [ELW]	$t\bar{t}j$	ratio to $\xi = 1$
cross section before cuts	6.45	3.24	1.81	66.0	1.67	106.7	$1.9 \cdot 10^{-2}$
2 $\tau$ s	0.44	0.22	0.12	37.0	0.94	7.44	$4.8 \cdot 10^{-3}$
Higgs rec. from taus + fatjet cuts	0.29	0.16	0.10	2.00	0.150	0.947	$5.1 \cdot 10^{-2}$
kinematic Higgs rec.	0.07	0.04	0.02	0.042	0.018	0.093	0.26
2b + $hh$ invariant mass + $p_{T,j}$ cut	0.010	0.006	0.004	<0.0001	0.0022	0.0014	1.54

# Higgs selfcoupling in HHjj+X



- Test for long. gauge boson scattering
- For kinematic distributions full loop recommended
- Gluon fusion dominating over WBF
- Analysis in  $\bar{b}b\tau^+\tau^-$

So far very rudimentary analysis:

	Signal with $\xi \times \lambda$			Background		$S/B$
	$\xi = 0$	$\xi = 1$	$\xi = 2$	$t\bar{t}jj$	Other BG	ratio to $\xi = 1$
tau selection cuts	0.212	0.091	0.100	3101.0	57.06	$0.026 \times 10^{-3}$
Higgs rec. from taus	0.212	0.091	0.100	683.5	31.92	$0.115 \times 10^{-3}$
Higgs rec. from $b$ jets	0.041	0.016	0.017	7.444	0.303	$1.82 \times 10^{-3}$
2 tag jets	0.024	0.010	0.012	5.284	0.236	$1.65 \times 10^{-3}$
incl. GF after cuts/re-weighting	0.181	0.099	0.067	5.284	0.236	1/61.76

Very bad S/B, but can be improved a lot...

# New Physics in HH

## Resonant enhancement

see also [Chen et al. 1312.7212]

- SUSY,  $H \rightarrow hh$

Measurement of rel. CS  $Hhh$  and  $hhh$  translates directly to measurement of  $\alpha$  and  $\beta$

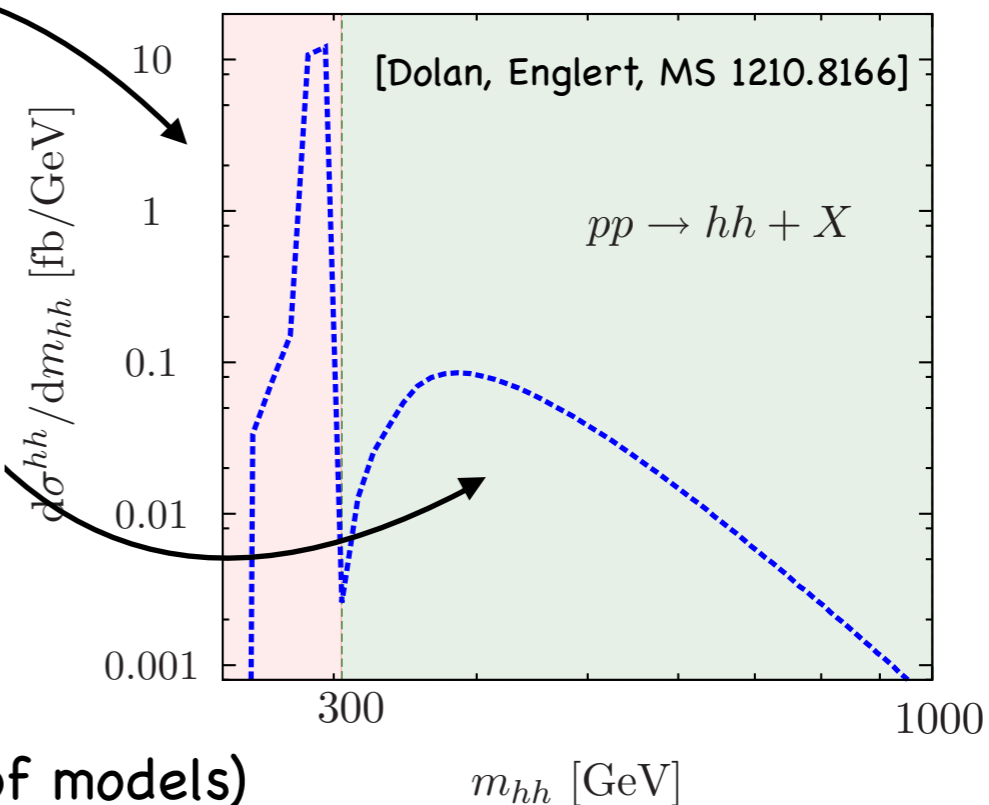
$$\lambda_{hhh} = 3 \cos 2\alpha \sin(\beta + \alpha)$$

$$\lambda_{Hhh} = 2 \sin 2\alpha \sin(\beta + \alpha) - \cos 2\alpha \cos(\beta + \alpha)$$

$$\lambda_{HHh} = -2 \sin 2\alpha \cos(\beta + \alpha) - \cos 2\alpha \sin(\beta + \alpha)$$

(same in Higgs portal type of models)

Assuming decoupling limit such that  $M_H > 2 M_h$  and  $BR(H \rightarrow hh) = 45\%$

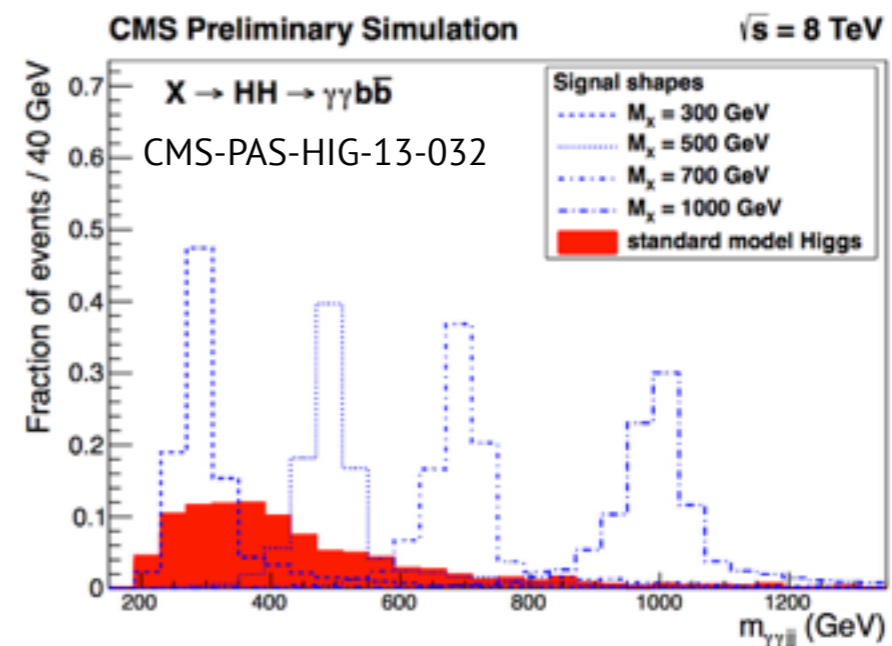


- E-dim,  $G \rightarrow hh \rightarrow 4b$

see [Gouzevitch et al. 1303.6636]

- Higgs portal

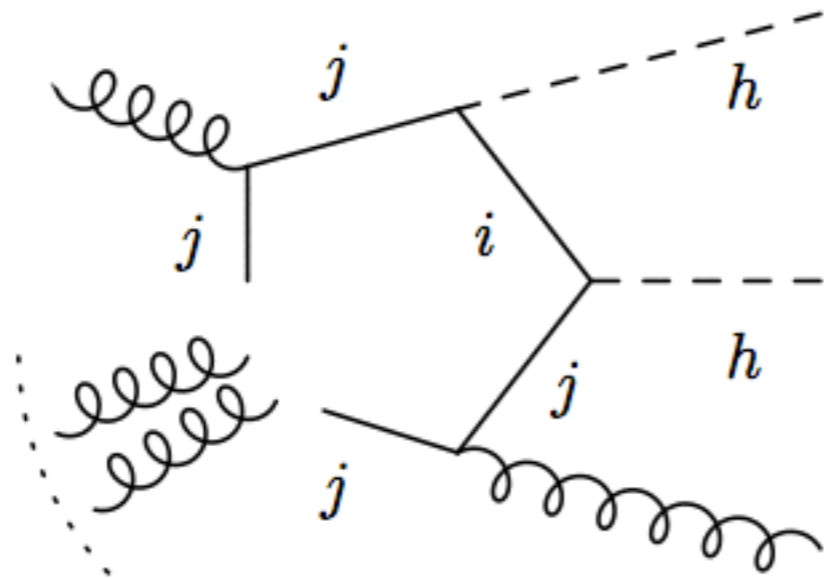
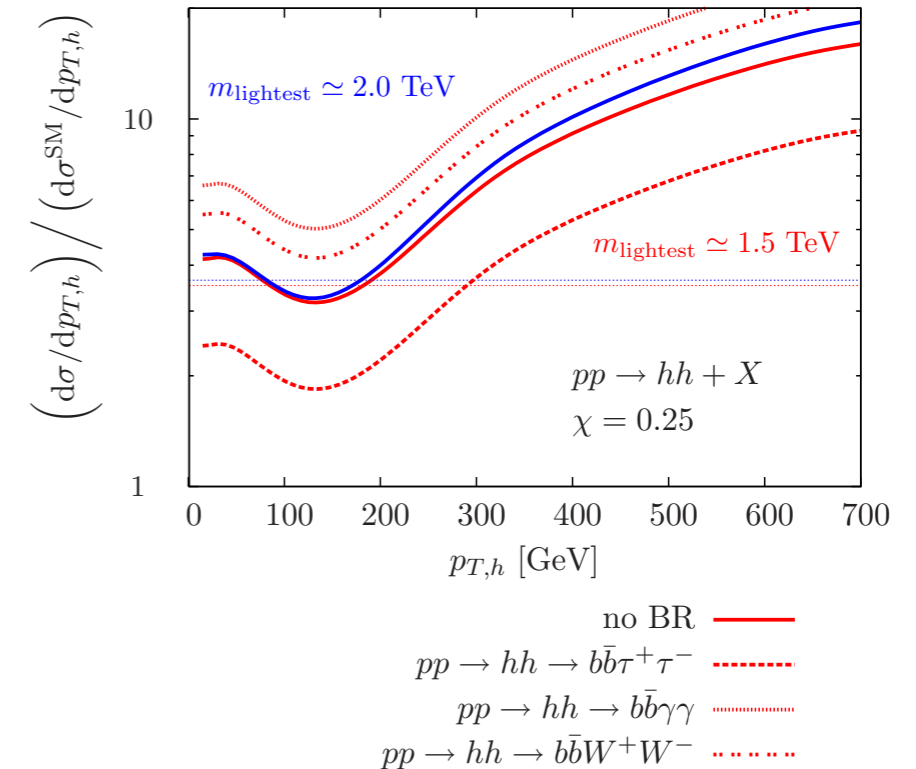
see [No, Ramsey-Musolf 1310.6035]



# New Physics in HH

## Continuous/Loop enhancement

- Composite Higgs
- 4th generation  
see [Kribs, Plehn, Tait, MS 0706.3718]
- Other theories modifying  $hh\bar{t}_i t_j$  or  $h\bar{t}_i t_j$



Usually high- $p_T$  region  
enhanced over SM

# A portal to the Higgs width?

New idea how to constrain Higgs width at hadron colliders:

[Caola, Melnikov PRD 88]

## Constraining the Higgs boson width with $ZZ$ production at the LHC

Fabrizio Caola<sup>1,\*</sup> and Kirill Melnikov<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of  $pp \rightarrow ZZ$  cross-section at the LHC in a broad range of  $ZZ$  invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of  $ZZ$  invariant masses from 100 to 800 GeV, we find  $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$  MeV, at the 95% confidence level. Restricting  $ZZ$  invariant masses to  $M_{ZZ} \geq 300$  GeV range, we estimate that this bound can be improved to  $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$  MeV.

# A portal to the Higgs width?

New idea how to constrain Higgs width at hadron colliders:

[Caola, Melnikov PRD 88]

## Constraining the Higgs boson width with $ZZ$ production at the LHC

Fabrizio Caola<sup>1,\*</sup> and Kirill Melnikov<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of  $pp \rightarrow ZZ$  cross-section at the LHC in a broad range of  $ZZ$  invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of  $ZZ$  invariant masses from 100 to 800 GeV, we find  $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$  MeV, at the 95% confidence level. Restricting  $ZZ$  invariant masses to  $M_{ZZ} \geq 300$  GeV range, we estimate that this bound can be improved to  $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$  MeV.



Measurement done in CMS-PAS-HIG-14-002 and presented at Moriond '14  
ATLAS is working to perform same measurement

# A portal to the Higgs width?

New idea how to constrain Higgs width at hadron colliders:

[Caola, Melnikov PRD 88]

## Constraining the Higgs boson width with $ZZ$ production at the LHC

Fabrizio Caola<sup>1,\*</sup> and Kirill Melnikov<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of  $pp \rightarrow ZZ$  cross-section at the LHC in a broad range of  $ZZ$  invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of  $ZZ$  invariant masses from 100 to 800 GeV, we find  $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$  MeV, at the 95% confidence level. Restricting  $ZZ$  invariant masses to  $M_{ZZ} \geq 300$  GeV range, we estimate that this bound can be improved to  $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$  MeV.



Measurement done in CMS-PAS-HIG-14-002 and presented at Moriond '14  
ATLAS is working to perform same measurement



**QUANTUM DIARIES**

Tevatron and LHC experiments so far. The week went on to include a spectacular CMS [result on the Higgs width](#).

# A portal to the Higgs width?

New idea how to constrain Higgs width at hadron colliders:

[Caola, Melnikov PRD 88]

## Constraining the Higgs boson width with $ZZ$ production at the LHC

Fabrizio Caola<sup>1,\*</sup> and Kirill Melnikov<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of  $pp \rightarrow ZZ$  cross-section at the LHC in a broad range of  $ZZ$  invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of  $ZZ$  invariant masses from 100 to 800 GeV, we find  $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$  MeV, at the 95% confidence level. Restricting  $ZZ$  invariant masses to  $M_{ZZ} \geq 300$  GeV range, we estimate that this bound can be improved to  $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$  MeV.



Measurement done in CMS-PAS-HIG-14-002 and presented at Moriond '14

ATLAS is working to perform same measurement



**ATLAS Experiment Blog**  
Higgs boson width.) This new measurement shows a remarkable sensitivity and constrains the Higgs boson width to be below 17 MeV, more than two orders of magnitude better than the previous limits! Standard Model 2.5 :  
QUA  
Tevatron and  
width.



# A portal to the Higgs width?

New idea how to constrain Higgs width at hadron colliders:

[Caola, Melnikov PRD 88]

## Constraining the Higgs boson width with $ZZ$ production at the LHC

Fabrizio Caola<sup>1,\*</sup> and Kirill Melnikov<sup>1,†</sup>

<sup>1</sup>*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of  $pp \rightarrow ZZ$  cross-section at the LHC in a broad range of  $ZZ$  invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of  $ZZ$  invariant masses from 100 to 800 GeV, we find  $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$  MeV, at the 95% confidence level. Restricting  $ZZ$  invariant masses to  $M_{ZZ} \geq 300$  GeV range, we estimate that this bound can be improved to  $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$  MeV.



Measurement done in CMS-PAS-HIG-14-002 and presented at Moriond '14  
ATLAS is working to perform same measurement



ATLAS

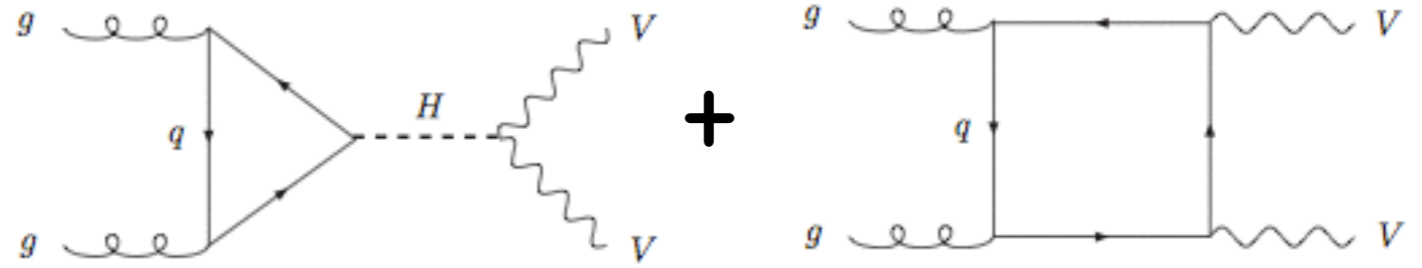
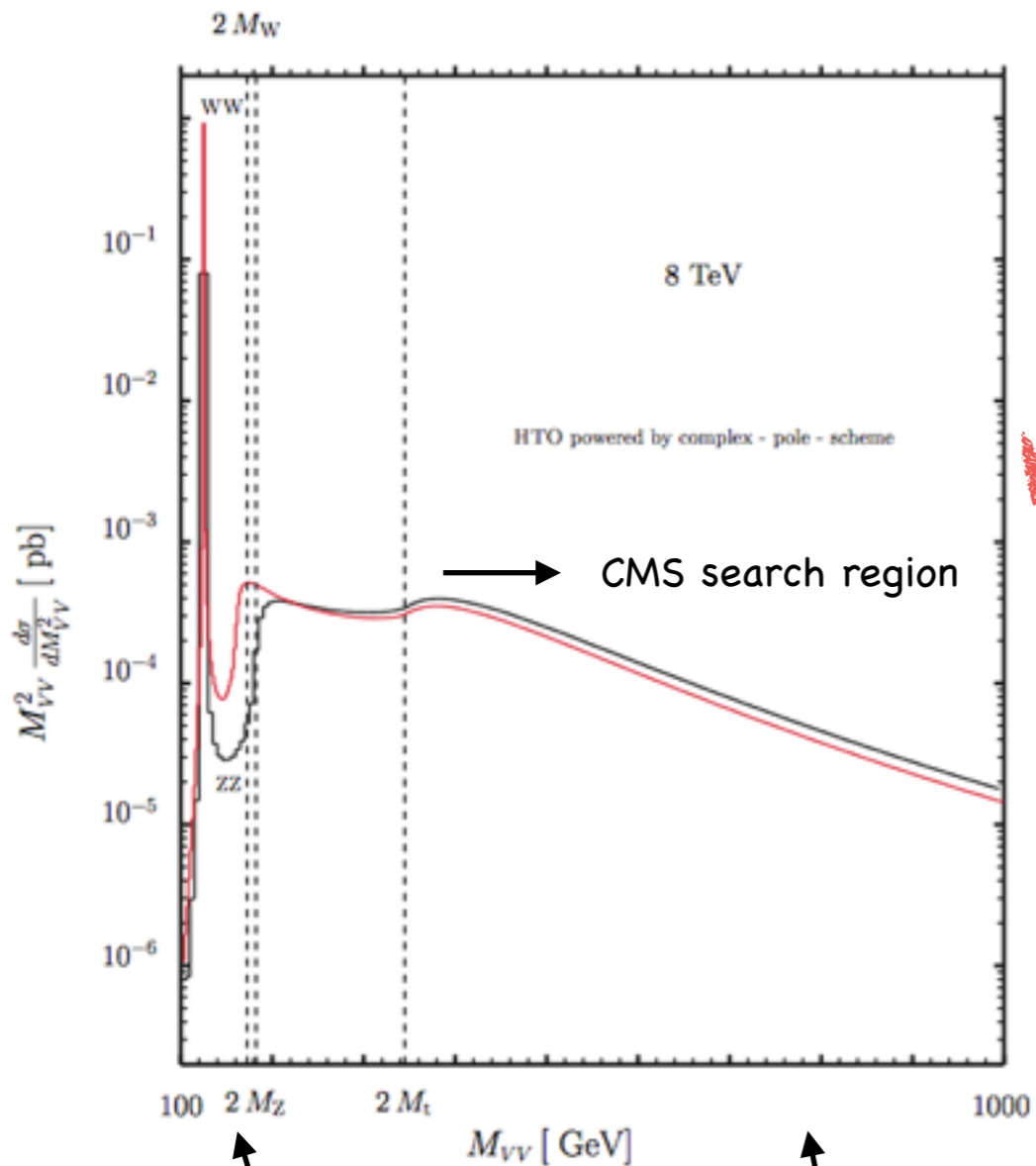
**LIFE AND PHYSICS**  
**JON BUTTERWORTH**  
HOSTED BY THE GUARDIAN

[WIKI](#)

### How wide is a Higgs?

In accord with Heisenberg's uncertainty principle, short-lived particles have uncertain mass. So the Higgs boson, which gives mass to other particles, is uncertain about its own mass. New results from the CMS experiment at the CERN LHC have started to tell us how uncertain

# CMS Measurement



- measure  $g_{ggH}^2 g_{HZZ}^2$  off-peak using angular correlations of 4l decay products
- use in on-peak relation to measure width after fixing signal strength

$$\mu_{i,j} = \sigma_{h,i} \times \text{BR}_j \sim \frac{\Gamma_i \Gamma_j}{\Gamma_h}$$

$$\sigma_{\text{on-peak}}^{gg \rightarrow H \rightarrow ZZ} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{\Gamma_H}$$

$$\sigma_{\text{off-peak}}^{gg \rightarrow H \rightarrow ZZ} \sim g_{ggH}^2 g_{HZZ}^2$$

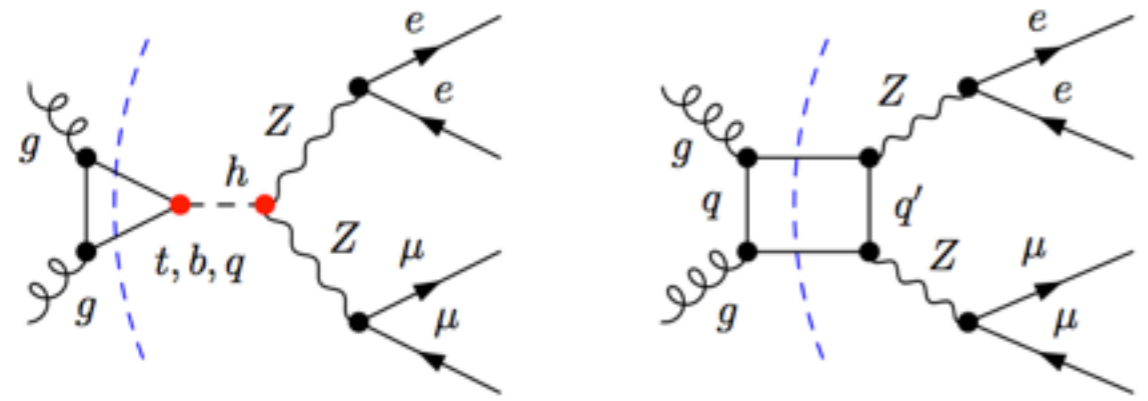
Obs.(exp.) @95% C.L:

$$\Gamma_H < 4.2 (8.5) \Gamma_H^{\text{SM}}$$

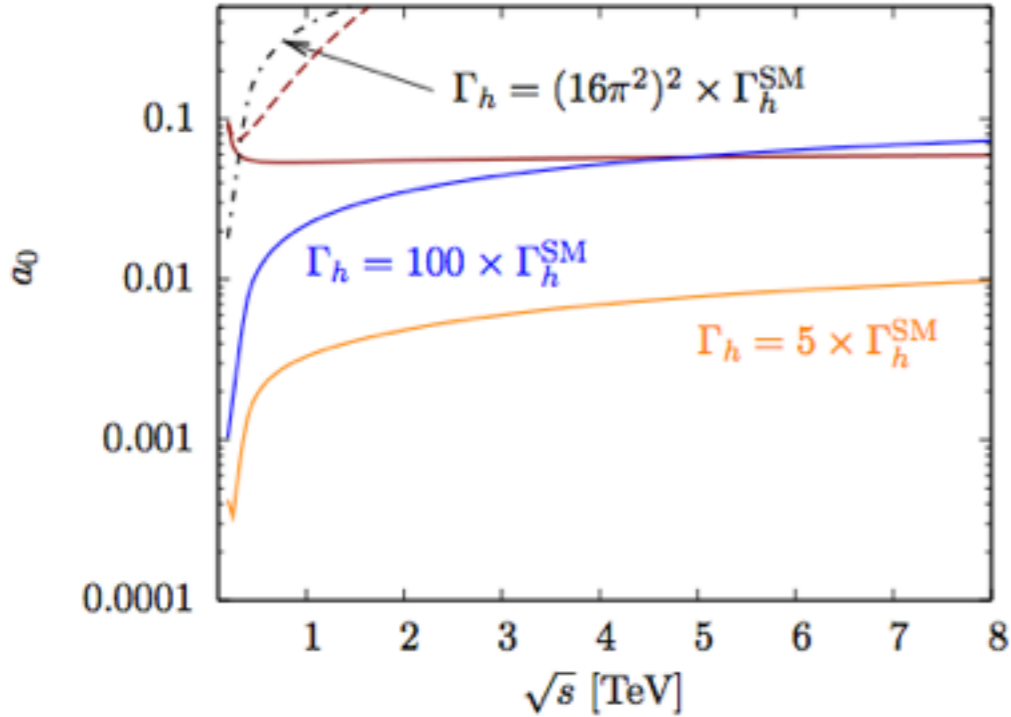
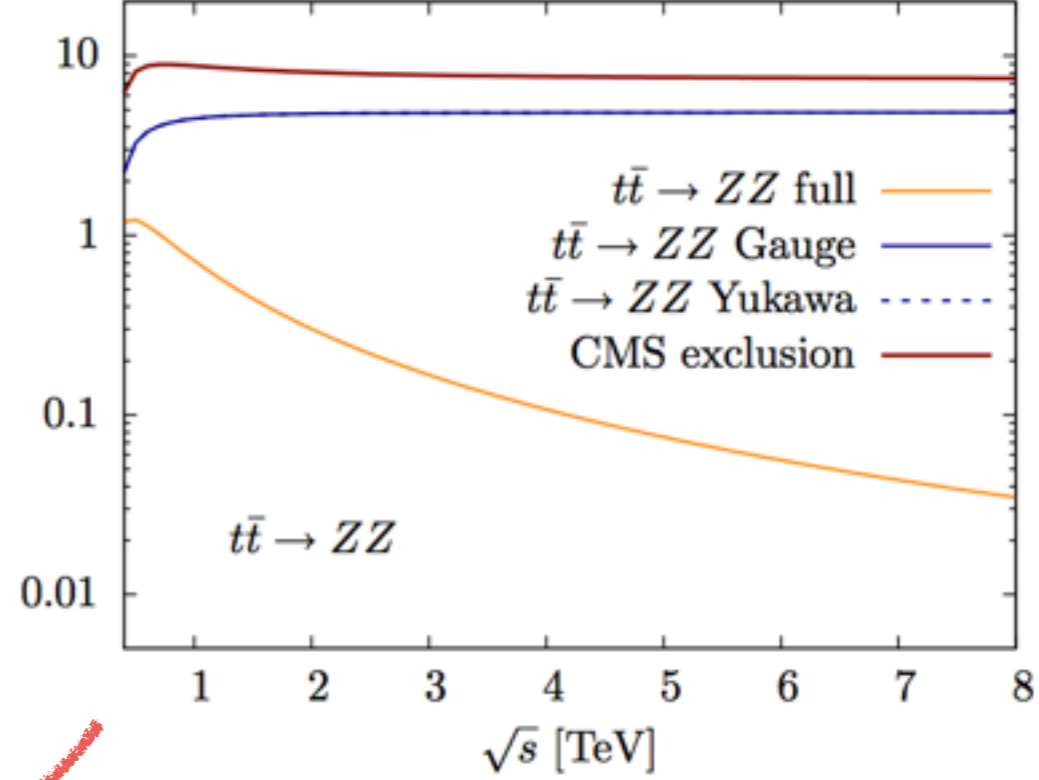
$$\Gamma_H < 17.4 (35.3) \text{ MeV}$$



1. Is measurement well defined?



modifying couplings violates unitarity of fermion-gauge interactions



while CMS test-hypotheses certainly ill-defined not killer at accessible energies

## 2. But how model independent is constraint? [Englert, MS] (submitted)

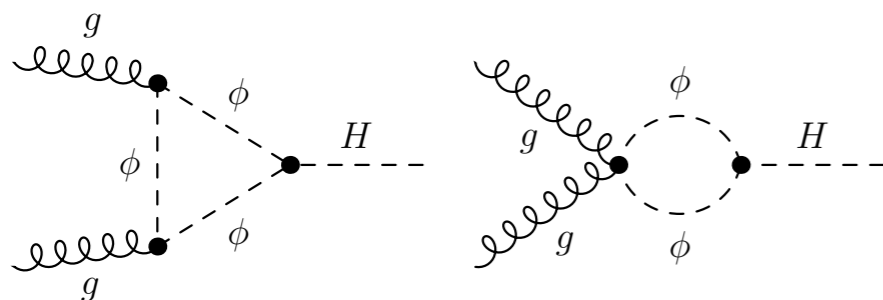
We have seen SM on-shell and off-shell region intimately related by unitarity requirements for fermion-gauge interactions

Direct correlation of on-shell  $g_{ggh}^2 g_{hZZ}^2$  and off-shell  $g_{ggh}^2(\sqrt{s}) g_{hZZ}^2(\sqrt{s})$  necessary ingredient for width measurement  $\rightarrow$  **can be broken by BSM effects**

Consider Higgs-portal (toy) model:  $\mathcal{L}_\phi = |D_\mu \phi|^2 - \tilde{m}_\phi^2 |\phi|^2 - \lambda |\phi|^2 |H|^2 + \dots$

where  $\phi$  scalar only charged under  $SU(3)_C$

$$m_\phi^2 = \tilde{m}_\phi^2 + \lambda v^2 \quad \text{free parameter}$$



$\swarrow$  off-shell CS

$$g_{ggh}(m_h) > g_{ggh,SM} \rightarrow \Gamma > \Gamma_{SM} \quad \text{for } \mu \sim 1$$

Despite increased on-shell coupling (and Higgs width) negligible contribution in off-shell region

Note, shown here only simplest toy model

$m_\phi$	$\mu$ ( $h$ peak)	$\Gamma_h/\Gamma_h^{SM}$	$\bar{\sigma}/\bar{\sigma}^{SM}$ [ $m(4\ell) \geq 330$ GeV] <sup>a</sup>
70 GeV	$\simeq 1.0$	$\simeq 5$	-2%
170 GeV	$\simeq 1.0$	$\simeq 4.7$	+80%
170 GeV	$\simeq 1.0$	$\simeq 1.7$	+6%

<sup>a</sup>We impose the cut set used by CMS [17] without the MELA cut [34].

# Summary

- Increasing interest and wide efforts in HH final state
- Might be possible to measure coupling at LHC 14 TeV  
→ a combination of all accessible final states will be necessary
- Still final states and reconstructions difficult to simulate:  
(ir)reducible backgrounds, taus, bs, photons

Experimentalists → Need finally solid analyses

- Very few HH+(jets) samples: full loop, BSM

Theorists → Need flexible MC

- .....
- Higgs width measurement extremely nice idea but not quite model independent (task for ILC)
  - But off-shell measurement still important:
    - New Physics constraints
    - CP Higgs
    - ....