

Multi-Reference In-medium Similarity Renormalization Group for the Nuclear Matrix Elements of Neutrinoless Double Beta Decay

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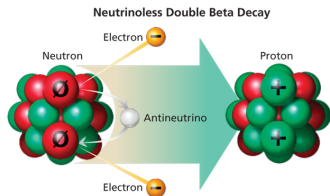
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DOE topical collaboration “Nuclear Theory for Double-Beta Decay and
Fundamental Symmetries”, UMass Amherst, Feb.3, 2017

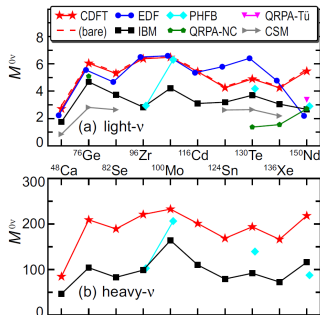
Nuclear Matrix Elements of $0\nu\beta\beta$ decay



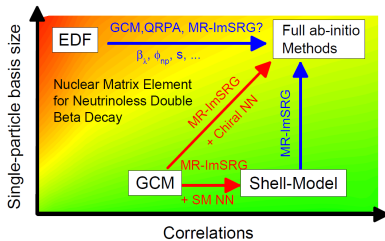
$$M_i^{0\nu} = \langle 0_F^+ | \hat{O}_i^{0\nu} | 0_I^+ \rangle$$

- **Decay mechanism:** transition operator (limited to 1B current mostly, L/H ν)
- **Nuclear structure:** wave functions of initial and final nuclei
Model-dependence (factor of 2-3)

Engel & Menendez, arXiv:1610.06548v1 [nucl-th]



Song, JMY, Ring & Meng, PRC(2017)



ab initio approaches for heavy deformed nuclei

- Rapid developments in *ab initio* approaches (CC, SCGF, IMSRG) for nuclei
- Limitation: spherical or light nuclei
- Multi-Reference IMSRG: a promising approach for heavy deformed nuclei

IMSRG: suppression H_{odd} continuously

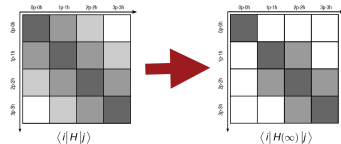
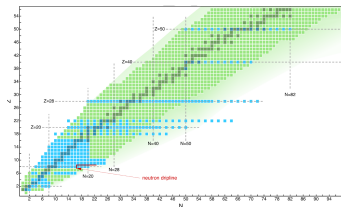
Unitary transformation $H(s) = U(s)H_0U^\dagger(s)$

$$\frac{dH(s)}{ds} = [\eta(s), H(s)],$$

$$\frac{dU(s)}{ds} = \eta(s)U(s)$$

Formal solution:

$$U(s) = \mathcal{S} \exp\left[\int_0^s ds' \eta(s')\right]$$



Hergert, Bogner, Morris, Schwenk, & Tsukiyama, PR

(2016)

IMSRG: Magnus expansion

- The Magnus expansion: rewriting $U(s) \equiv e^{\hat{\Omega}(s)}$ leads to the following ODE

$$\frac{d\hat{\Omega}(s)}{ds} = \sum_{k=0}^{\infty} \frac{B_k}{k!} \text{ad}_{\hat{\Omega}}^k \eta(s), \quad \Omega(0) = 0 \quad (1)$$

$\text{ad}_{\Omega}^0 \eta(s) = \eta(s)$, $\text{ad}_{\Omega}^k \eta(s) = [\Omega(s), \text{ad}_{\Omega}^{k-1}]$ and B_k are the Bernoulli numbers.

$$\hat{\Omega}(s) = \int_0^s \eta(s_1) ds_1 + \frac{1}{2} \int_0^s ds_1 \int_0^{s_1} ds_2 [\eta(s_1), \eta(s_2)] + \dots \quad (2)$$

The unitarity of $U(s)$ is guaranteed by the anti-hermitian $\Omega(s)$.

- Operator $\hat{O}(s)$ (BCH expansion)

$$\begin{aligned} \hat{O}(s) &= e^{\hat{\Omega}(s)} \hat{O}_0 e^{-\hat{\Omega}(s)} = \sum_{k=0}^{\infty} \frac{1}{k!} \text{ad}_{\hat{\Omega}(s)}^k \hat{O}_0 \\ &= \hat{O}_0 + [\hat{\Omega}(s), \hat{O}_0] + \frac{1}{2!} [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{O}_0]] + \dots \end{aligned} \quad (3)$$

IMSRG: Brillouin Generator

- One-body term:

$$\eta_l^k(s) \equiv \langle \Phi | [\hat{H}(s), \tilde{A}_l^k] | \Phi \rangle \sim \lambda^{1B}, \lambda^{2B} \quad (4)$$

- Two-body term:

$$\eta_{mn}^{kl}(s) \equiv \langle \Phi | [\hat{H}(s), \tilde{A}_{mn}^{kl}] | \Phi \rangle \sim \lambda^{1B}, \lambda^{2B}, \lambda^{3B} \quad (5)$$

$\lambda^{1B}, \lambda^{2B}, \lambda^{3B}$: irreducible 1B, 2B, and 3B density matrices

$$\lambda_j^i = \rho_j^i, \quad (6a)$$

$$\lambda_{kl}^{ij} = \rho_{kl}^{ij} - \hat{A}(\lambda_k^i \lambda_l^j), \quad (6b)$$

$$\lambda_{lmn}^{ijk} = \rho_{lmn}^{ijk} - \hat{A}(\lambda_l^i \lambda_m^j \lambda_n^k + \lambda_l^i \lambda_m^k \lambda_n^j). \quad (6c)$$

- Hierarchy in irreducible DME: $\lambda^{1B} \gg \lambda^{2B} \gg \lambda^{3B} \gg \dots$
- Convergence: $\eta(\infty) = 0$ and $\left. \frac{d\Omega(s)}{ds} \right|_{s=\infty} = 0$.

Brief summary of last talk at MSU

- $\lambda_j^i, \lambda_{kl}^{ij}$ from GCM ($\beta_2, \gamma, \phi_{pn}$)
- NN Interaction: KB3G
- Minimum: $E(^{48}\text{Ti}) = -23.88$ MeV (SM: 23.66 MeV).

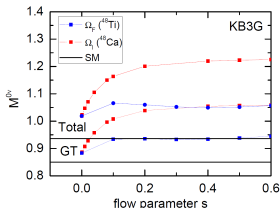
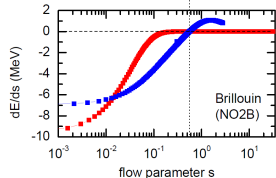
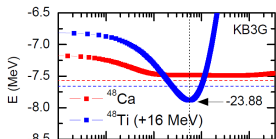
MR-IMSRG(2): does not work for ^{48}Ti ?

NME for $0\nu\beta\beta$

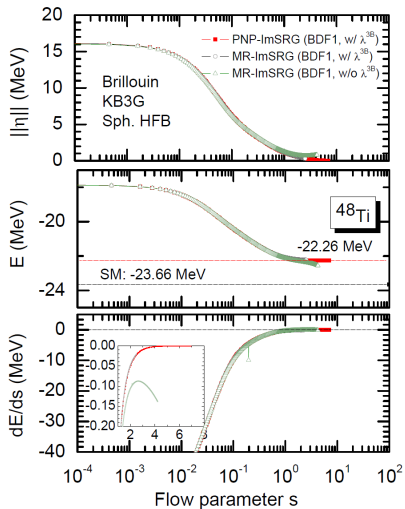
$$M^{0\nu} = \langle 0_F^+ | e^{\hat{\Omega}_F(s)} \hat{O}^{0\nu} e^{-\hat{\Omega}_I(s)} | 0_I^+ \rangle$$

	SM	GCM	ImSRG(2)(Ω_I)	ImSRG(2)(Ω_F)
GT	0.848	0.883	1.058	0.941
Fermi	0.146	0.207	0.233	0.172
Tensor	-0.058	-0.071	-0.067	-0.060
Total	0.936	1.019	1.224	1.053

Missing the irreducible λ^{3B} terms in $\eta^{(2B)}$?



MR-IMSRG(2): Benchmark Calc. from a Spherical State

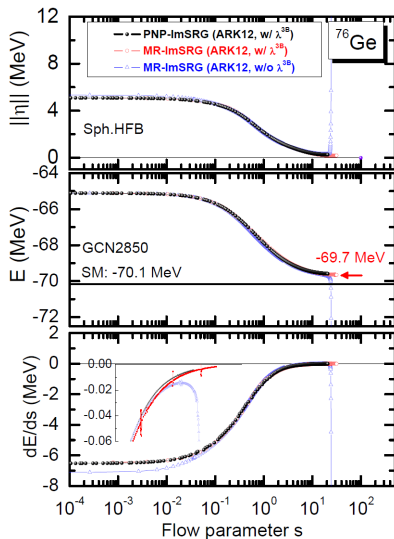


Numerical Details

- Reference state: spherical HFB
- NN Interaction: shell-model KB3G
- Techniques: PNP-ImSRG and MR-ImSRG

Better convergence w/ λ^{3B} in $\eta(s)$

MR-ImSRG(2): Benchmark Calc. from a Spherical State



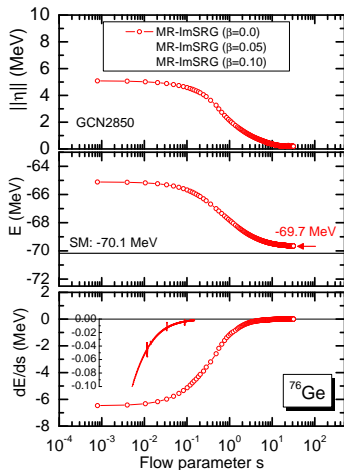
Numerical Details

- Reference state: spherical HFB
- NN Interaction: shell-model GCN2850
- Techniques: PNP-ImSRG and MR-ImSRG

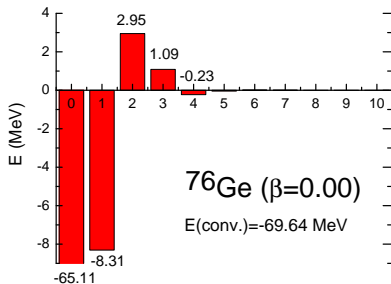
Better convergence w/ λ^{3B} in $\eta(s)$

MR-IMSRG(2): Benchmark Calc. from a Deformed State

$$E(s) = E(0) + \frac{1}{1!} [\hat{\Omega}(s), \hat{H}_0]^{OB} + \frac{1}{2!} [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]^{OB} + \frac{1}{3!} [\hat{\Omega}(s), [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]]^{OB} + \dots$$

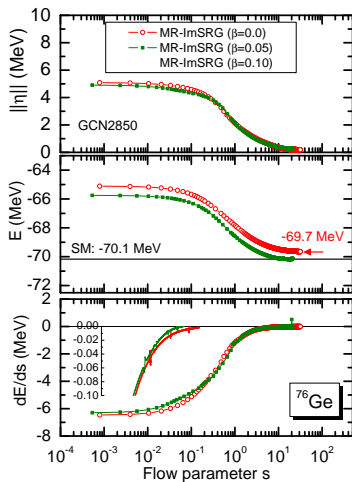


- λ^{3B} to energy ($\beta = 0.00$ case):
 $\sim +1 \times 10^{-3}$ MeV

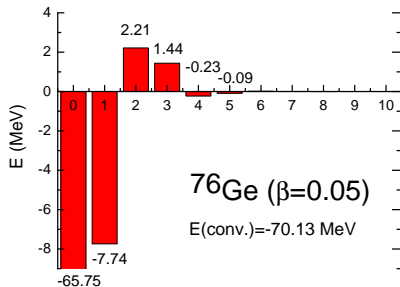


MR-IMSRG(2): Benchmark Calc. from a Deformed State

$$E(s) = E(0) + \frac{1}{1!} [\hat{\Omega}(s), \hat{H}_0]^{OB} + \frac{1}{2!} [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]^{OB} + \frac{1}{3!} [\hat{\Omega}(s), [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]]^{OB} + \dots$$

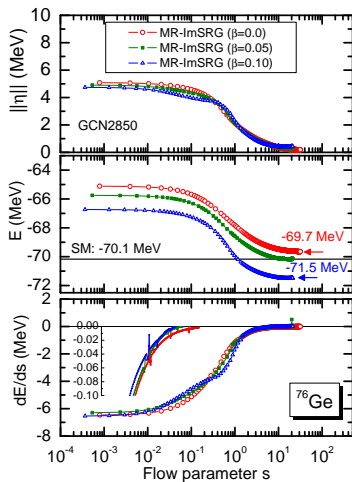


- λ^{3B} to energy ($\beta = 0.05$ case):
 $\sim +4 \times 10^{-3}$ MeV

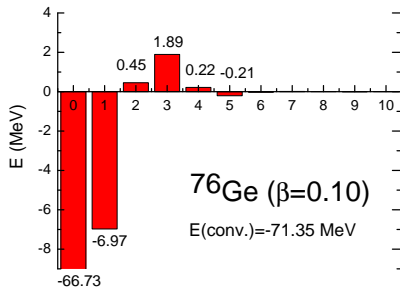


MR-IMSRG(2): Benchmark Calc. from a Deformed State

$$E(s) = E(0) + \frac{1}{1!} [\hat{\Omega}(s), \hat{H}_0]^{OB} + \frac{1}{2!} [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]^{OB} + \frac{1}{3!} [\hat{\Omega}(s), [\hat{\Omega}(s), [\hat{\Omega}(s), \hat{H}_0]]]^{OB} + \dots$$



- λ^{3B} to energy ($\beta = 0.10$ case):
 $\sim +0.15$ MeV. $[\hat{\Omega}(s), \hat{H}_0]^{3B}$?



Summary and outlook

Summary

- The MR-IMSRG(2) based on shell-model interactions **works well for near-spherical Ref. states, but not for large deformed ones.**
- The **more correlation** is included in the Ref. state, the **more important** is the high-rank irreducible density ($\lambda^{2B}, \lambda^{3B}, \dots$).
- Extension of the MR-IMSRG(2) to MR-IMSRG(2*) or MR-IMSRG(3) is needed for deformed nuclei.

Outlook

- MR-IMSRG based on chiral NN interaction with the MR-IMSRG(2*).
- Calculation of the NME for $0\nu\beta\beta$ with the MR-IMSRG(2*).

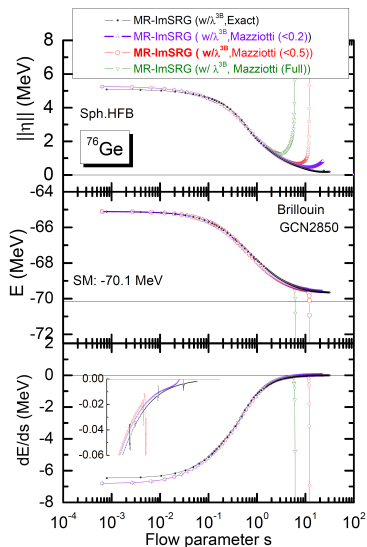
Acknowledgement

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Thanks for your attention



MR-ImSRG: Benchmark Calculation from a Spherical State



Mazziotti's prescription for λ^{3B}

$$\lambda_{stu}^{pqr} \simeq -[n_{stu}^{pqr}]^{-1} \frac{1}{4} \sum_a \hat{A}(\lambda_{st}^{pa} \lambda_{au}^{qr}) \quad (7)$$

with

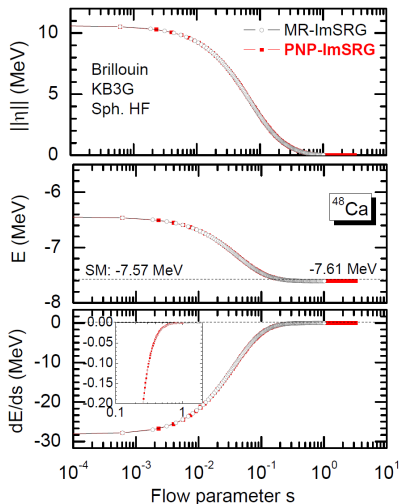
$$n_{stu}^{pqr} = \sum_{a=p,q,r,s,t,u} \lambda_a^a - 3 \quad (8)$$

D. A. Mazziotti, PRA (1999)

Numerical Details

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- NN Interaction: shell-model GCN2850
- Techniques: MR-ImSRG

MR-IMSRG(2): Benchmark Calc. from a Spherical State

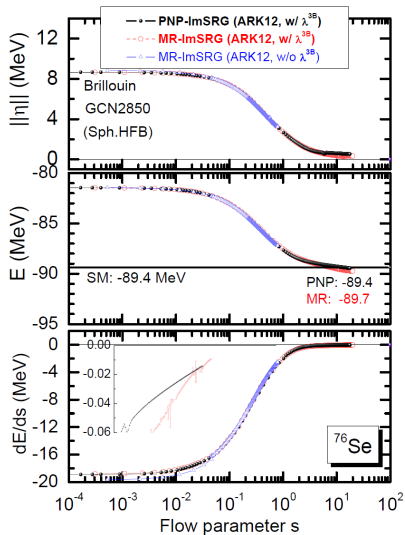


Numerical Details

- Reference state: spherical HF state: $\lambda^{2B}, \lambda^{3B}$ are zero.
- NN Interaction: shell-model KB3G
- Techniques: PNP-ImSRG and MR-ImSRG

MR-IMSRG(2) works for this case

MR-ImSRG(2): Benchmark Calc. from a Spherical State



Numerical Details

- Reference state: spherical HFB
- NN Interaction: shell-model GCN2850
- Techniques: PNP-ImSRG and MR-ImSRG

Better convergence w/ λ^{3B} in $\eta(s)$

MR-ImSRG(2): Benchmark Calc. from a Deformed State

