



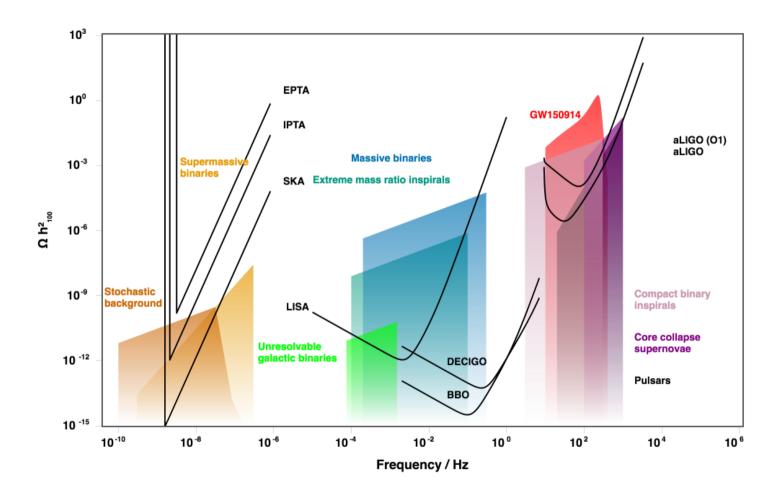


Gravitational waves from a first order electroweak phase transition

David J. Weir, University of Helsinki UMass Amherst, 8 April 2017

https://tinyurl.com/acfi-gws

Lots of sources...



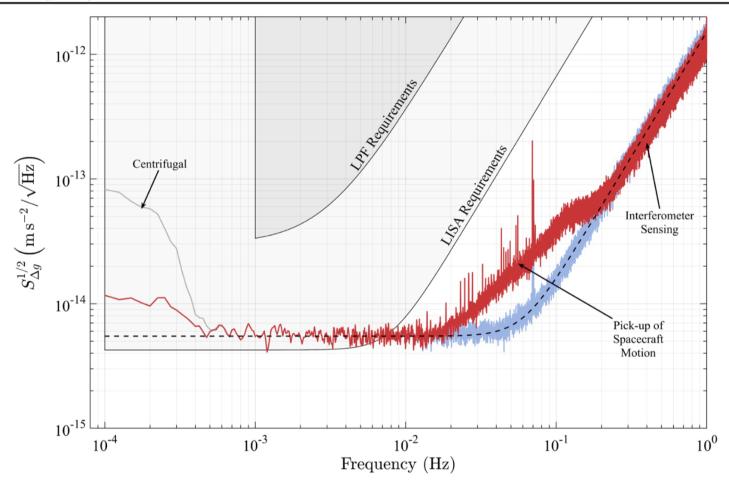
Source: http://rhcole.com/apps/GWplotter/

LISA Pathfinder

PRL 116, 231101 (2016)

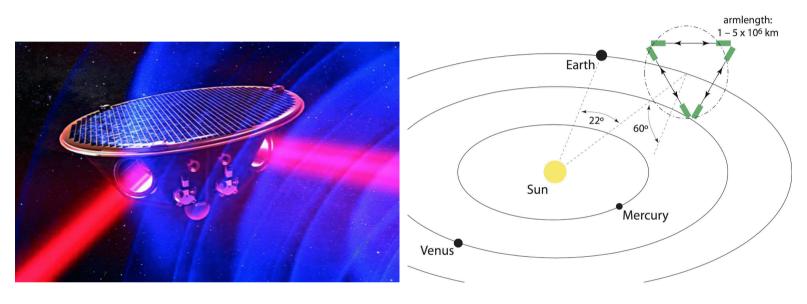
PHYSICAL REVIEW LETTERS

week ending 10 JUNE 2016



Exceeded design expectations by factor of five!

What's next: LISA



- LISA: three arms (six laser links), 2.5 M km separation
 Launch as ESA's third large-scale mission (L3) in (or before) 2034
 Proposal officially submitted earlier this year 1702.00786

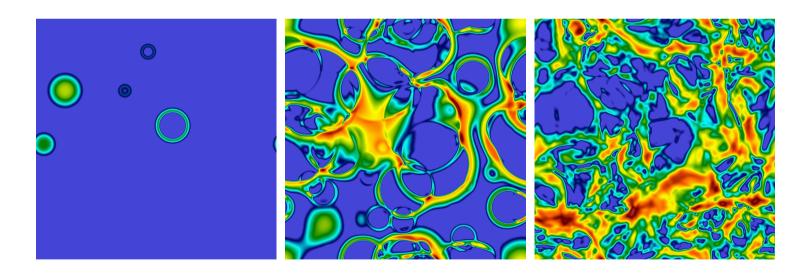
From the proposal:

SI7.2: Measure, or set upper limits on, the spectral shape of the cosmological stochastic GW background

OR7.2: Probe a broken power-law stochastic background from the early Universe as predicted, for example, by first order phase transitions [21] (other spectral shapes are expected, for example, for cosmic strings [22] and inflation [23]). Therefore, we need the ability to measure $\Omega = 1.3 \times 10^{-11} \left(f/10^{-4} \, \text{Hz} \right)^{-1}$ in the frequency ranges $0.1 \, \text{mHz} < f < 2 \, \text{mHz}$ and $2 \, \text{mHz} < f < 20 \, \text{mHz}$, and $\Omega = 4.5 \times 10^{-12} \left(f/10^{-2} \, \text{Hz} \right)^3$ in the frequency ranges $2 \, \text{mHz} < f < 20 \, \text{mHz}$ and $0.02 < f < 0.2 \, \text{Hz}$.

First order thermal phase transition:

- 1. Bubbles nucleate and grow
- 2. Expand in a plasma create shock waves
- 3. Bubbles + shocks collide violent process
- 4. Sound waves left behind in plasma
- 5. Turbulence; expansion



Thermal phase transitions

Standard Model is a crossover

Kajantie et al.; Karsch et al.; ...

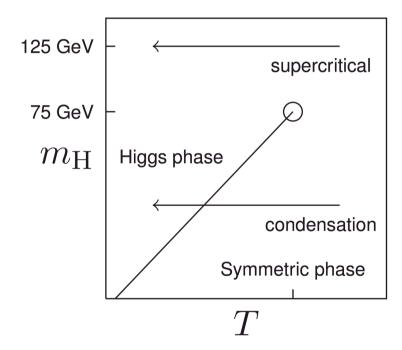
• First order possible in extensions (xSM, 2HDM, ...)

Andersen et al., Kozaczuk et al., Carena et al.,

Bödeker et al., Damgaard et al., Ramsey-Musolf et al.,

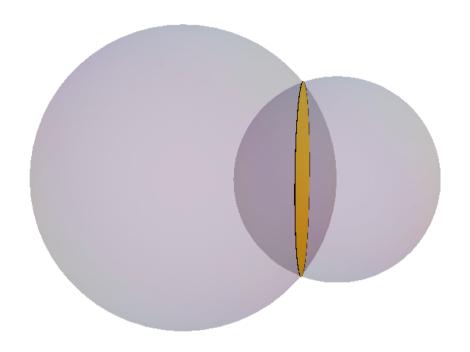
Cline and Kainulainen, ...

- Baryogenesis?
- GW PS ⇔ model information?



What the metric sees at a thermal phase transition

- Bubbles nucleate and expand, shocks form, then:
 - 1. $h^2\Omega_{\phi}$: Bubbles + shocks collide 'envelope phase'
 - 2. $h^2\Omega_{sw}$: Sound waves set up 'acoustic phase'
 - 3. $h^2\Omega_{\text{turb}}$: [MHD] turbulence 'turbulent phase'
- Sources add together to give observed GW power: $h^2\Omega_{\rm GW}\approx h^2\Omega_{\rm d} + h^2\Omega_{\rm sw} + h^2\Omega_{\rm turb}$



Kosowsky, Turner and Watkins; Kamionkowski, Kosowsky and Turner

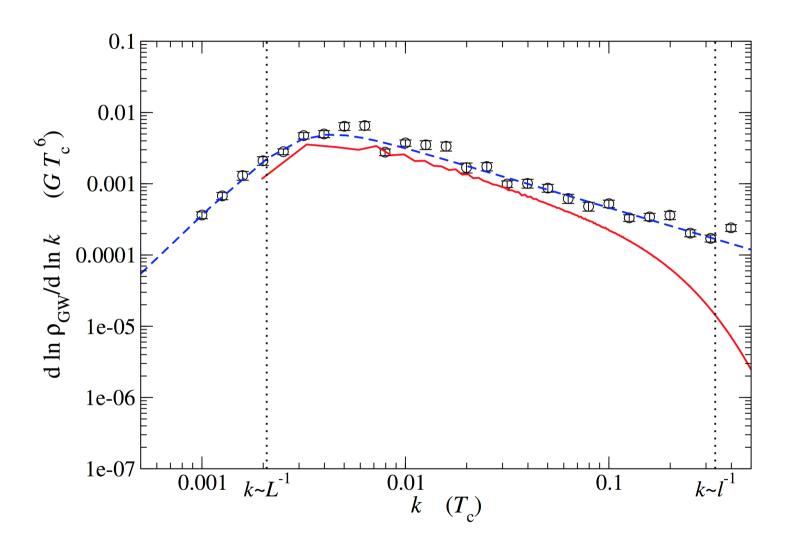
- Thin, hollow bubbles, no fluid
- Stress-energy tensor $\propto R^3$ on wall
- Solid angle: overlapping bubbles → GWs
- Simple power spectrum:
 - One length scale (average radius R_*)
 - Two power laws $(\omega^3, \sim \omega^{-1})$
 - Amplitude
 - \Rightarrow 4 numbers define spectral form

NB: Used to be applied to shock waves (fluid KE), now only use for bubble wall (field gradient energy)

4-5 numbers parametrise the transition:

- α_{T_*} , vacuum energy fraction
- $v_{\rm w}$, bubble wall speed
- κ_{ϕ} , conversion 'efficiency' into gradient energy $(\nabla \phi)^2$
- Transition rate:
 - H_* , Hubble rate at transition
 - β , bubble nucleation rate
 - \rightarrow ansatz for $h^2\Omega_{\phi}$

[only matters for vacuum/runaway transitions]



Coupled field and fluid system

Ignatius, Kajantie, Kurki-Suonio and Laine

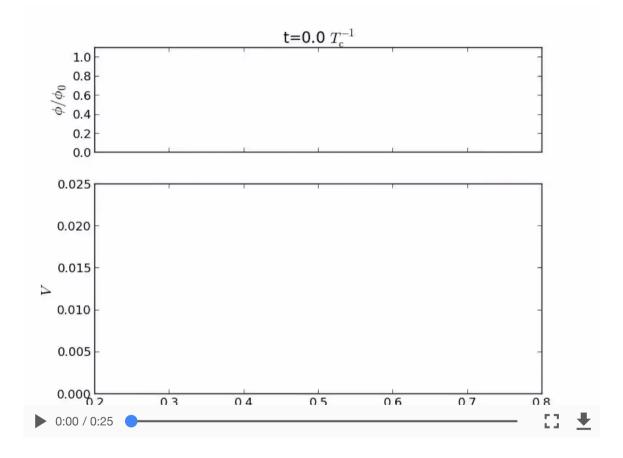
- Scalar ϕ and ideal fluid u^{μ} :
 - Split stress-energy tensor $T^{\mu\nu}$ into field and fluid bits $\partial_{\mu}T^{\mu\nu} = \partial_{\mu}(T^{\mu\nu}_{\text{field}} + T^{\mu\nu}_{\text{fluid}}) = 0$
 - Parameter η sets the scale of friction due to plasma

$$\partial_{\mu}T_{\text{field}}^{\mu\nu} = \tilde{\eta}\frac{\phi^{2}}{T}u^{\mu}\partial_{\mu}\phi\partial^{\nu}\phi \qquad \partial_{\mu}T_{\text{fluid}}^{\mu\nu} = -\tilde{\eta}\frac{\phi^{2}}{T}u^{\mu}\partial_{\mu}\phi\partial^{\nu}\phi$$

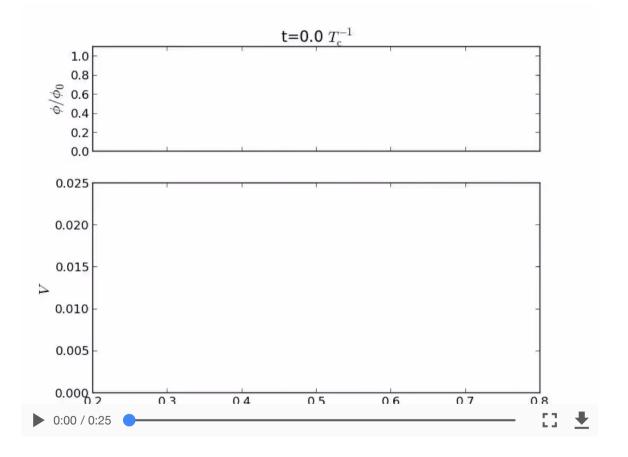
- $V(\phi, T)$ is a 'toy' potential tuned to give latent heat \mathcal{L}
- ullet $\beta \Leftrightarrow$ number of bubbles; $\alpha_{T_*} \Leftrightarrow \mathcal{L}$, $v_{\text{wall}} \Leftrightarrow \tilde{\eta}$

Begin in spherical coordinates: what sort of solutions does this system have?

Velocity profile development: small $\tilde{\eta} \Rightarrow$ detonation (supersonic wall)

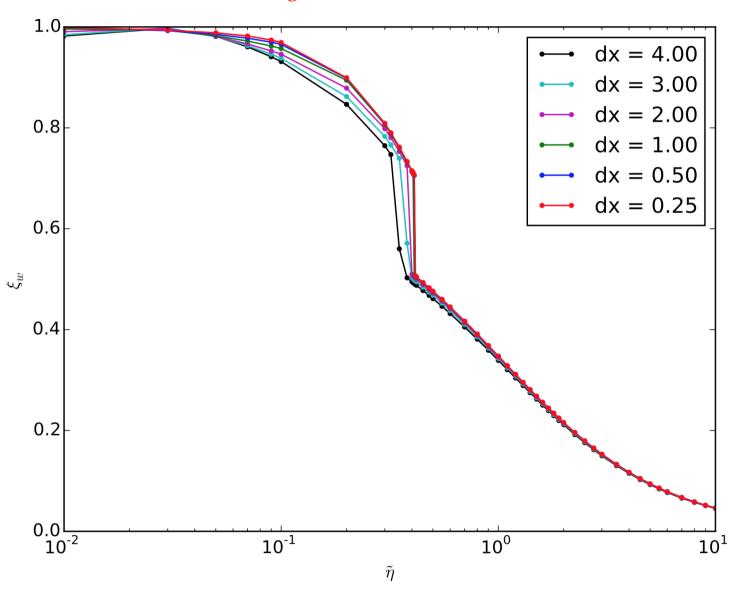


Velocity profile development: large $\tilde{\eta} \Rightarrow$ deflagration (subsonic wall)

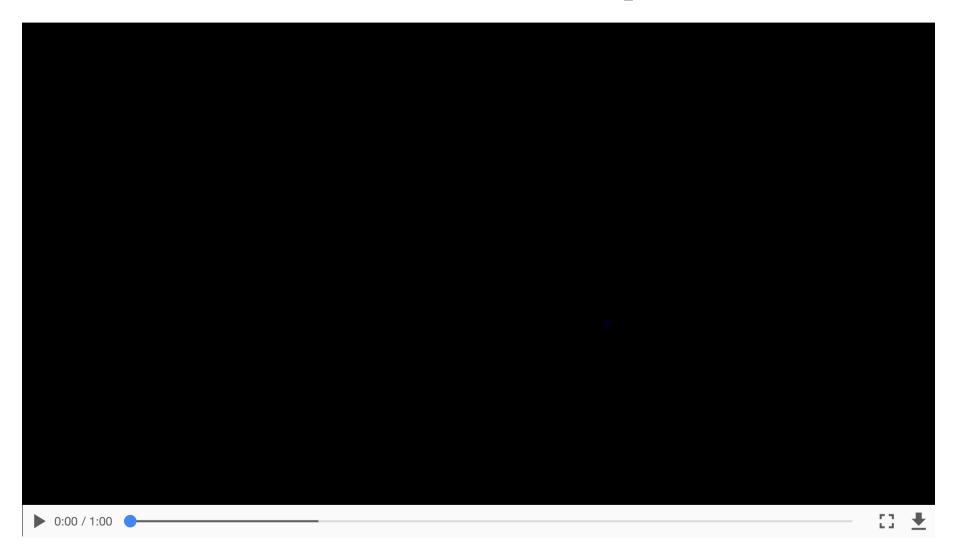


$v_{\rm w}$ as a function of $\tilde{\eta}$

Cutting [Masters dissertation]

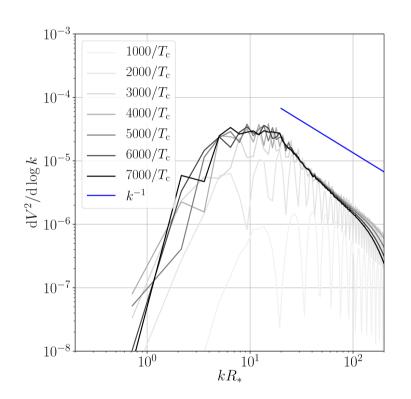


Simulation slice example

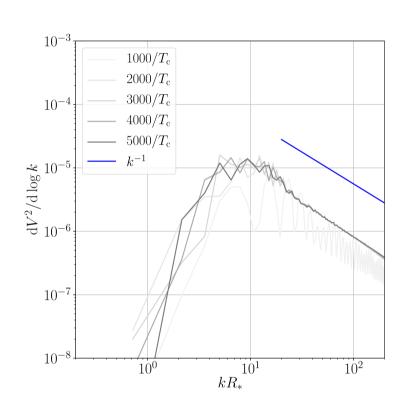


Velocity power spectra and power laws

Fast deflagration



Detonation

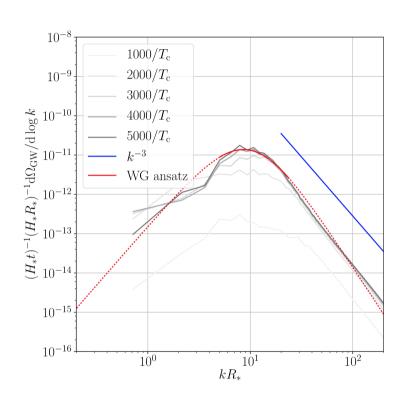


- Weak transition: $\alpha_{T_*} = 0.01$
- Power law behaviour above peak is between k^{-2} and k^{-1}
- "Ringing" due to simultaneous nucleation, unimportant

GW power spectra and power laws

Fast deflagration

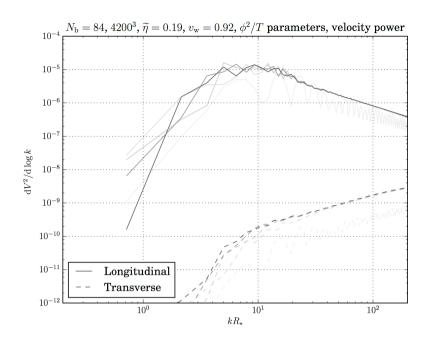
Detonation



- Causal k^3 at low k, approximate k^{-3} or k^{-4} at high k
- Curves scaled by *t*: source until turbulence/expansion

 \rightarrow power law ansatz for $h^2\Omega_{\rm sw}$

Transverse versus longitudinal modes – turbulence?



- Short simulation; weak transition (small α): linear; most power in longitudinal modes \Rightarrow acoustic waves, turbulent
- Turbulence requires longer timescales R_*/U_{f}
- Plenty of theoretical results, use those instead

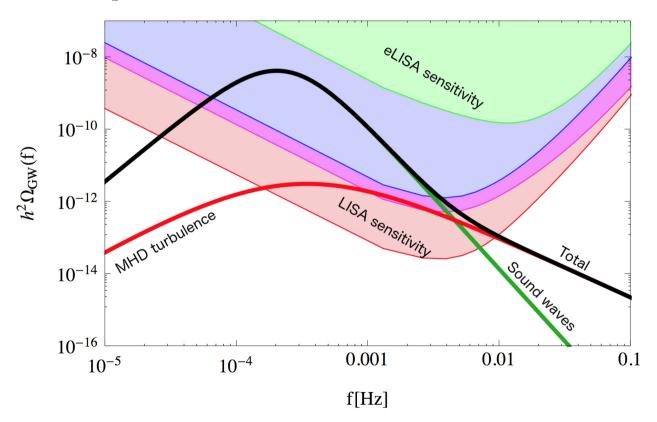
Kahniashvili et al.; Caprini, Durrer and Servant; Pen and Turok; ...

 \rightarrow power law ansatz for $h^2\Omega_{\rm turb}$

Putting it all together - $h^2\Omega_{\rm gw}$ 1512.06239

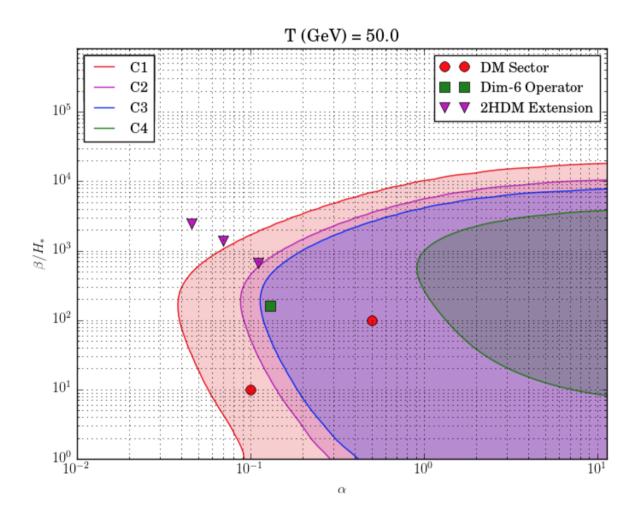
- Three sources, $\approx h^2 \Omega_{\phi}$, $h^2 \Omega_{\rm sw}$, $h^2 \Omega_{\rm turb}$
- Know their dependence on T_* , α_T , $v_{\rm W}$, β Espinosa, Konstandin, No, Servant
- Know these for any given model, predict the signal...

(example,
$$T_* = 100 \text{GeV}$$
, $\alpha_{T_*} = 0.5$, $v_w = 0.95$, $\beta/H_* = 10$)



Putting it all together - physical models to GW power spectra

Model
$$\longrightarrow$$
 (T_* , α_{T_*} , v_w , β) \longrightarrow this plot

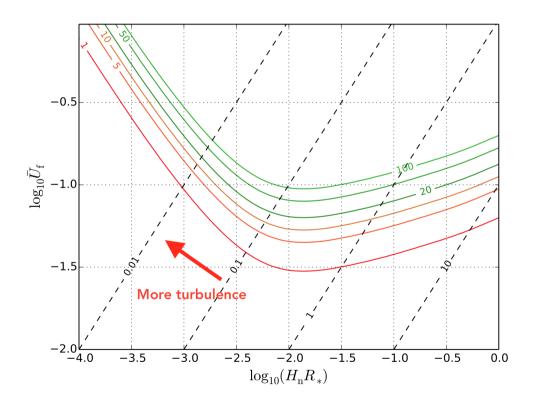


... which tells you if it is detectable by LISA (see 1512.06239)

Detectability from acoustic waves alone

- In many cases, sound waves dominant
- Parametrise by RMS fluid velocity $U_{\rm f}$ and bubble radius R_* (quite easily obtained Espinosa, Konstandin, No and Servant)

Sensitivity plot:



The pipeline



- 1. Choose your model (e.g. SM, xSM, 2HDM, ...)
- 2. Dim. red. model Kajantie et al.
- 3. Phase diagram (α_{T_*}, T_*) ; lattice: Kajantie et al.
- 4. Nucleation rate (β) ; lattice: Moore and Rummukainen
- 5. Wall velocities (v_{wall})
 Moore and Prokopec; Kozaczuk
- 6. GW power spectrum Ω_{gw}
- 7. Sphaleron rate

Very leaky, even for SM!

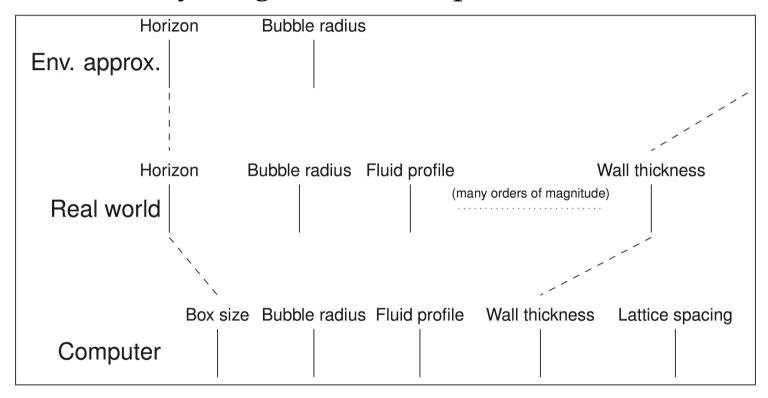
Questions, requests or demands...

- Turbulence
 - MHD or no MHD?
 - Timescales $H_*R_*/U_{\rm f} \sim 1$, sound waves and turbulence?
 - More simulations needed?
- Interaction with baryogenesis
 - Competing wall velocity dependence of BG and GWs?
 - Sphaleron rates in extended models?
- The best possible determinations for xSM, 2HDM, Σ SM, ...
 - What is the phase diagram?
 - Nonperturbative nucleation rates?

Implementation extra slides

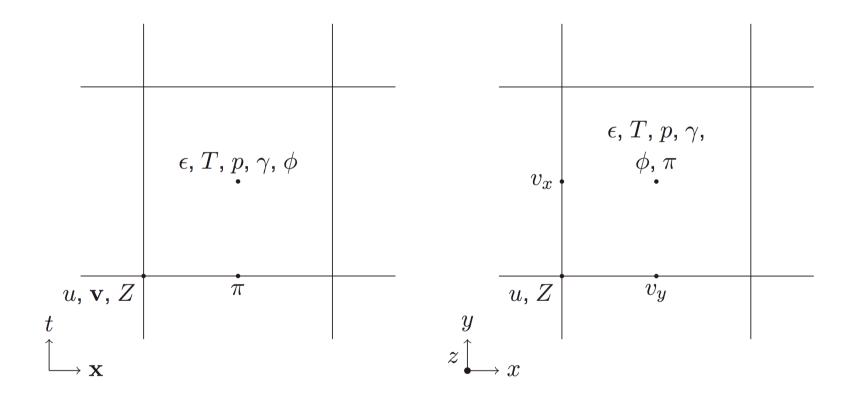
Dynamic range issues

- Most realtime lattice simulations in the early universe have a single [nontrivial] length scale
- Here, many length scales important



• Recently completed simulations with 4200^3 lattices, $\delta x = 2/T_c \rightarrow \text{approx 1M CPU hours each (17.6M total)}$

Implementation: special relativistic hydrodynamics Different things live in different places...



With this discretisation, evolution is second-order accurate!

Extra, extra slides [click here]